

Primordial Black Holes

and their effect on the cosmic ionization history

Massimo Ricotti

ricotti@astro.umd.edu

University of Maryland

Collaborators:

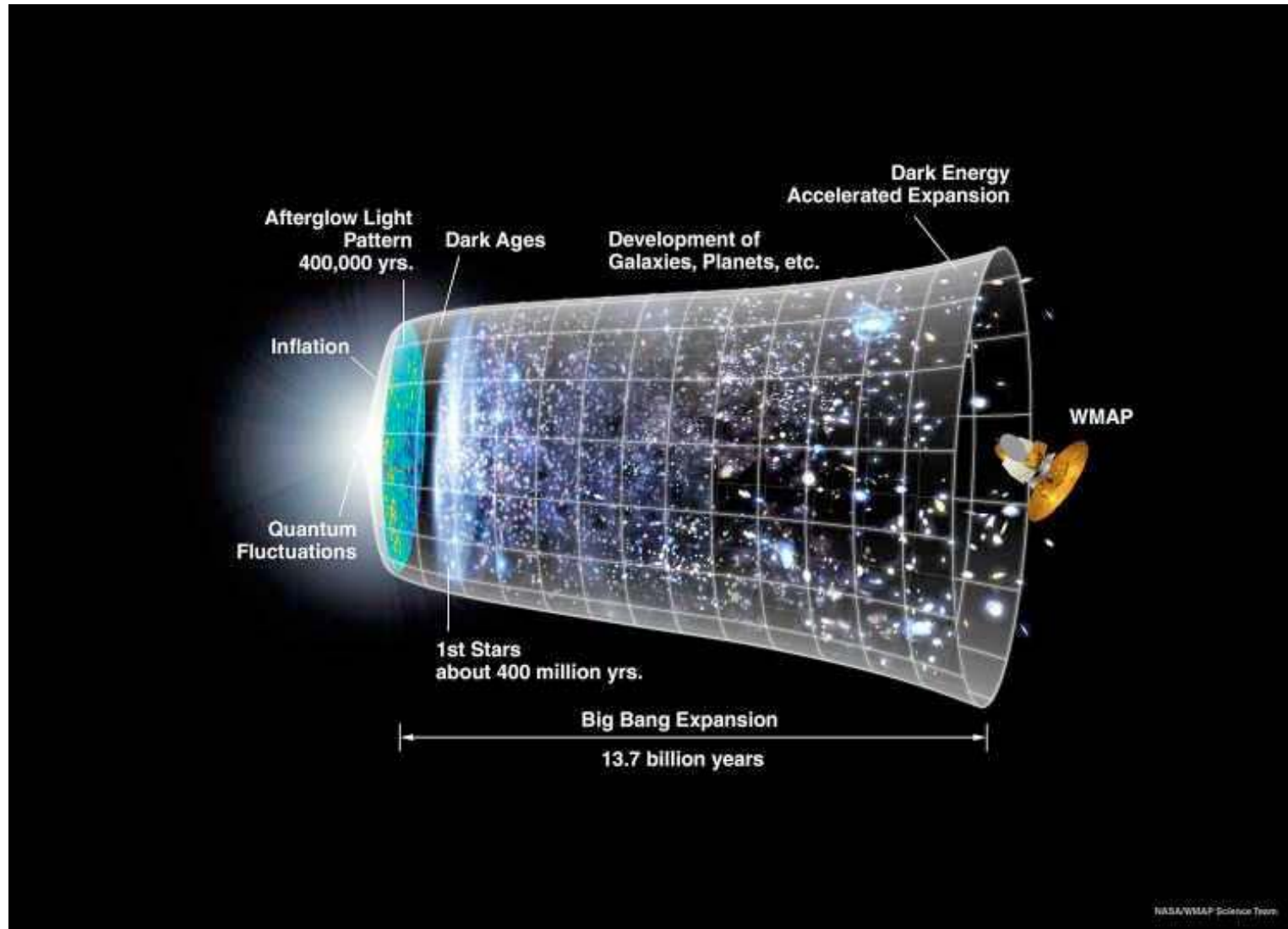


Jerry Ostriker (Princeton), Katie Mack (Princeton)

Outline

- What are PBHs and why we do we care?
- Modelling accretion onto PBHs:
 - Growth of DM halos, modified Bondi solutions
 - PBHs motions and feedback processes
- Effect on the ionization history
- Effect on the CMB, new obs. constraints and Implications

Hot Big Bang theory



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1. Collapsed radiation (relativistic matter)
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3. Form in quasi-linear regime: 50%
4. Tiny collapsed fraction during rad. era may produce all the dark matter!

Newtonian limit easy to understand: During radiation era

$$c_s = c/\sqrt{3} \text{ hence } R_J = R_{\text{Sch}} = R_h$$

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Collapsed fraction depends on the power spectrum of
 initial density fluctuations and the cosmic equation of state:

- $\delta_{cr} = \delta_{cr}(w)$ where $P = w$ is the cosmic EOS
- $\delta(M, z) \sim \exp[-(\delta_{cr}/2 \delta(M, z))^2]$ (assuming Gaussian fluctuations)

Radiation is redshifted away, PBHs are not:

$$\rho_{\text{pbh}} = \frac{(1 + z_f)}{(1 + z_{\text{eq}})}$$

$$f_{\text{pbh}} = \frac{\rho_{\text{pbh}}}{\rho_{\text{m}}} = \frac{M_{\text{pbh}}}{1 M} \frac{1}{10^9}^{1/2}$$

Radiation is redshifted away, PBHs are not:

$$\rho_{\text{pbh}} = \frac{\rho_{\text{rad}}}{(1 + z_f)/(1 + z_{\text{eq}})}$$

$$f_{\text{pbh}} = \frac{\rho_{\text{pbh}}}{\rho_{\text{dm}}} = \frac{M_{\text{pbh}}^{-1/2}}{10^{-9} M}$$

● Example:

- During QCD phase transition at $t = 10^{-5}$ sec
- $M_{\text{pbh}} = M_h = 1 M$
- if $f_{\text{pbh}} = 10^{-9}$: all the dark matter is made of PBHs

Some models for PBH formation

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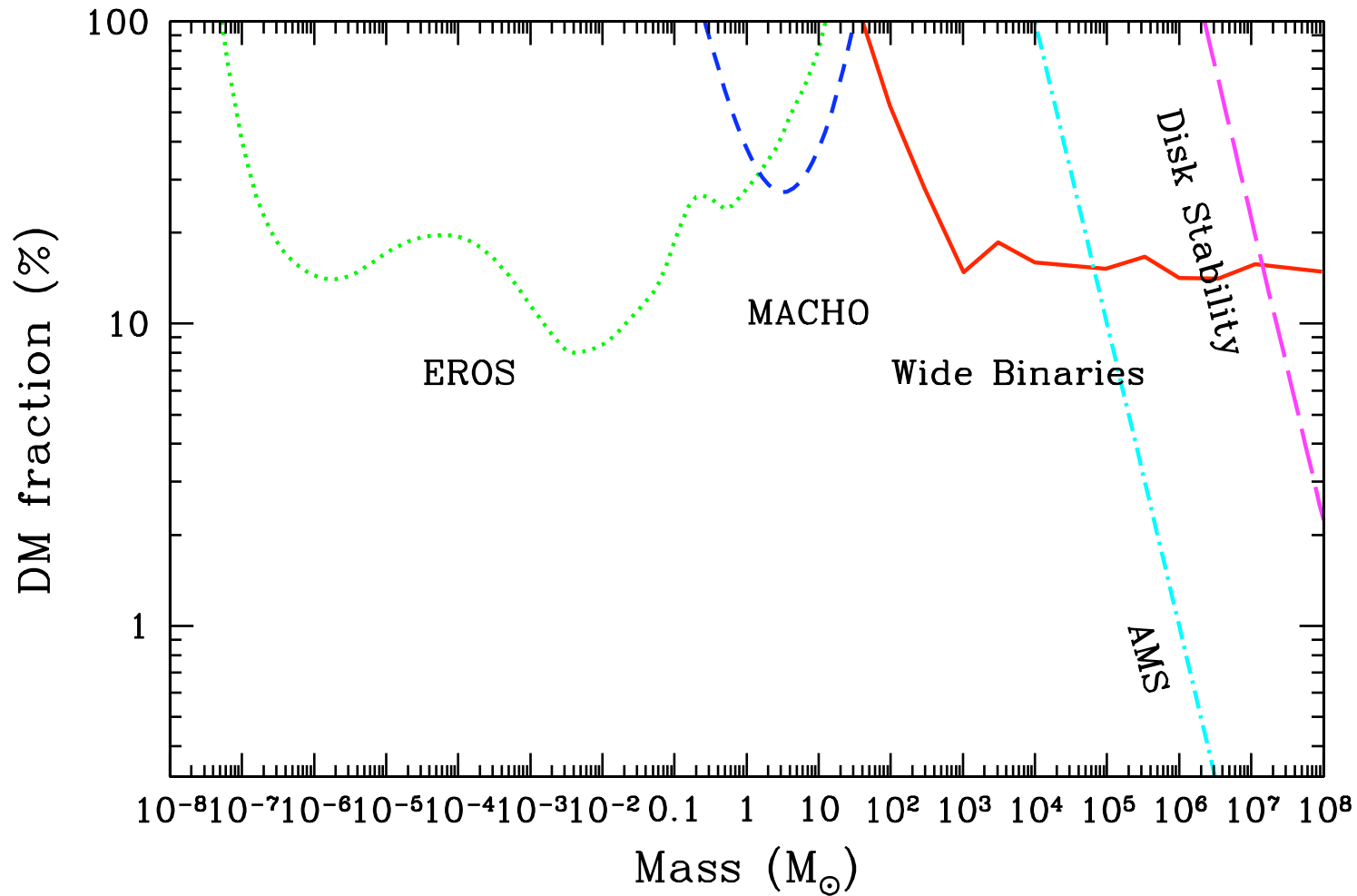
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3. Collapse of cosmic string loops (*e.g.*, Polnarev & Zemboricz 88; Hawking 89; Brandenberger & Wichoski 98)
4. Bubble collisions (*e.g.*, Crawford & Schramm 82; La & Steinhardt 89)
5. Collapse of domain walls (Berezin et al 83; Ipser & Sikivie 84; Rubin et al. 00)

Do PBHs exist?

