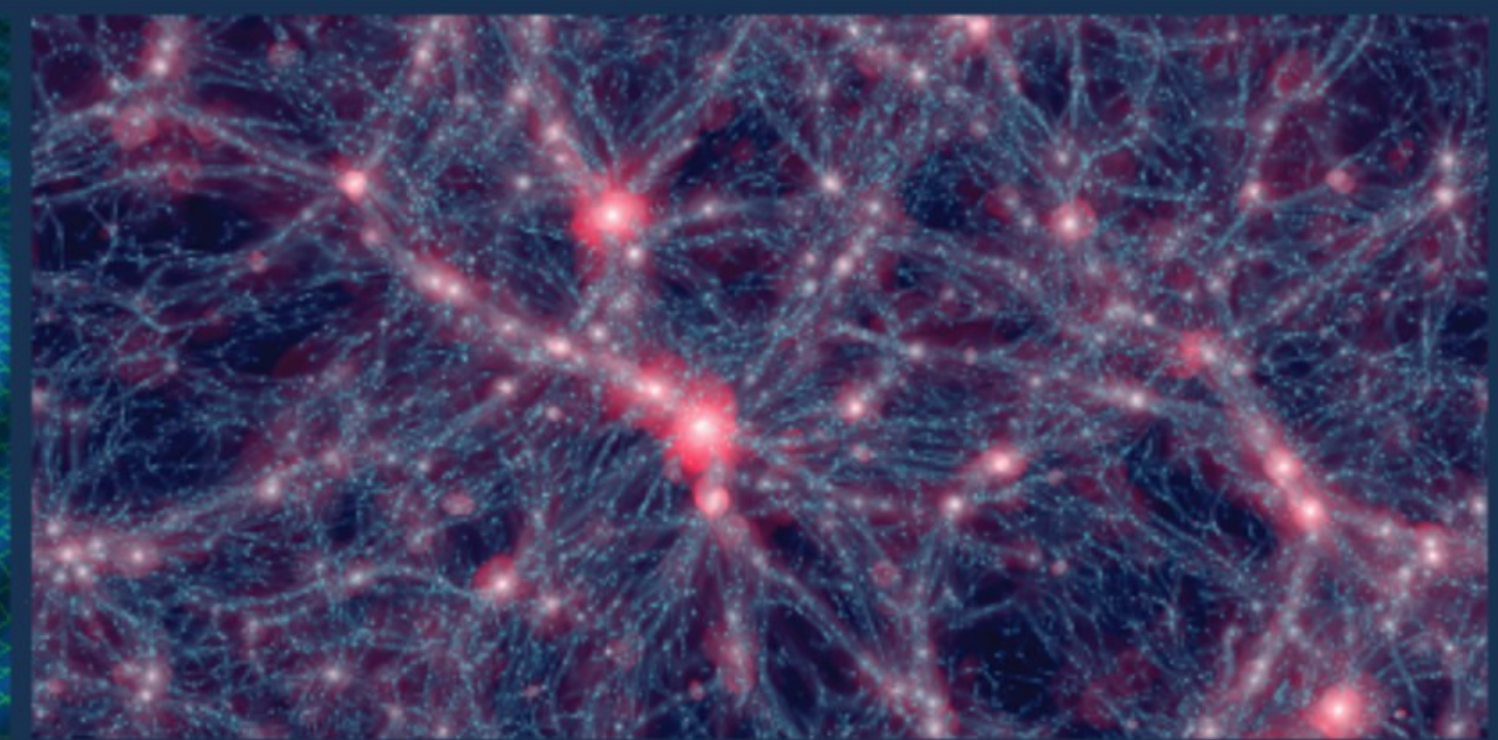
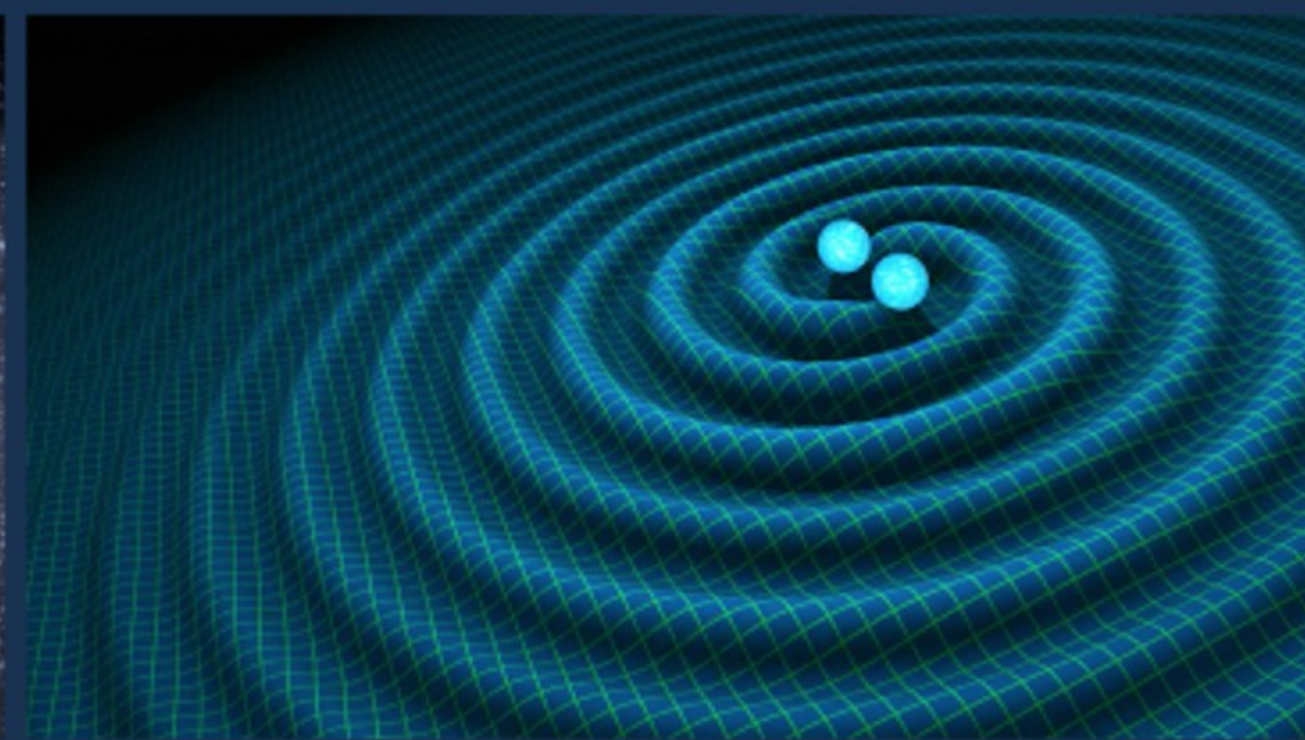
