Einstein-Cartan Theory and Averaging

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- Riemann-Cartan Spacetime and Teleparallel Gravity
- 2 Averaged Connection
- Poincaré Gauge Theory
- 4 Conclusions and Outlook

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 \Rightarrow Such a manifold is called a **Riemann-Cartan spacetime** U_4 .

$$\label{eq:Kappa} {\mathcal{K}^{\lambda}}_{\mu\nu} = \mathring{\Gamma}^{\lambda}{}_{\mu\nu} - {\Gamma^{\lambda}}_{\mu\nu}.$$

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$$\Rightarrow$$

$$(L_4,g) \stackrel{Q=0}{\longrightarrow} U_4 \stackrel{T=0}{\longrightarrow} V_4 \stackrel{R=0}{\longrightarrow} R_4$$

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General Relativity

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TELEPARALLEL GRAVITY

In terms of Weitzenböck connection

$$\mathcal{L} = \frac{e}{16\pi G} T$$

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$$\Rightarrow \frac{du_{\lambda}}{ds} - \Gamma_{\mu\lambda\nu}u^{\mu}u^{\nu} = K_{\lambda\mu\nu}u^{\mu}u^{\nu}$$

$$\langle T^{\mu}{}_{\nu} \rangle(x) = \frac{1}{V_{\Sigma}} \int_{\Sigma} P^{\mu}{}_{\mu'}(x, x') P_{\nu}{}^{\nu'}(x, x') T^{\mu'}{}_{\nu'}(x') e(x') d^4x'$$

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- where $e(x') = \det(e^{i}_{\mu}(x')) = \sqrt{-g(x')}$
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Averaging Process and Averaged Connection

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$$\bar{\mathcal{T}}^{\lambda}{}_{\mu\nu} = \frac{1}{2}(\bar{\mathsf{\Gamma}}^{\lambda}{}_{\nu\mu} - \bar{\mathsf{\Gamma}}^{\lambda}{}_{\mu\nu}) = \frac{1}{2}(\left\langle \mathcal{T}^{\lambda}{}_{\mu\nu} \right\rangle - \mathcal{T}^{\lambda}{}_{\mu\nu})$$

⇒ The averaged geometry is a Riemann-Cartan geometry

Hehl, von der Heyde, Kerlick, and Nester:

• Lagrangian invariant under Poincaré gauge transformations

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• Metric energy-momentum tensor $\sigma^{\mu\nu}=2\delta\mathcal{L}/\delta g_{\mu\nu}$ and spin energy potential $\mu_{\lambda}^{\nu\mu} = \delta \mathcal{L}/\delta T^{\lambda}_{\mu\nu}$

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Lagrangian invariant under Poincaré gauge transformations

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- Metric energy-momentum tensor $\sigma^{\mu\nu}=2\delta\mathcal{L}/\delta g_{\mu\nu}$ and spin energy potential $\mu_{\lambda}^{\nu\mu} = \delta \mathcal{L}/\delta T^{\lambda}_{\mu\nu}$
- Spin angular momentum tensor $\tau_{\lambda}^{\nu\mu} = \delta \mathcal{L}/\delta K^{\lambda}_{\mu\nu}$ and total energy-momentum tensor $\sum_{i}^{\mu} = \delta \mathcal{L}/\delta e^{i}_{\mu}$

Hehl, von der Hevde, Kerlick, and Nester:

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- They are related according to

$$\mu^{\lambda\mu\nu} = -\tau^{\lambda\mu\nu} + \tau^{\mu\nu\lambda} - \tau^{\nu\lambda\mu}$$

and

$$\Sigma^{\mu\nu} = \sigma^{\mu\nu} - (\nabla_{\lambda} + 2T^{\rho}_{\lambda\rho})\mu^{\mu\nu\lambda}$$

Matter equation

$$\delta \mathcal{L}/\delta \Psi = 0$$

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• 1st field equation

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2nd field equation

$$T^{\lambda}{}_{\mu\nu} + \delta^{\lambda}_{\mu} T^{\rho}{}_{\nu\rho} + \delta^{\lambda}_{\nu} T^{\rho}{}_{\mu\rho} = 8\pi G \tau_{\mu\nu}{}^{\lambda}$$

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⇒ Combined field equation

$$\mathring{R}^{\mu\nu} - \frac{1}{2}\mathring{R}g^{\mu\nu}
= 8\pi G \sigma^{\mu\nu} + 8\pi G \left(-2\tau^{\mu\rho}{}_{\lambda}\tau^{\nu\lambda}{}_{\rho} + 2\tau^{\mu\rho}{}_{\rho}\tau^{\nu\lambda}{}_{\lambda} - 2\tau^{\mu\rho\lambda}\tau^{\nu}{}_{\rho\lambda} \right)
+ \tau^{\rho\lambda\mu}\tau_{\rho\lambda}{}^{\nu} + \frac{1}{2}g^{\mu\nu} \left(2\tau_{\sigma}{}^{\rho}{}_{\lambda}\tau^{\sigma\lambda}{}_{\rho} - 2\tau_{\sigma}{}^{\rho}{}_{\rho}\tau^{\sigma\lambda}{}_{\lambda} + \tau^{\sigma\rho\lambda}\tau_{\sigma\rho\lambda} \right)$$