- PBHs with mass $< 10^{15}$ g evaporate in $t < t_H$ (Hawking 1975)
 - Abundance of PBHs with mass $1 \text{ g} < M < 10^{15} \text{ g}$ is $< 10^{-20} - 10^{-22}$ (e.g., Carr 2003)
- More massive PBHs are poorly constrained:
 - They may constitute the bulk of the dark matter
 - MACHO collaboration: 20% of Milky-Way halo is in compact objects with $M = 0.1 - 1 M_\odot$ (but 2000 result, non confirmed by later data)



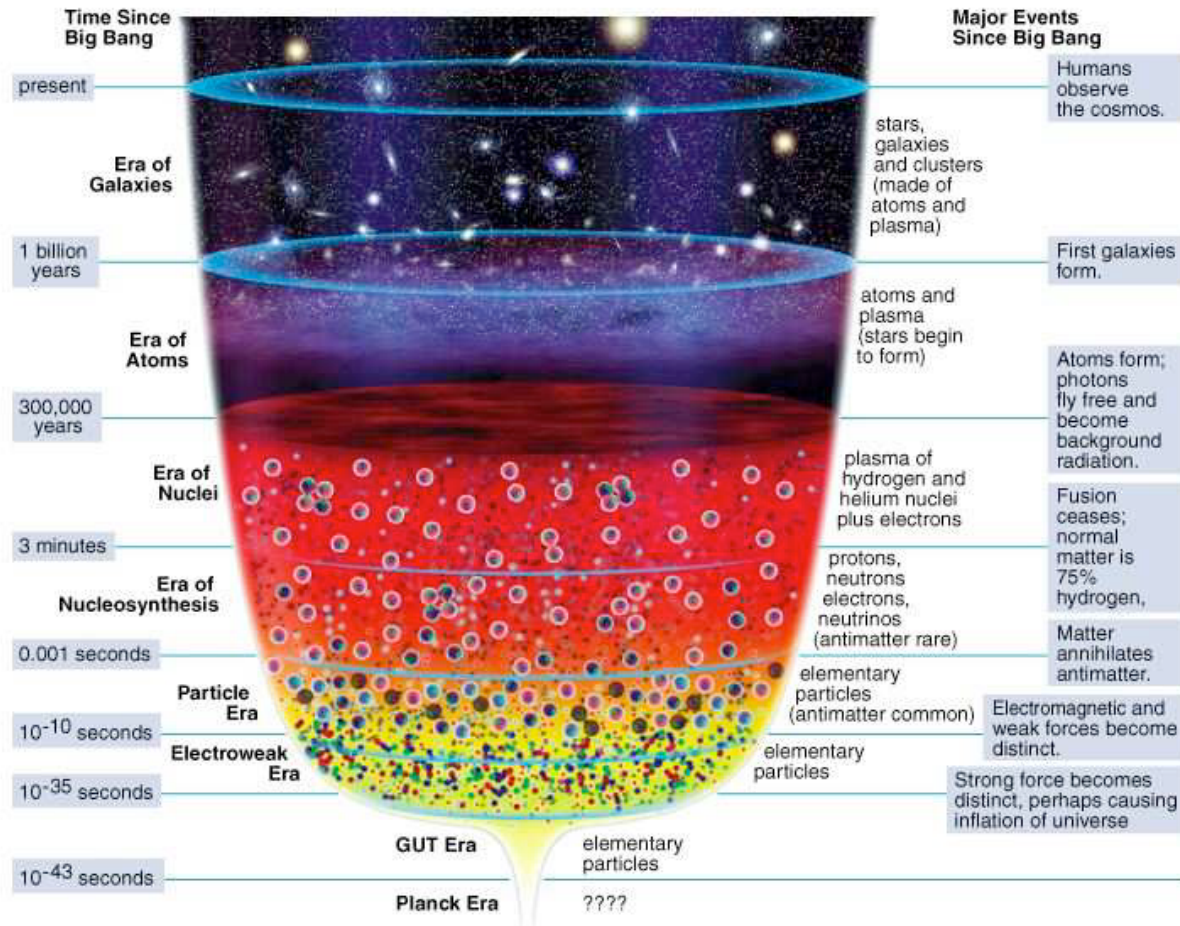
Refs: MACHO collaboration [e.g., Alcock et al. (1998, 2000, 2001); Hamadache et al. 2006]; EROS collaboration; Lacey &



Netriker 85; Moore 93; Carr 94; Afshordi, McDonald & Spergel

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2. The dark matter can be made of PBHs
3. Produce MACHOS, IMBH and ULXs ?
4. Seeds for supermassive Black Holes?



Physics of PBHs accretion: Outline

Gas accretion onto PBHs produce X-rays and affect the ionization history of the IGM

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3. Proper motion: gas and dark matter are not perfectly coupled (Silk damping)
4. Angular momentum of accreted gas and dark matter
5. Feedback processes (global and local radiative feedbacks)

1. Clothing Dark Matter Halo

1. PBHs seed accumulation of dark halo ($f_{\text{pbh}} < 1$)
2. The gas accretion rate onto the PBH is increased!

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- growth of order unity during radiation era
- $M_h = 3M_{\text{pbh}} \frac{1+z}{1000}^1$
- Self-similar secondary infall solution (*e.g.*, Bertschinger 1985)
- Truncated power-law density profile with $\gamma = 2.25$
- Truncation at $r_h = r_{\text{ta}}/3$

2. Bondi type accretion

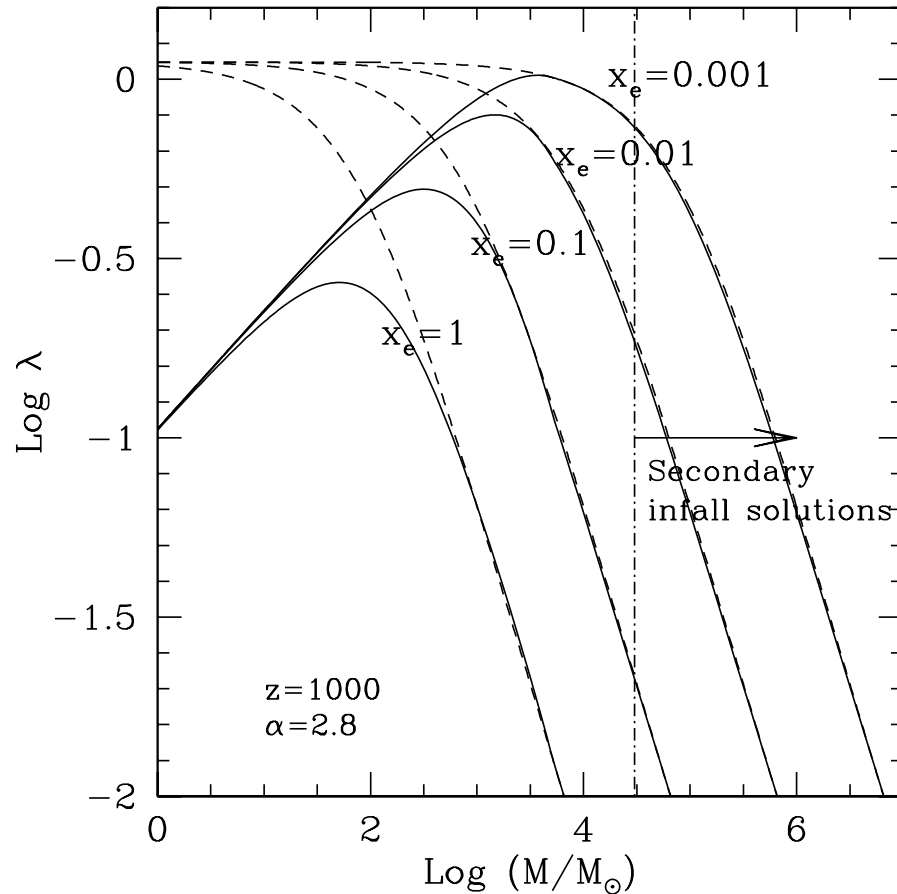
- Spherical accretion ()
- Steady-state ($M < 2 \cdot 10^4 M_\odot$,)
- Viscosity (Compton drag)

$$\frac{dv}{dt} = \frac{4}{3} \frac{x_e \tau U_{\text{cmb}}}{m_p c} v = -v$$

- Hubble expansion ($M > 2 \cdot 10^4 M_\odot$,)
- Clothing dark halo (power-law density profile,)

Ref: Ricotti (2007)

Spherical accretion rate



- point mass potential (dashed curves)
- dark halo potential (solid curves)

3. Relative motion of PBHs and gas

$$\dot{M}_g \sim \frac{M^2}{v_{\text{eff}}^3}$$

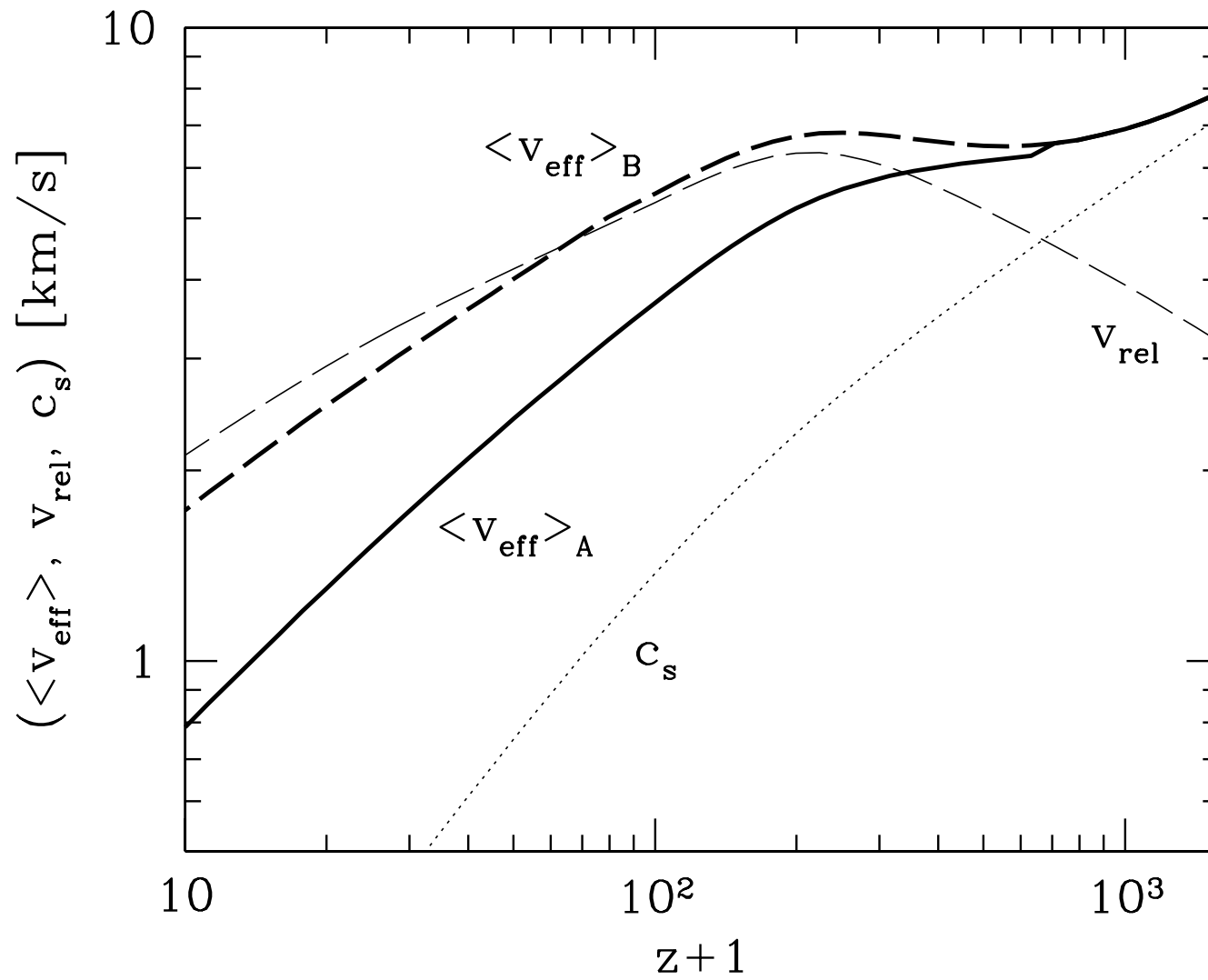
where $v_{\text{eff}} = (v_{\text{rel}}^2 + c_s^2)^{1/2}$