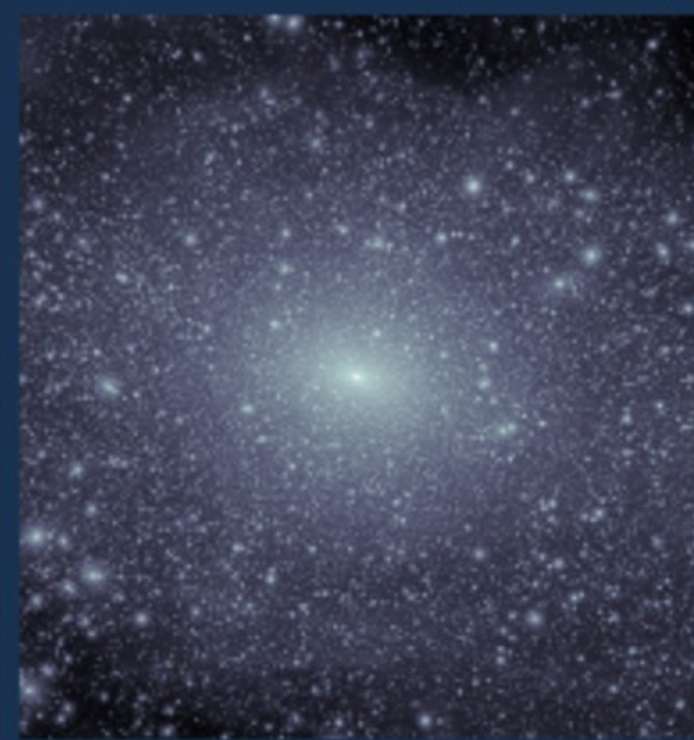