• Nonminimal coupling of the Riemannian connection to the matter field Ψ

$$\bar{\Gamma}^{\lambda}{}_{\mu\nu} = \mathring{\Gamma}^{\lambda}{}_{\mu\nu} + 8\pi G \left(\tau_{\mu\nu}{}^{\lambda} - \tau_{\nu}{}^{\lambda}{}_{\mu} + \tau^{\lambda}{}_{\mu\nu} + \delta^{\lambda}_{\mu}\tau_{\nu\sigma}{}^{\sigma} - g_{\mu\nu}\tau^{\lambda\sigma}{}_{\sigma} \right)$$

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- ⇒ Apply to cosmological perturbation theory

Cosmological Perturbation Theory

Linear scalar perturbations about a FLRW background

$$ds^2 = -(1+2\phi)dt^2 + a^2(t)(1-2\phi)\delta_{ij}dx^idx^j$$

• Perturbed energy momentum tensor $T_{\mu\nu} = \bar{\rho}(1+\delta)u_{\mu}u_{\nu}$

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Cosmological Perturbation Theory

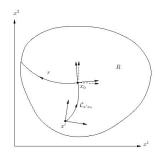
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- G_0^0 component: energy constraint equation $\nabla^2 \phi(t_0) = 4\pi G a^2 \bar{\rho}(t_0) \delta(t_0)$
- G_i^0 component: momentum constraint equation $\dot{a}(\partial_i \phi) + a(\partial_i \dot{\phi}) = -4\pi G a \bar{\rho} v_i$
- Trace of the G_i^i component: evolution equation

$$\ddot{\phi} + 4H\dot{\phi} + 3H^2\phi + 2\frac{\ddot{a}}{a}\phi - 2H^2\phi = 0$$

How to Choose the Tetrad Field



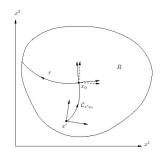
Parallel transport along geodesics $C_{x_0x'}$ realized by Wegner-Wilson line operator

$$V(x',x_0;\mathcal{C}_{x_0x'})=\mathcal{P}\exp\left[-\int_{\mathcal{C}_{x_0x'}}dz^{\mu}\;\mathring{\Gamma}_{\mu}(z)
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Poincaré Gauge Theory

where $\check{\Gamma}_{\mu}(x)$ are four matrices with components $(\mathring{\Gamma}_{\mu}(x))^{\lambda}_{\nu} = \mathring{\Gamma}^{\lambda}_{\mu\nu}(x)$

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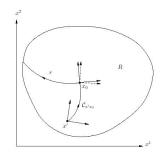
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Perturbed geodesic

$$z_P^i(s) = \frac{dz^i}{d au}(0)\frac{s}{a} + \frac{dz^i}{d au}(0)\int_0^{s/a}\phi(s')ds' + \delta^{ij}\epsilon_{jkl}v^k(\frac{s}{a})\frac{dz^l}{d au}(0)$$

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Perturbed connector

$$V_j{}^i(\tau,0,;\mathcal{C}_{0\tau}) = (1+\phi(\tau)-\phi(0)))\,\delta_j{}^i - \delta^{ik}\epsilon_{kjl}\frac{dv^l}{d\tau}(\tau)$$

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$$\delta(t_0) = A \exp(-\frac{r^2}{(2\sigma)^2})$$

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⇒ The velocity field becomes

$$v^{i} = \frac{\pi^{3/2}AG(2\sigma)^{3}\bar{\rho}(t_{0})}{6\pi Gt^{2/3}t_{0}^{1/3}} \frac{1}{r} \operatorname{erf}\left(\frac{r}{2\sigma}\right) - \frac{\pi^{3/2}AG(2\sigma)^{3}\bar{\rho}(t_{0})}{4\pi G} \left(\frac{4}{15}\frac{t^{2/3}}{t_{0}^{1/3}} + \frac{2}{5}\frac{t}{t_{0}^{2}}\right) \left(-\frac{x^{i}}{r^{3}} \operatorname{erf}\left(\frac{r}{2\sigma}\right) + \frac{1}{r}\frac{2}{\sqrt{\pi}} \exp\left(-\left(\frac{r}{2\sigma}\right)^{2}\right) \frac{x^{i}}{2\sigma r}\right)$$

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- In a cosmological setting, what is the magnitude of the additional term in the field equation and does it behave like dark energy?