1. Linear regime: Silk damping (i = baryons, dark matter)

$$V_i^2 = \frac{1.2}{2} \frac{m^2 H^2}{2} \int_0^\infty P_i(k) w_s^2(k, a) w_l^2(k, r_0) dk,$$

$$V_i^2 = \frac{1.2}{2} \frac{m^2 H^2}{2} \int_0^\infty P_i(k) w_s^2(k, a) [1 - w_l^2(k, r_0)] dk.$$

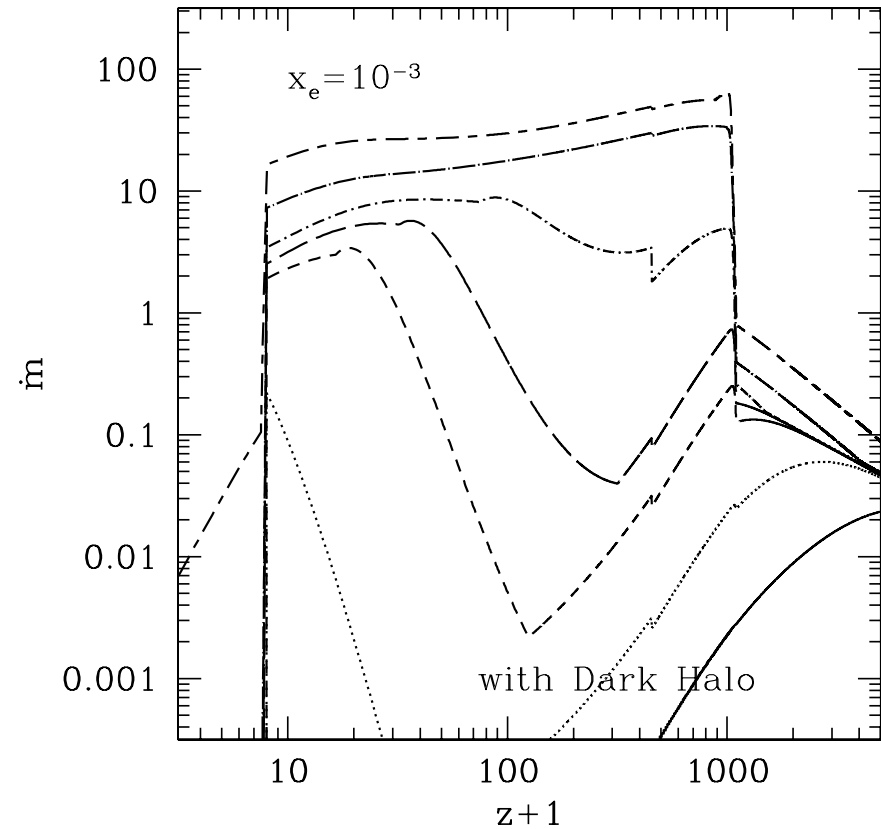
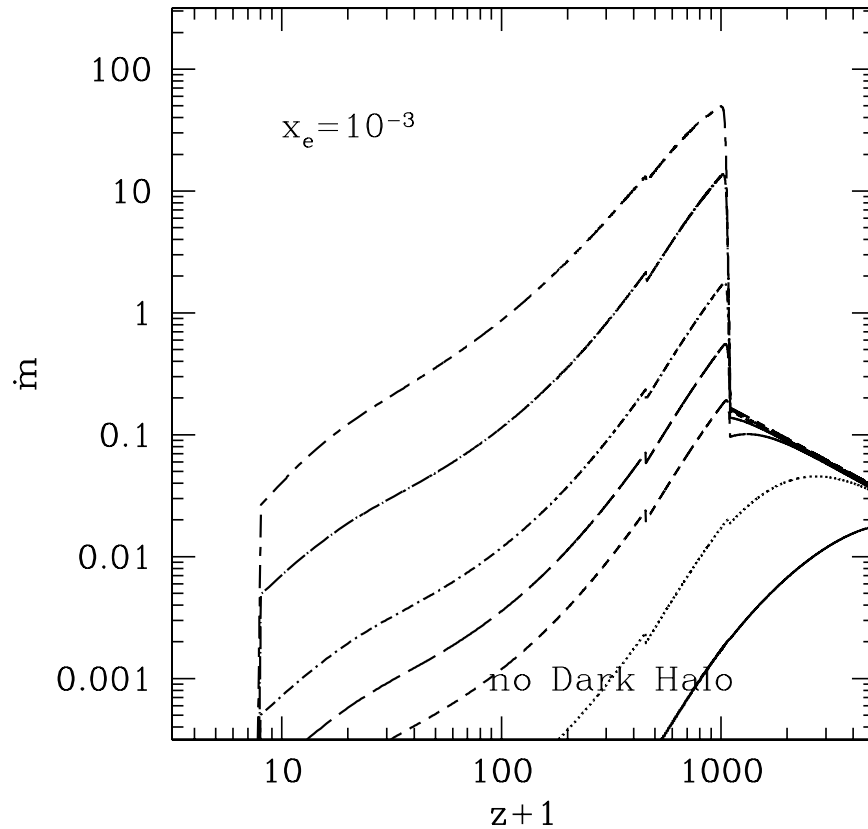
2. Non-linear regime: capture by mini-halos



Angular momentum of accreted material

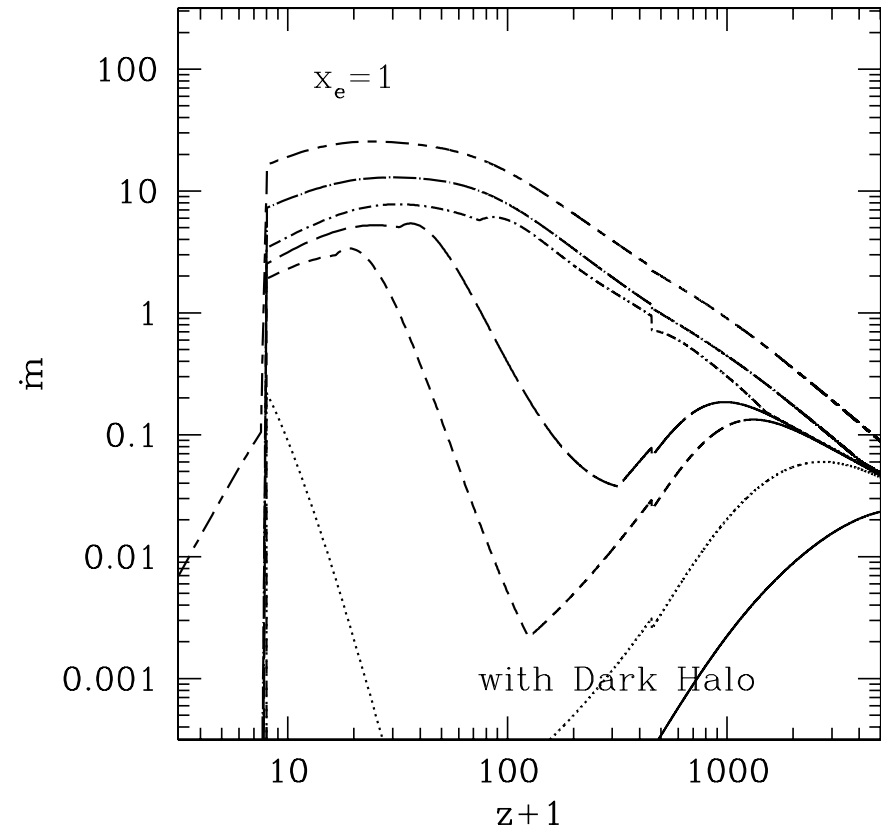
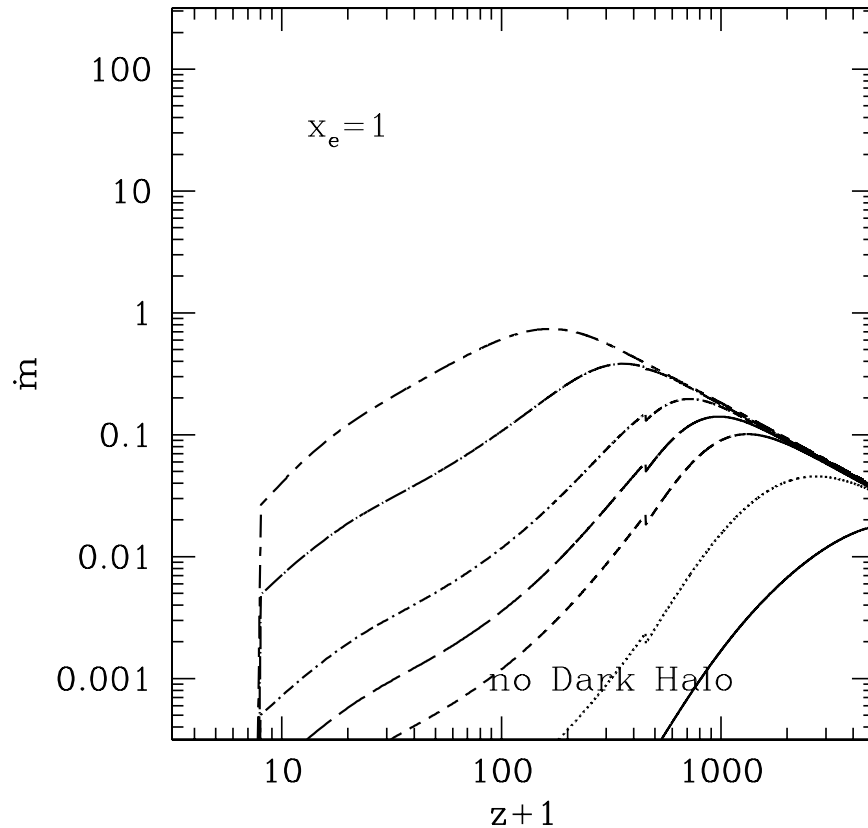
- Dark matter:
 - Quasi-spherical accretion
 - Ang. momentum sufficiently large to avoid direct accretion of DM into PBH
- Gas:
 - Spherical accretion for $M < 500 M_{\odot}$
 - Compton drag reduces further ang. momentum (Loeb 93; Umemura et al 93)

4. Accretion rate neglecting feedback



curves from bottom to top refer to masses of PBHs from $0.1 M$ to $10^5 M$ (factor of 10 spacing).