Home	Venue	Accomodation	Social events	Contact
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DARK MATTERS 11-13 December 2017
Joe Silk's 75th Birthday



Participants

Monday 11th December

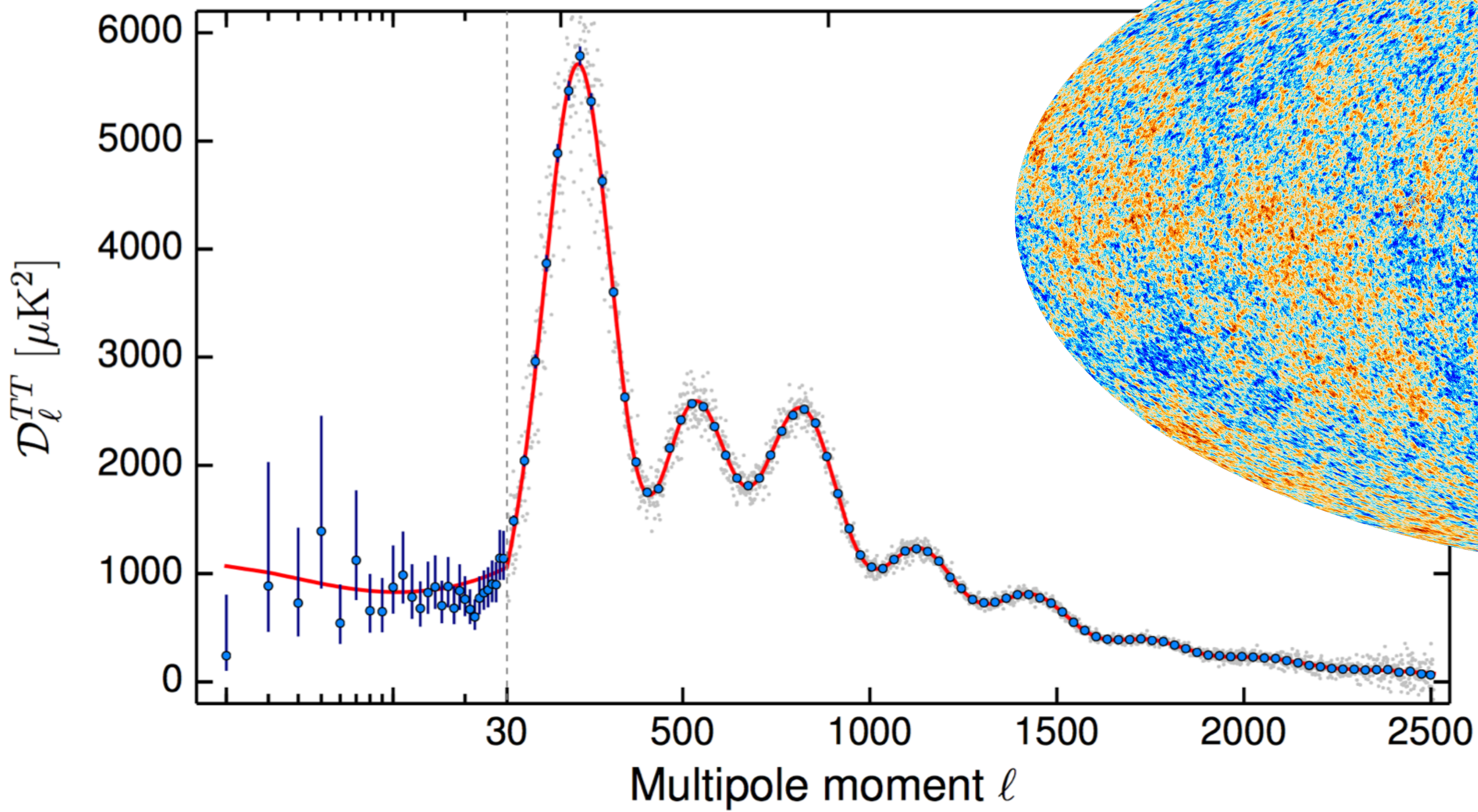
Tuesday 12th December

Wednesday 13th December

*Evidences for inflation
constraints on alternatives*

THANKS TO
JÉRÔME MARTIN
FOR HIS HELP

Planck 2015



$\Omega_{\kappa} = 0.000 \pm 0.005$

$n_s = 0.9639 \pm 0.0047$ almost scale invariant

$f_{NL}^{loc} = 0.8 \pm 5$
 $f_{NL}^{eq} = -4 \pm 43$
 $f_{NL}^{ort} = -26 \pm 21$

excluded
 gaussian signal

isocurvature $\lesssim 1\%$

$r < 0.11$

quantum vacuum fluctuations of a single scalar d.o.f



compatible with ***INFLATION***

Inflation: phase of accelerated expansion

... needs a fluid with negative pressure

$$\frac{\ddot{a}}{a} = -\frac{1}{6M_{\text{P}}^2} (\rho + 3p) > 0$$

Inflaton (compatible with the cosmological principle)