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Accretion Luminosity

Define dimensionless luminosity $l = L/L_{\text{Ed}}$ and accretion rate $\dot{m} = \dot{M}/\dot{M}_{\text{Ed}}$, then:

$l = \eta \dot{m}$, where η is the radiative efficiency

We assume:

- $l = 0.01\dot{m}^2$ if $\dot{m} < 1$ (spherical accretion)
- $l = f_{\text{duty}}(0.1\dot{m}) < f_{\text{duty}}$ if $\dot{m} > 1$ (thin disk)

5. Feedback processes

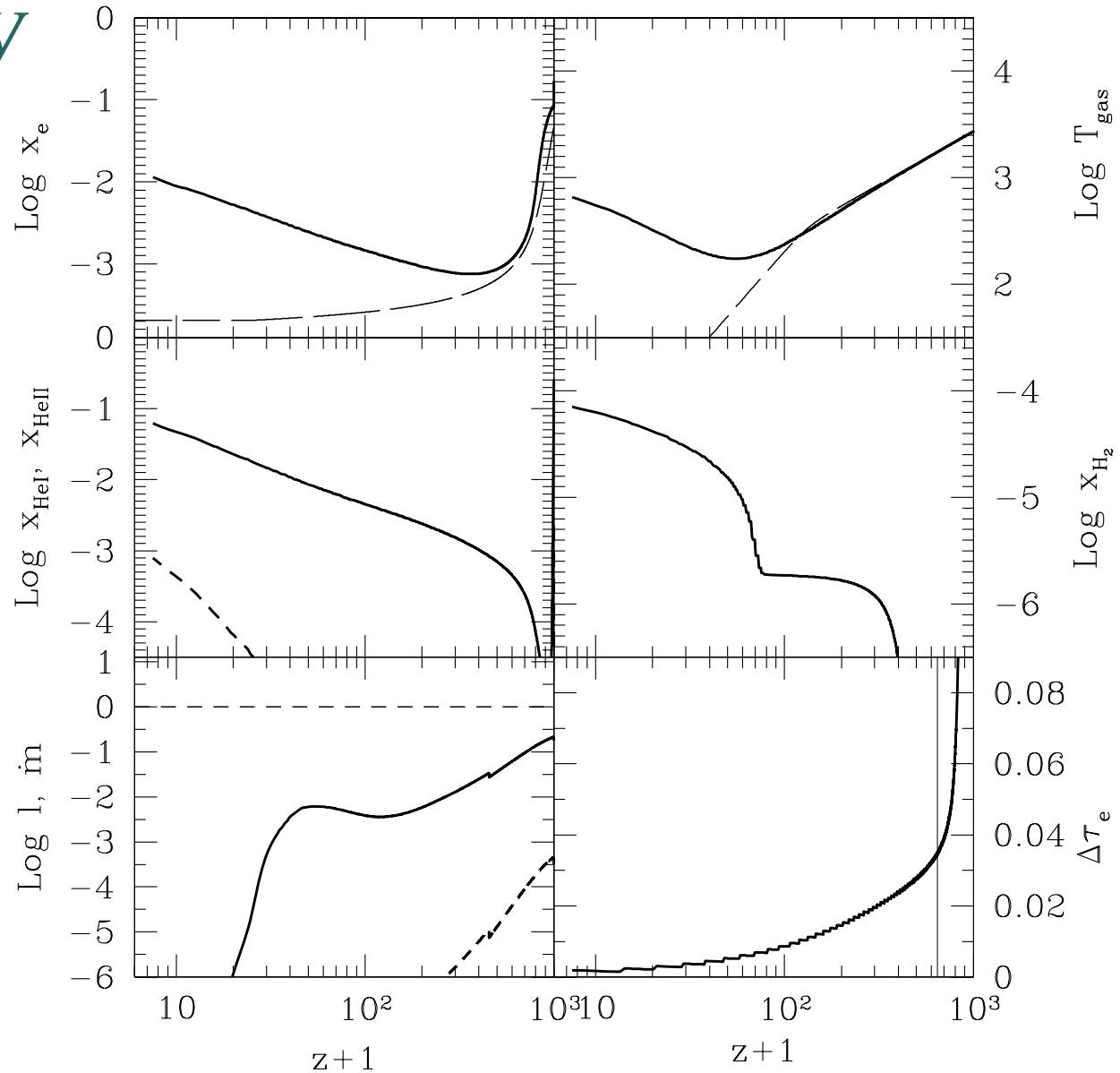
- Local feedbacks (typically negligible)
 - Size of H II region with respect to Bondi radius:
 - In most cases $r_{\text{H II}}/r_{\text{B}} < 1$
 - If $r_{\text{H II}}/r_{\text{B}} > 1$

$$I_t = \frac{I}{1 + t_{\text{off}}/t_{\text{on}}} = \frac{I}{1 + (r_{\text{H II}}/r_{\text{B}})^{1/3}} = f_{\text{duty}} I$$

- Temperature of H II region: $T_{\text{H II}} \gg T_{\text{cmb}}$
- Global feedback (X-ray heating):
Iterative semi-analytic code (Ricotti & Ostriker 04)

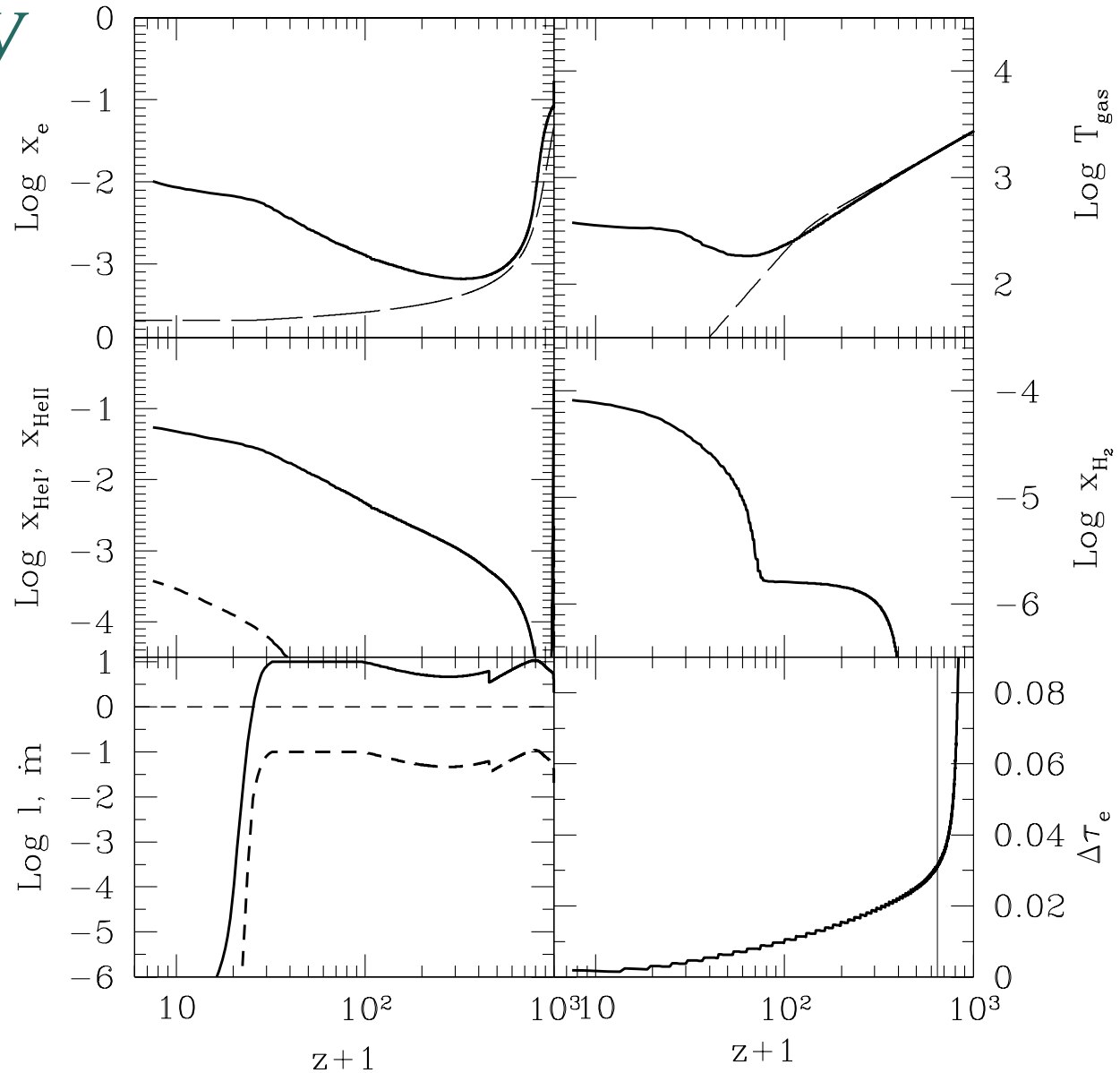
Cosmic ionization, thermal, and chemical history

Simulations:
 $M_{\text{pbh}} = 100,$
 $f_{\text{pbh}} = 10^{-4}$



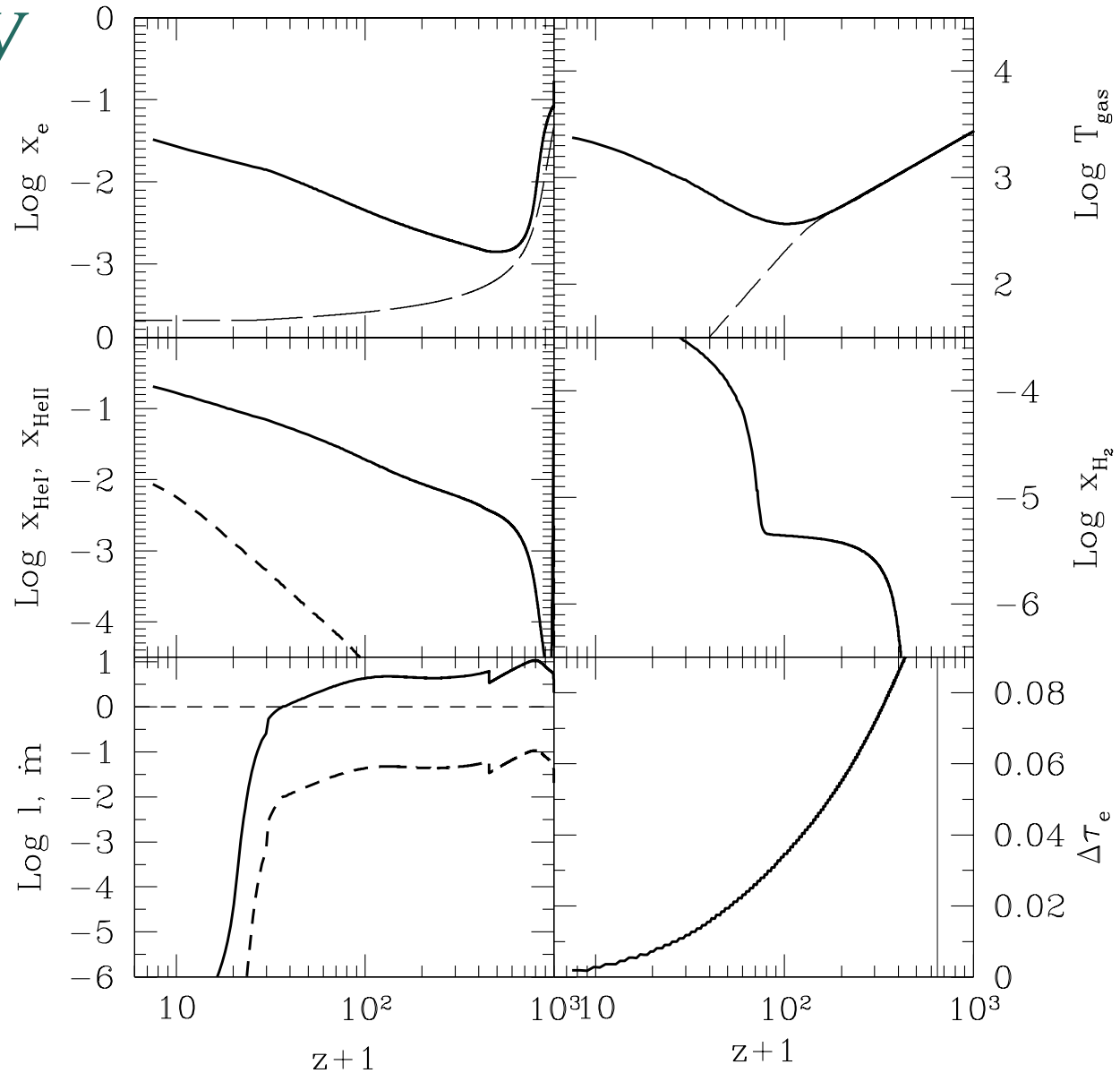
Cosmic ionization, thermal, and chemical history

Simulations:
 $M_{\text{pbh}} = 1000,$
 $f_{\text{pbh}} = 10^{-7}$

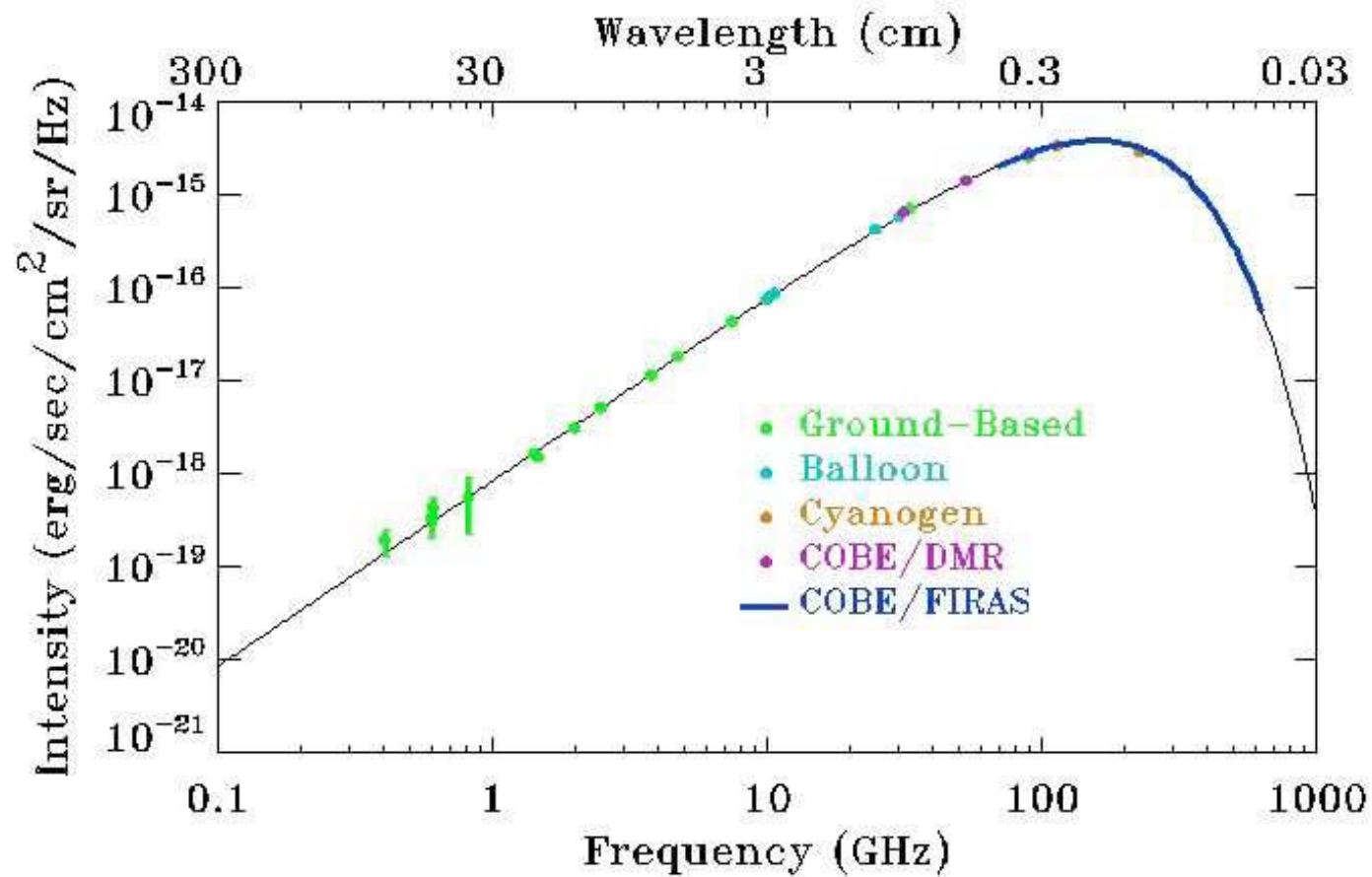


Cosmic ionization, thermal, and chemical history

Simulations:
 $M_{\text{pbh}} = 1000,$
 $f_{\text{pbh}} = 10^{-6}$



Spectrum of the CMB



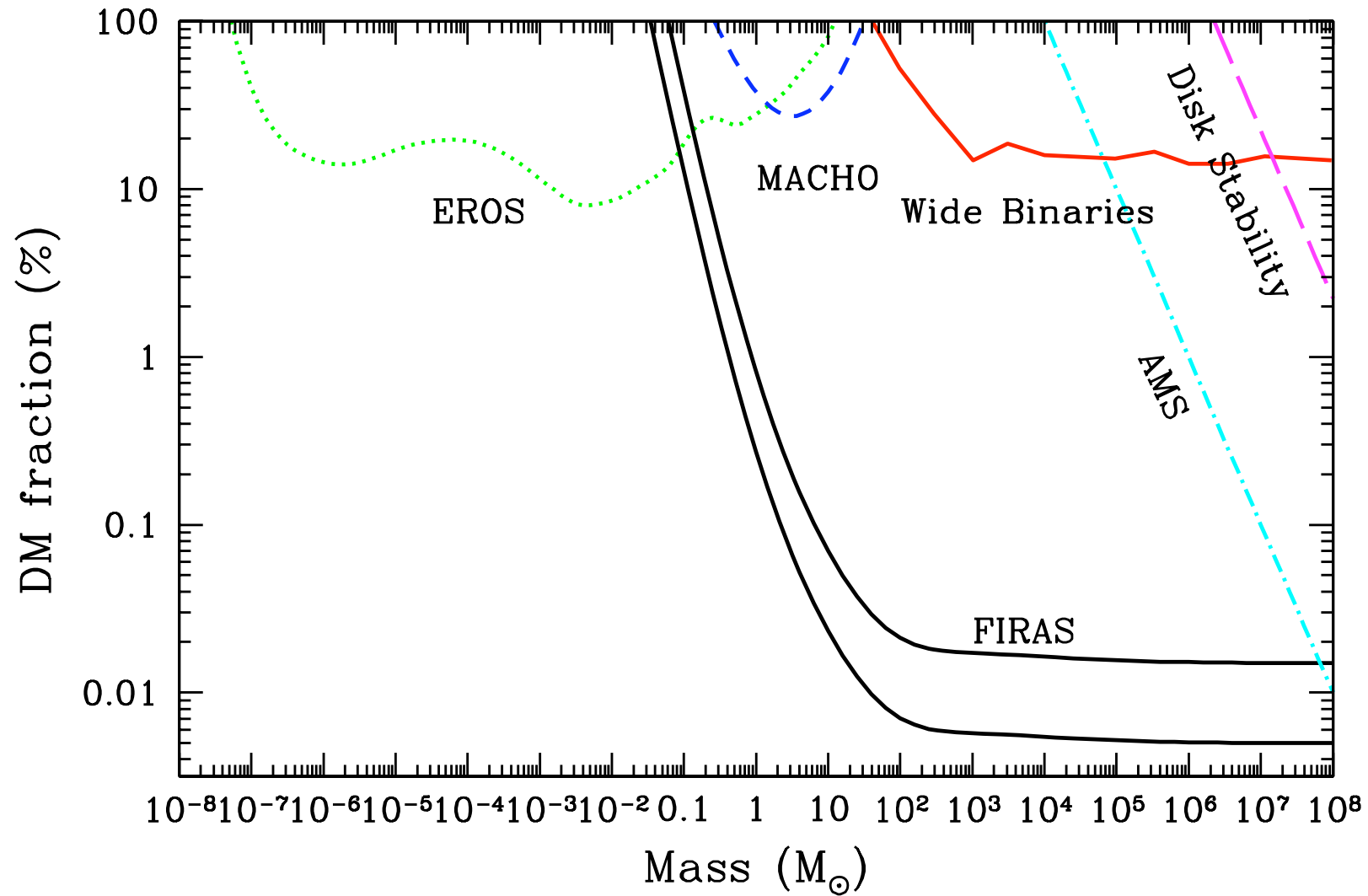
CMB spectral distortions

- $M_{\text{pbh}} < 10 M_{\odot}$ weakly affected by Compton drag even before recombination
- FIRAS $y = 1.5 \times 10^{-5}$ at 95% confidence
- 3 phases: $y = y_1 + y_2 + y_3$
 1. $z_{\text{rec}} < z < z_{\text{eq}}$: all energy injected absorbed by gas
 2. $z_{\text{dec}} < z < z_{\text{rec}}$: fraction of energy absorbed by gas
 3. $z < z_{\text{dec}}$: Compton heating becomes negligible