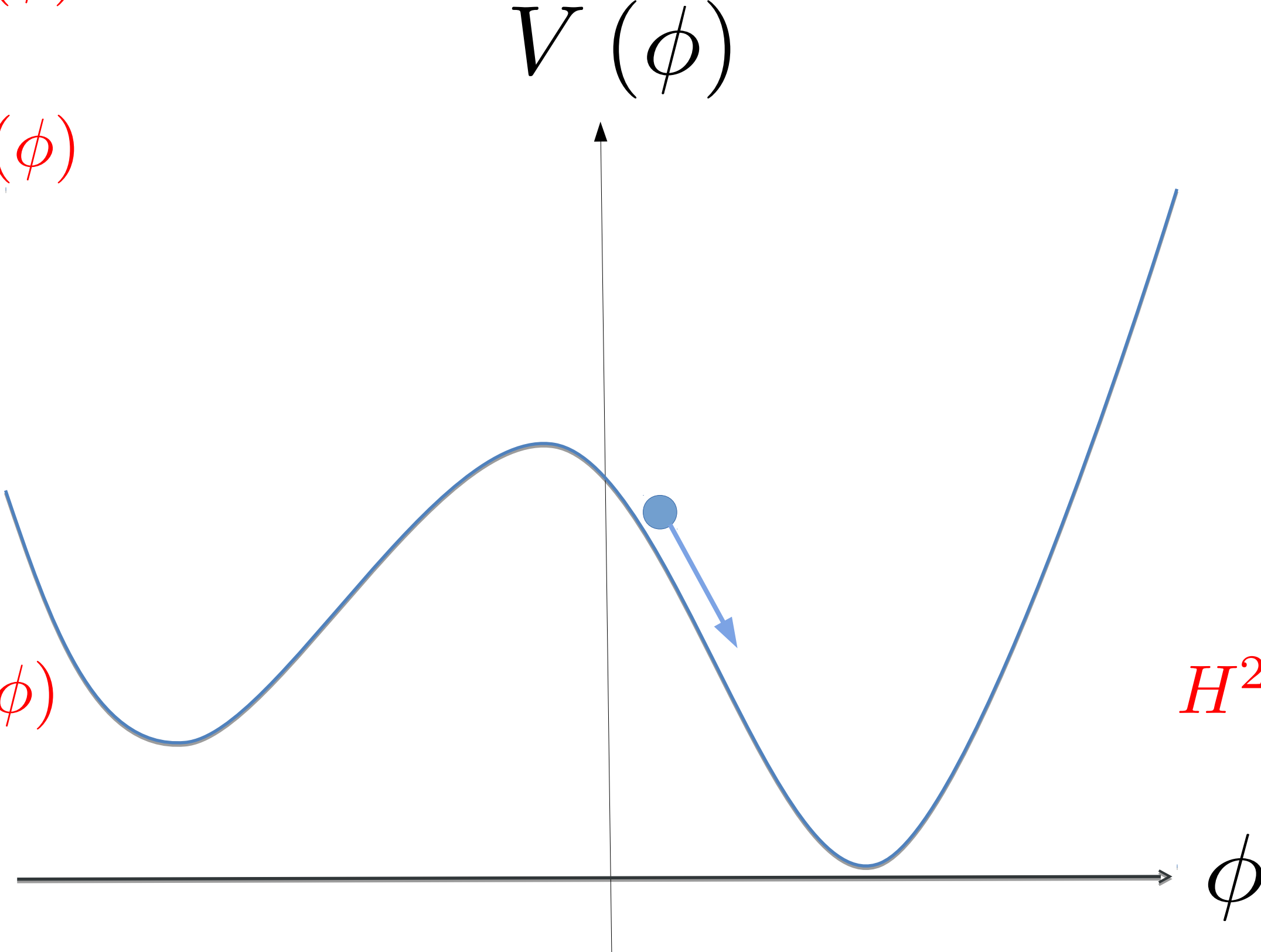
$$\rho = \frac{1}{2}\dot{\phi}^2 + V(\phi)$$

$$p = \frac{1}{2}\dot{\phi}^2 - V(\phi)$$

Slow-roll

$$\frac{1}{2}\dot{\phi}^2 \ll V(\phi)$$

$$\rho + 3p \simeq -2V(\phi)$$



$$\frac{1}{2}\dot{\phi}^2 \ll V(\phi)$$

$$p \simeq -\rho$$

$$\frac{d\rho}{dt} = -3H(\rho + p) \simeq 0$$

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{\rho}{3M_{\text{P}}^2} \simeq H_{\text{inf}}^2$$

$$a(t) \simeq e^{H_{\text{inf}}t}$$

Standard questions and inflationary answers

Singularity

Not solved... actually not addressed!

Horizon $d_H \equiv a(t) \int_{t_{\text{ini}}}^t \frac{d\tau}{a(\tau)}$ can be made as big as one wishes

Flatness $\frac{d}{dt} |\Omega - 1| = -2 \frac{\ddot{a}}{\dot{a}^3} \longrightarrow \ddot{a} > 0 \quad \& \quad \dot{a} > 0$
accelerated expansion (**inflation**)

Homogeneity & Isotropy

Initial Universe = very small patch

Accelerated expansion drives the shear to zero...

\implies vacuum state!

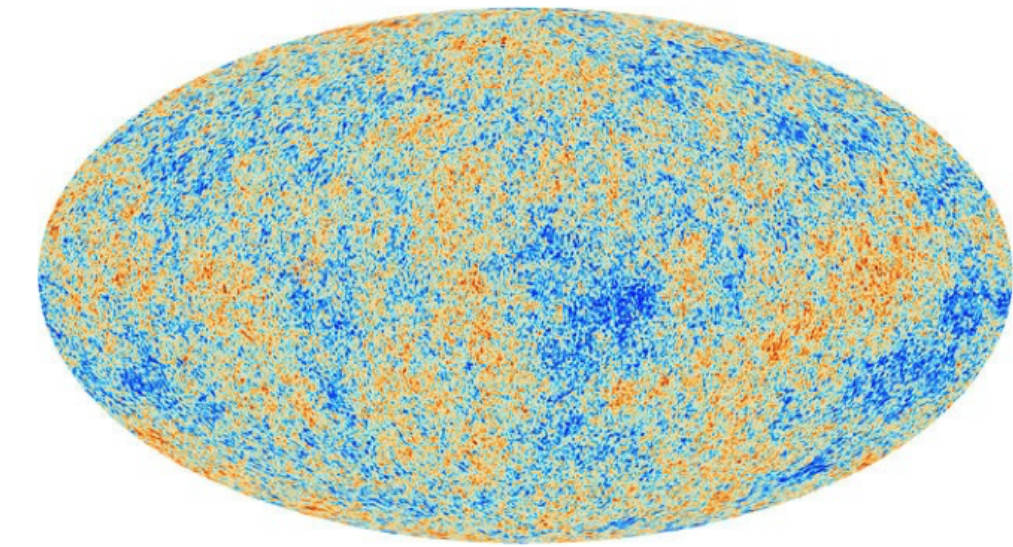
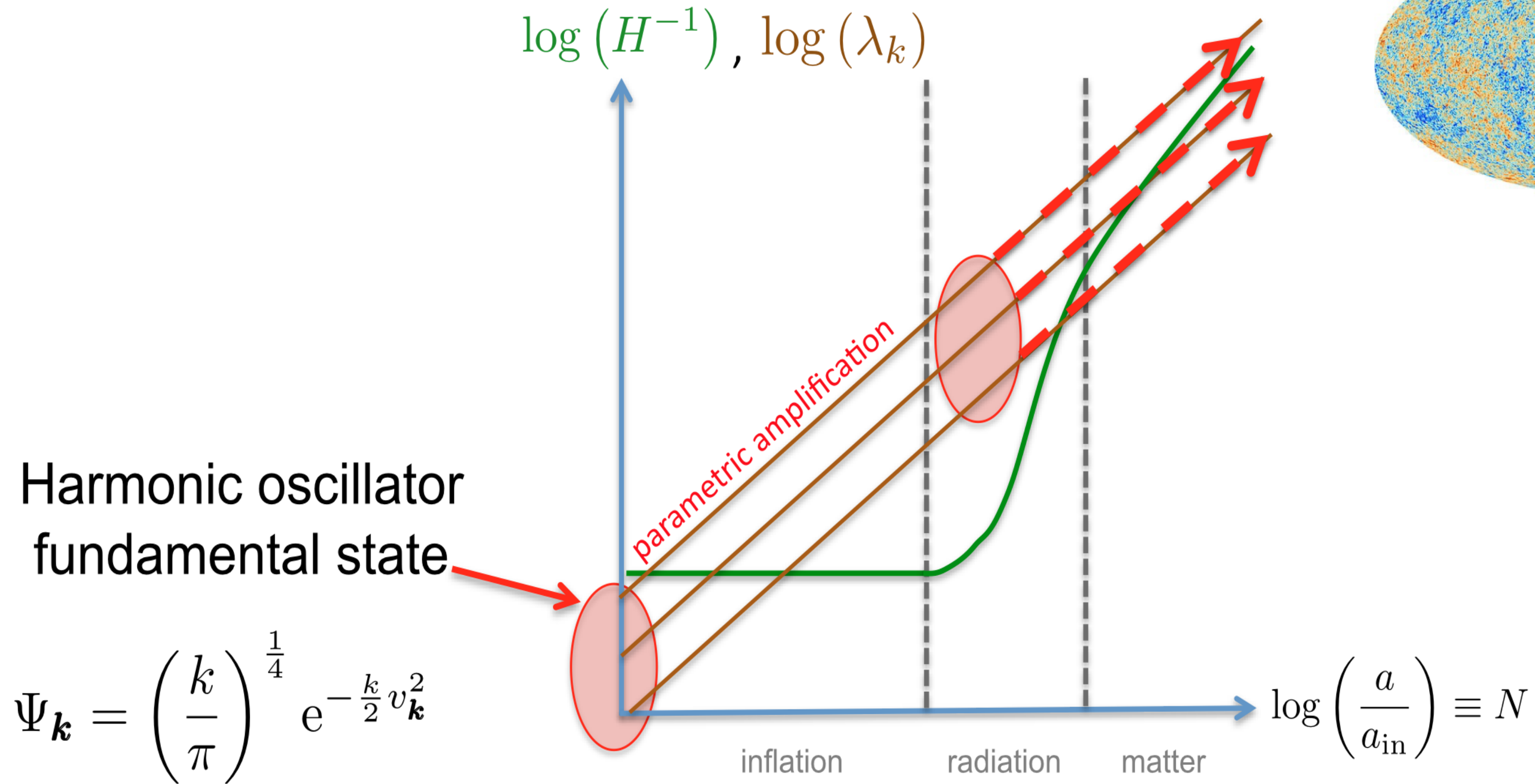
Perturbations