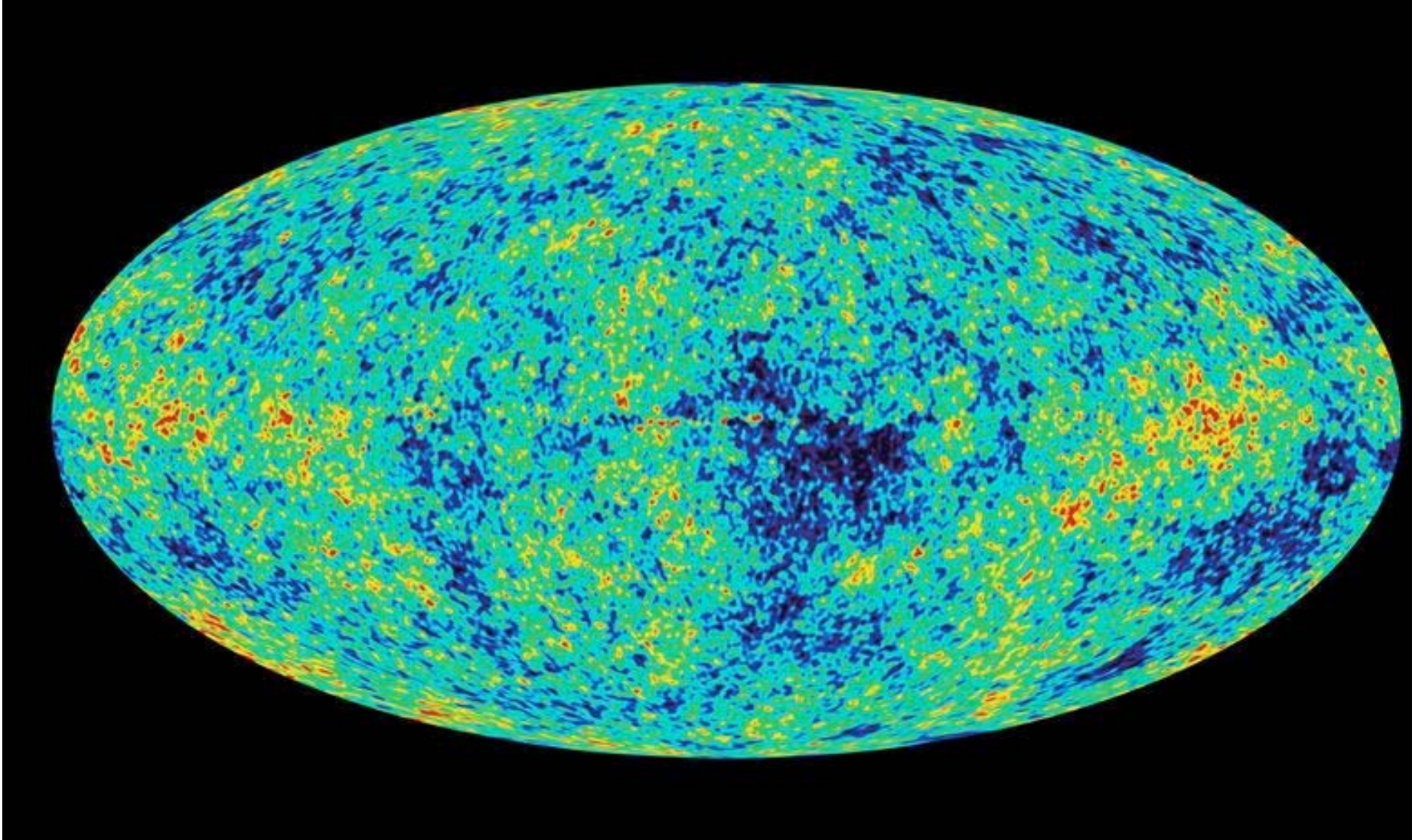
Constrain on maximum energy injection imposed by:

$$y = \frac{1}{4U(z_{\text{eq}})} \int_{z_{\text{eq}}}^{z_{\text{rec}}} \frac{dz}{aH(z)} \frac{dU(z)}{dt} \approx 1.5 \times 10^{-5}$$

Constraints on f_{PBH}

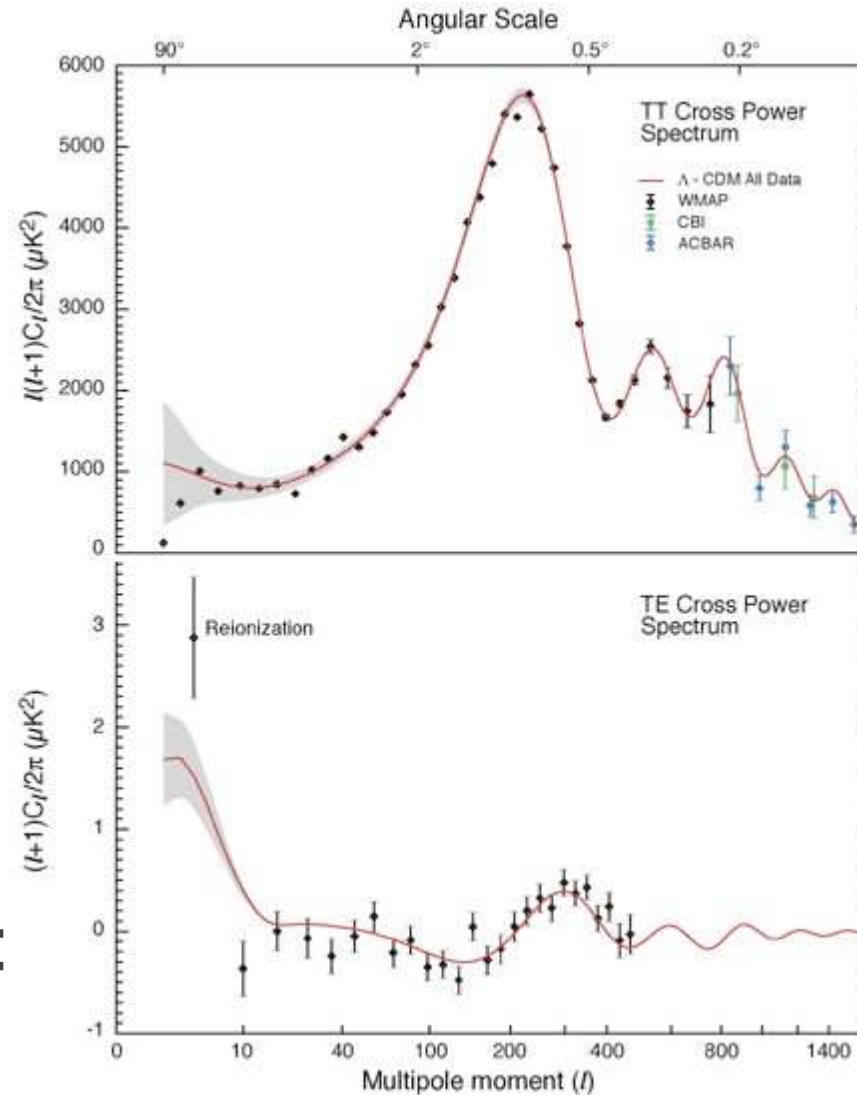


CMB anisotropies: WMAP



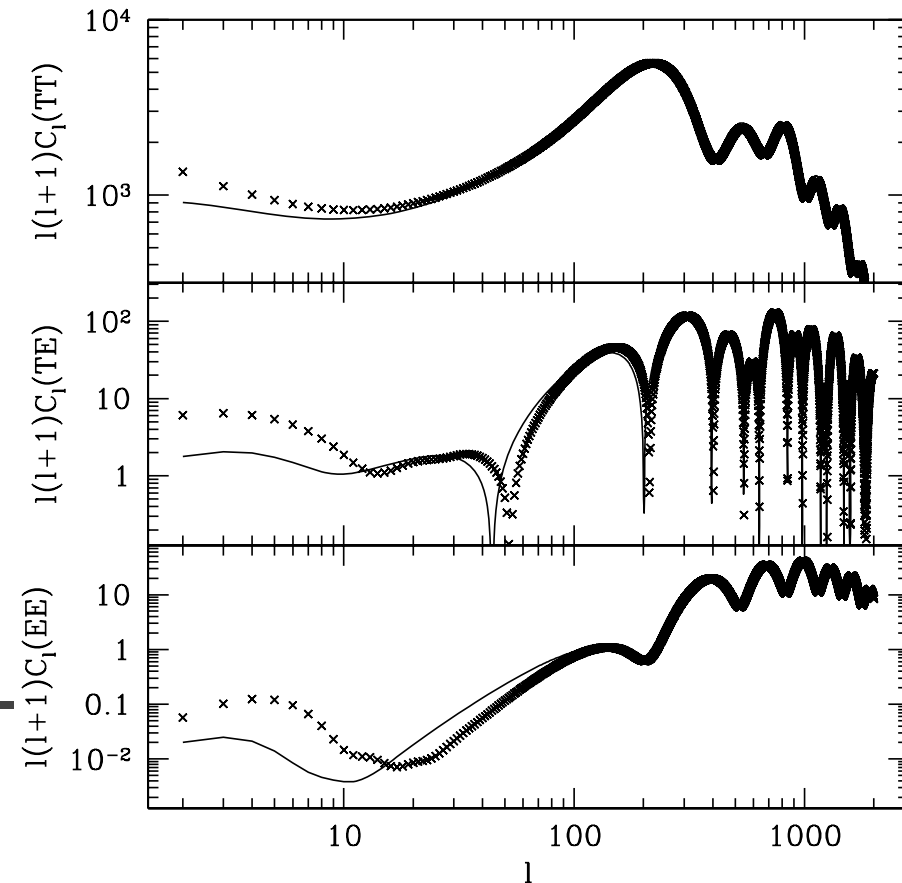
CMB anisotropies: spectrum

- Cosmological parameters:
 $\Omega_m = 0.24, \Omega_b = 0.046, h = 0.73$
- Initial density perturbations:
 $\sigma_8 = 0.74, n_s = 0.95$ (little power at small scales!)
- Ionization history:
 $z_{\text{rei}} = 12, \tau = 0.1$



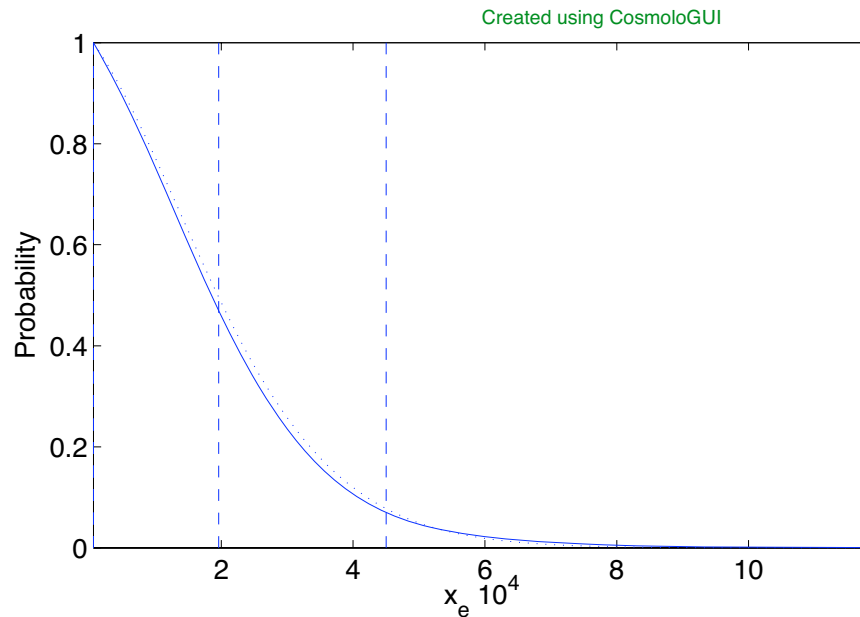
Effects of PBHs on CMB anisotropies

- Affects recombination
- Complementary and uncorrelated to reionization effects
- Affect small angular scales: $l > 10$