Bonus of the theory: predictions!!!

Others

dark matter/energy, baryogenesis,...

Almost constant Hubble radius $R_H = H_{\text{inf}}^{-1}$



Everything depends on the slow roll parameters

$$\epsilon_1 \simeq \frac{1}{2M_{\text{P}}^2} \left(\frac{V_{,\phi}}{V} \right)^2$$

$$\epsilon_2 \simeq \frac{2}{M_{\text{P}}^2} \left[\left(\frac{V_{,\phi}}{V} \right)^2 - \frac{V_{,\phi\phi}}{V} \right]$$

Evaluated at ϕ_* (pivot scale leaves Hubble radius):

$$\left. \begin{aligned} \mathcal{P}_\zeta &\propto \frac{H_{\text{inf}}^2}{M_{\text{P}}^2 \epsilon_{1*}} [1 + \mathcal{O}(k, \epsilon_*)] \\ \mathcal{P}_h &\propto \frac{H_{\text{inf}}^2}{M_{\text{P}}^2} [1 + \mathcal{O}(k, \epsilon_*)] \end{aligned} \right\} \text{Spectral indices} \left\{ \begin{aligned} n_{\text{S}} &= 1 + \frac{d \ln \mathcal{P}_\zeta}{d \ln k} = -2\epsilon_{1*} - \epsilon_{2*} \\ n_{\text{T}} &= \frac{d \ln \mathcal{P}_h}{d \ln k} = -2\epsilon_{1*} \end{aligned} \right.$$

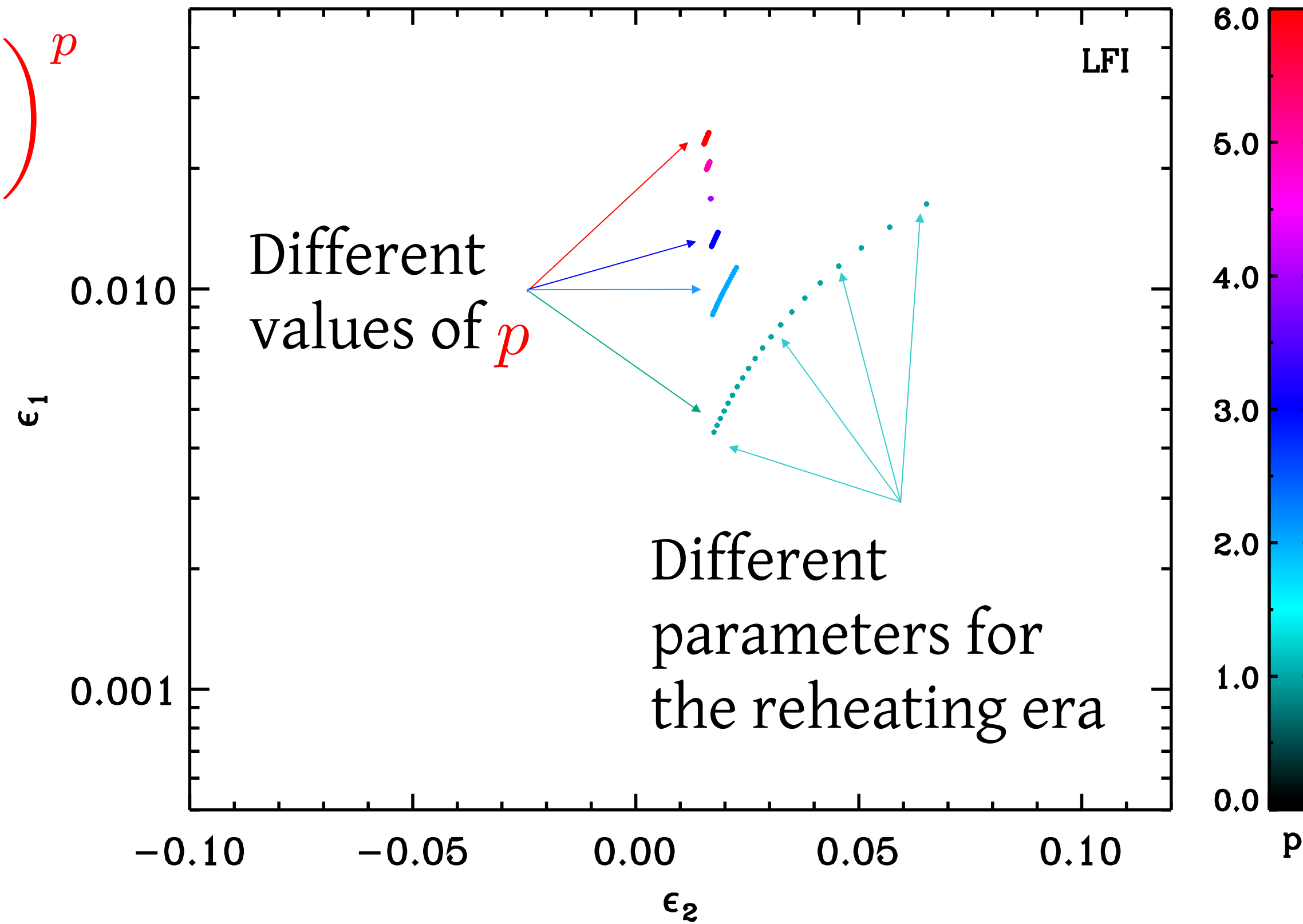
→ Consistency relation $r \equiv \frac{T}{S} = \frac{\mathcal{P}_\zeta}{\mathcal{P}_h} = 16\epsilon_{1*} = -8n_{\text{T}}$

+ running

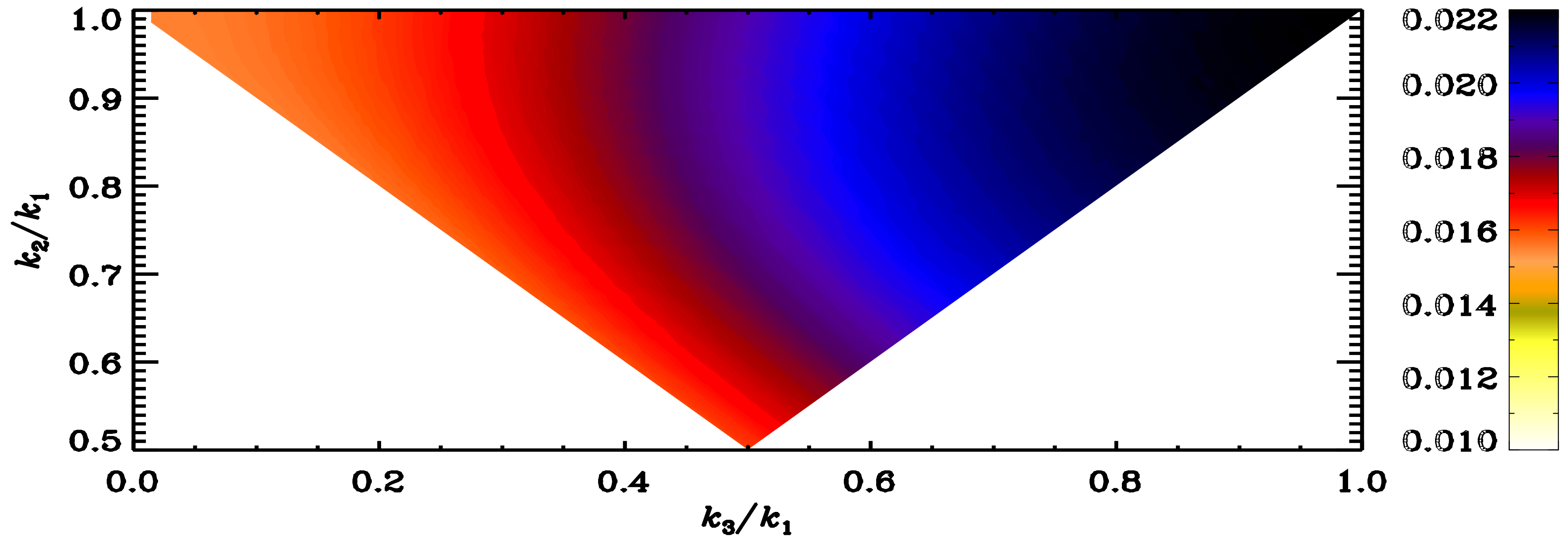
Various models, various predictions... (one example)

"Large field inflation"

$$V(\phi) = m^4 \left(\frac{\phi}{M_P} \right)^p$$



3 point correlation function...