Modified recombination history

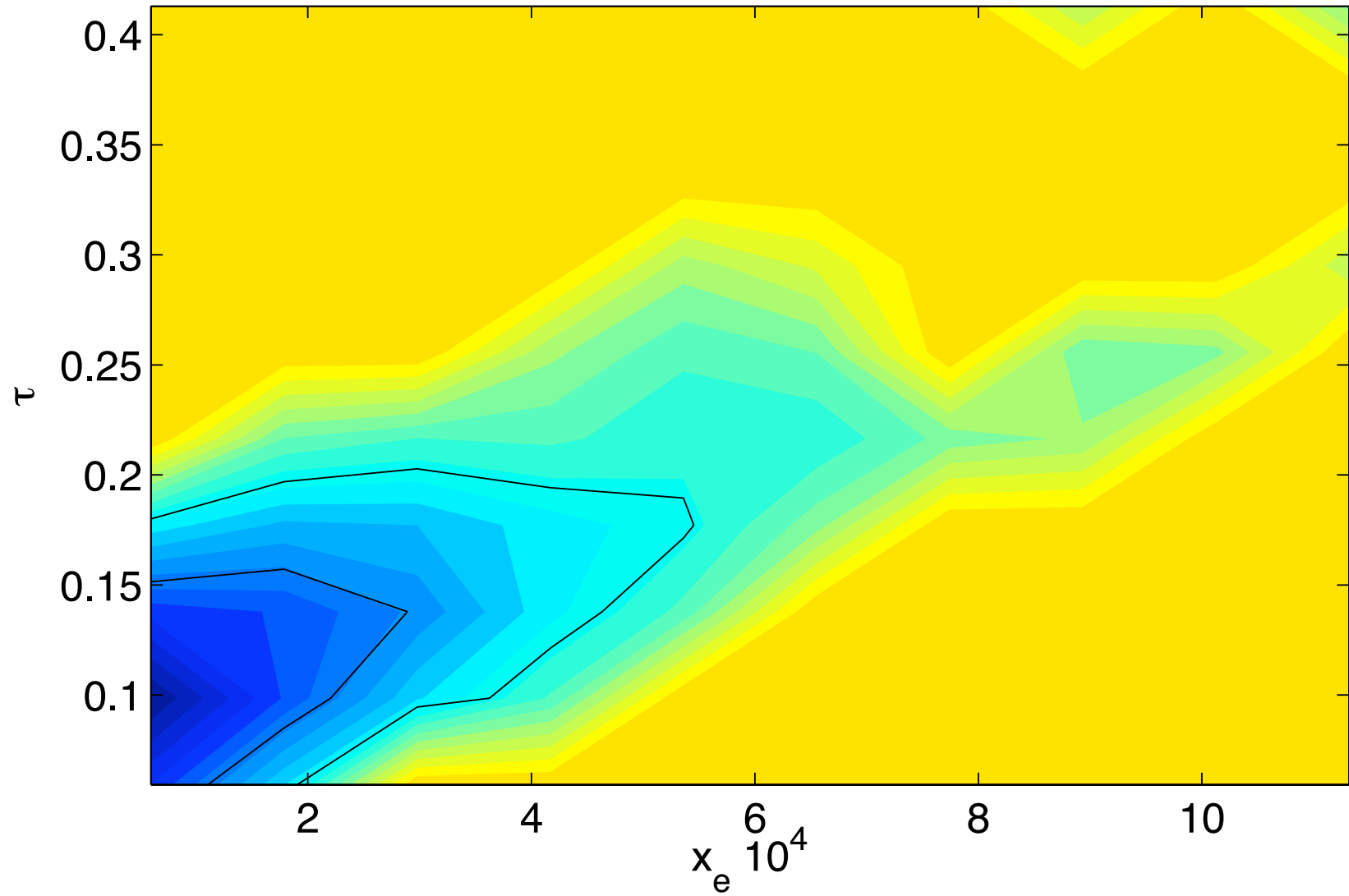
$$x_e(z) = x_{e,\text{rec}}(z) + \min \left(x_{e0} \frac{1+z}{1000}, 0.1 \right)$$



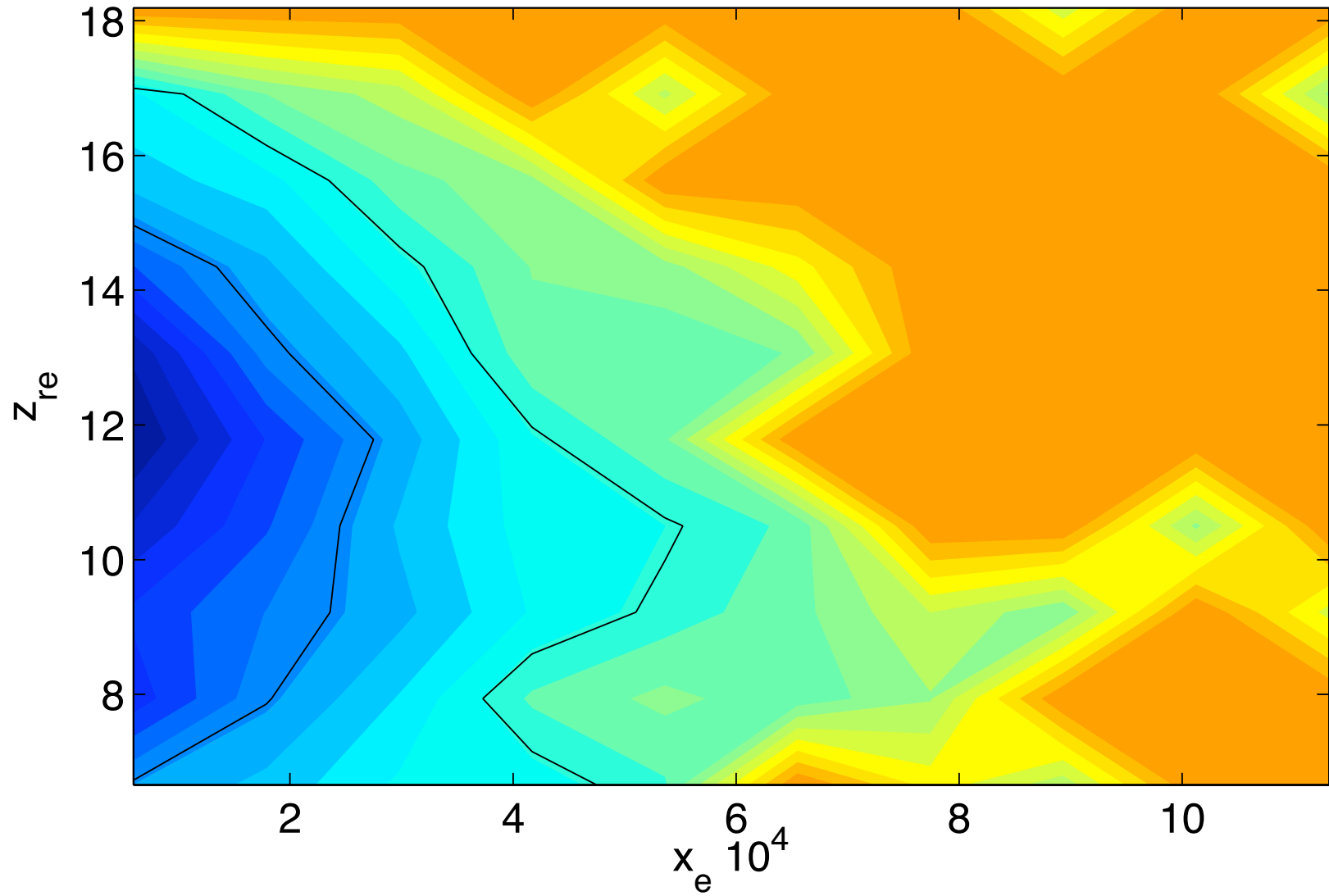
RECFAST: Seager, Sasselov & Scott 99;

COSMOMC: Lewis & Bridle 02

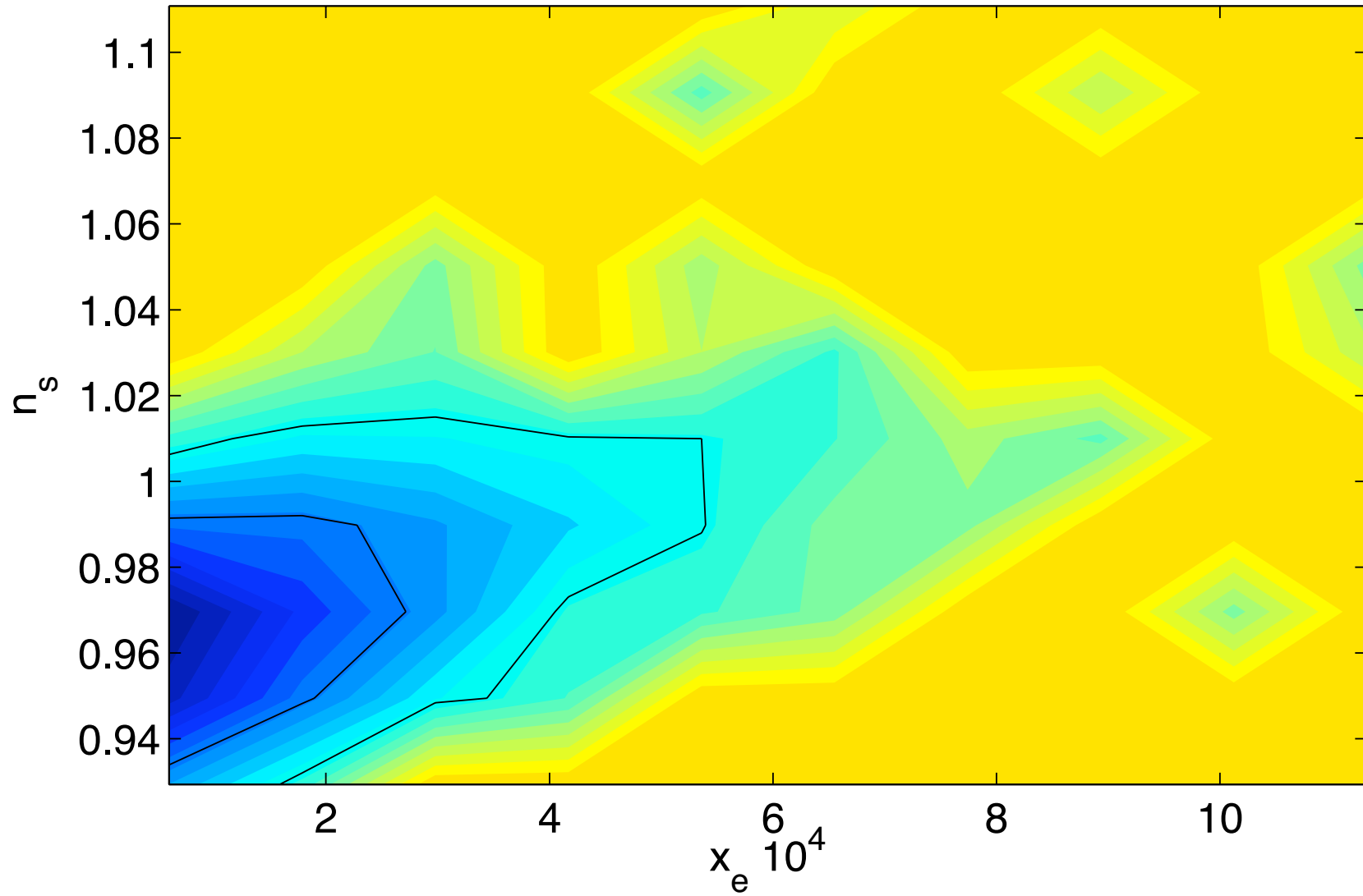
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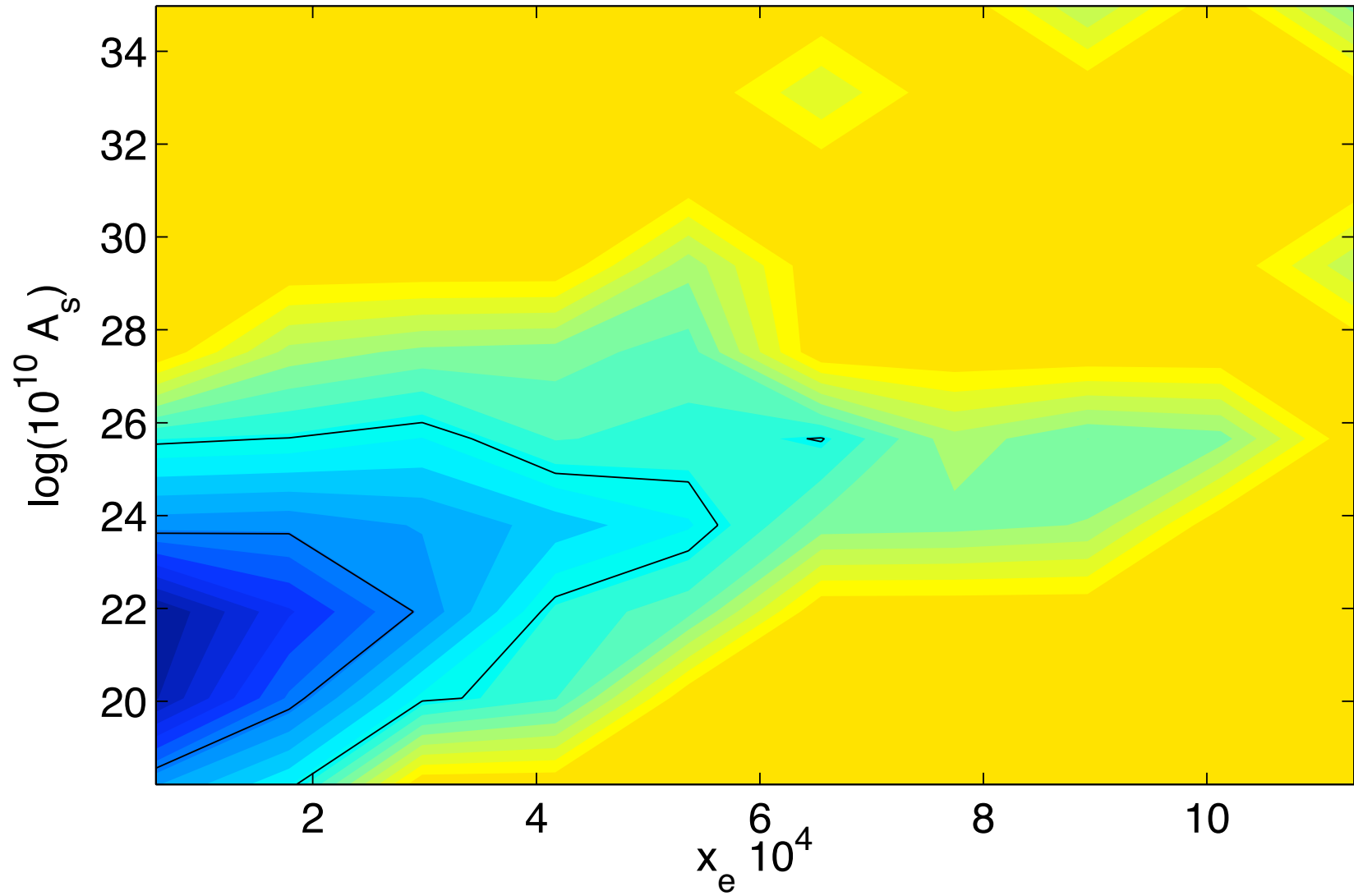
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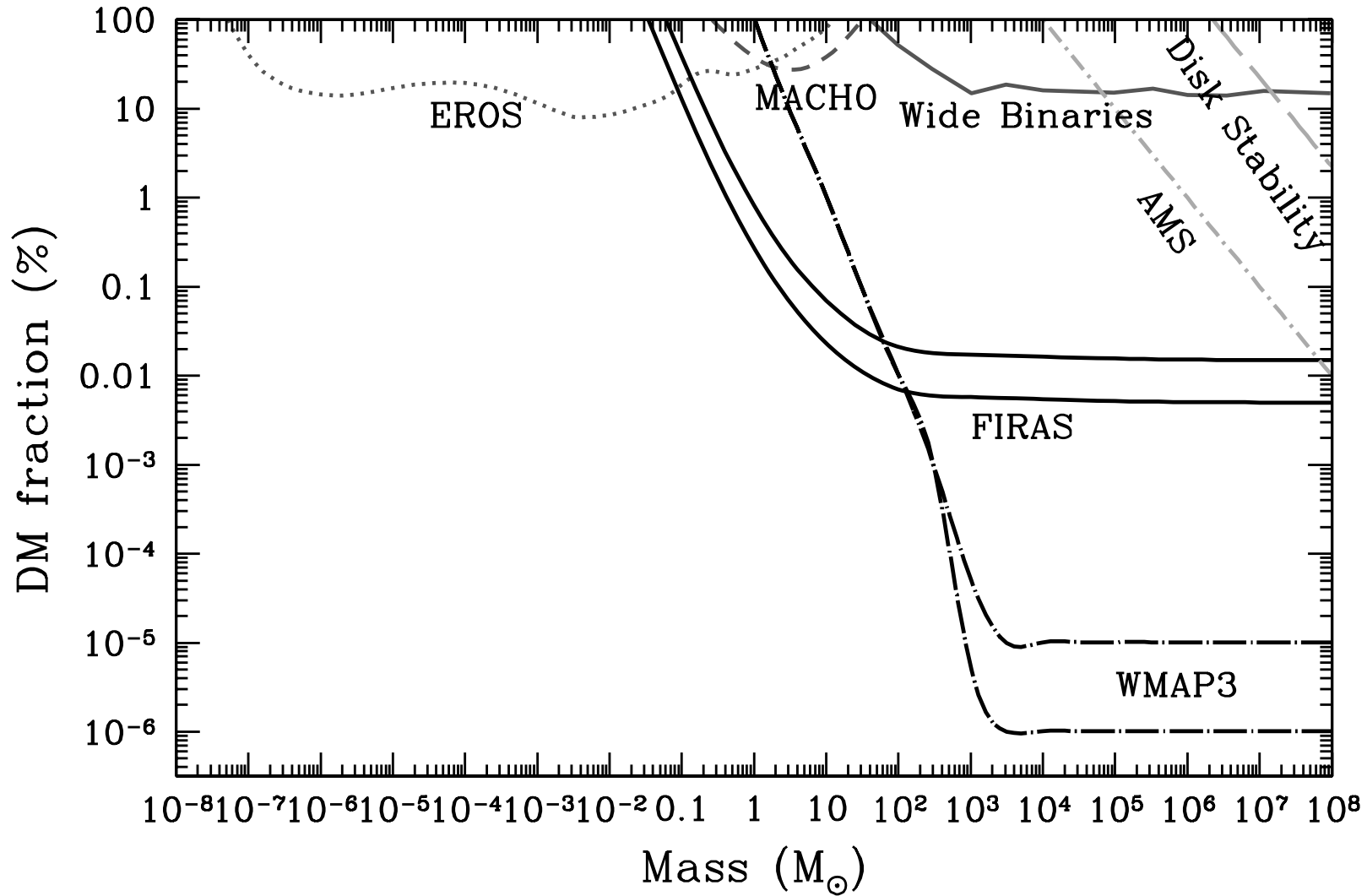
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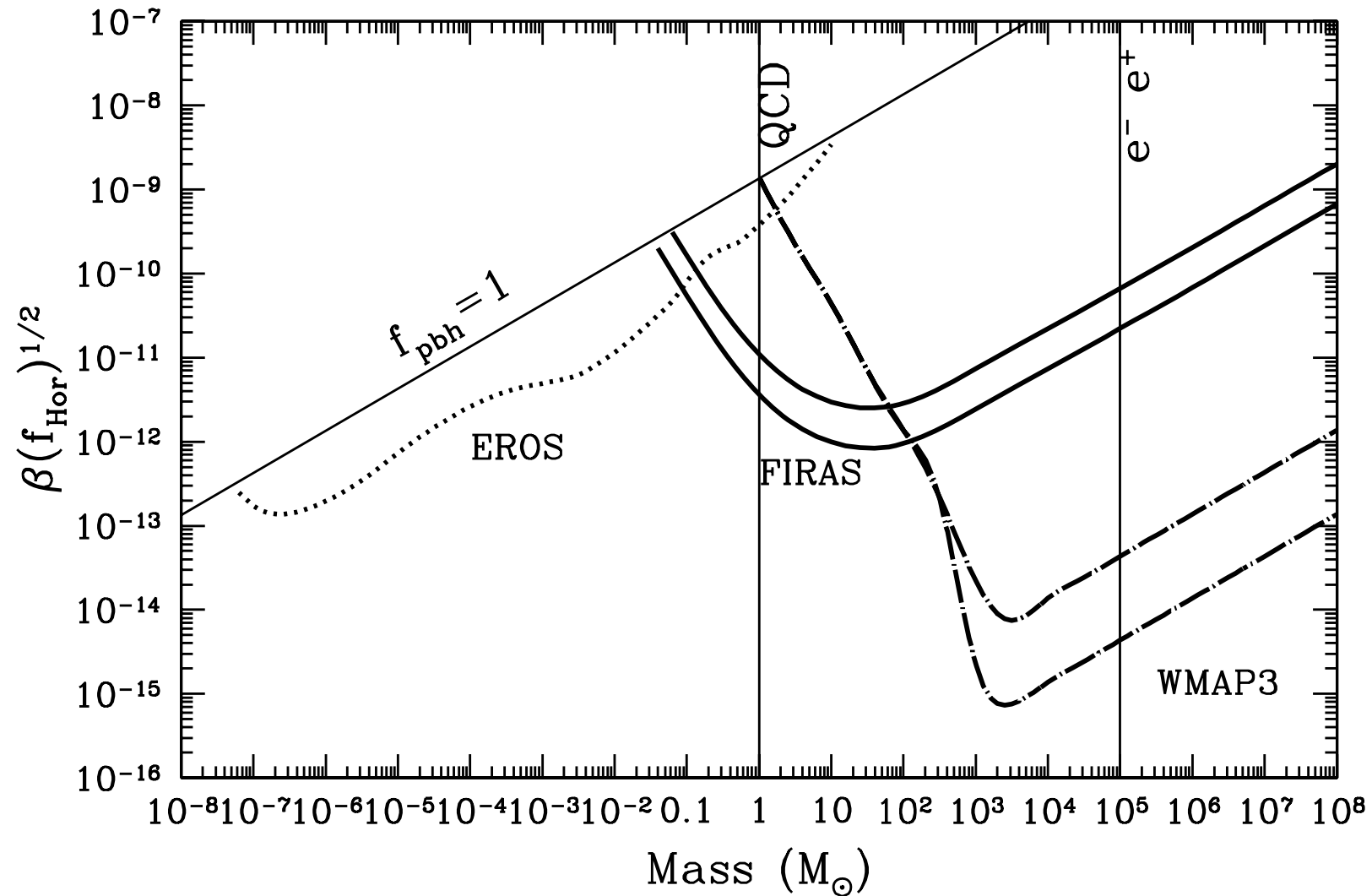
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