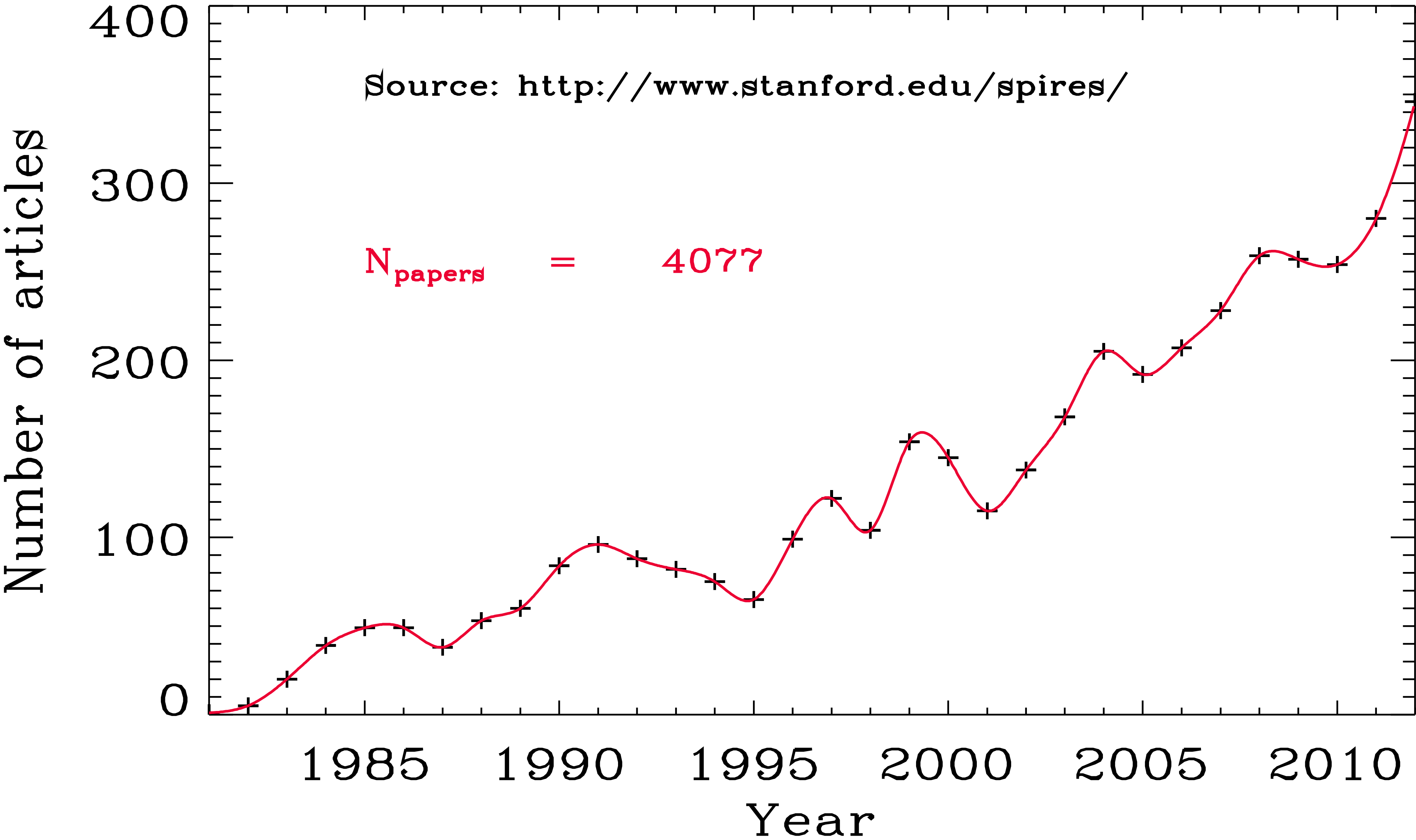


$$\langle \zeta_{\mathbf{k}_1} \zeta_{\mathbf{k}_2} \zeta_{\mathbf{k}_3} \rangle \propto f_{\text{NL}}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) \frac{\delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3)}{k_1^3 k_2^3 k_3^3} [k_1^3 \mathcal{P}_\zeta(k_2) \mathcal{P}_\zeta(k_3) + 2 \text{ perms.}]$$

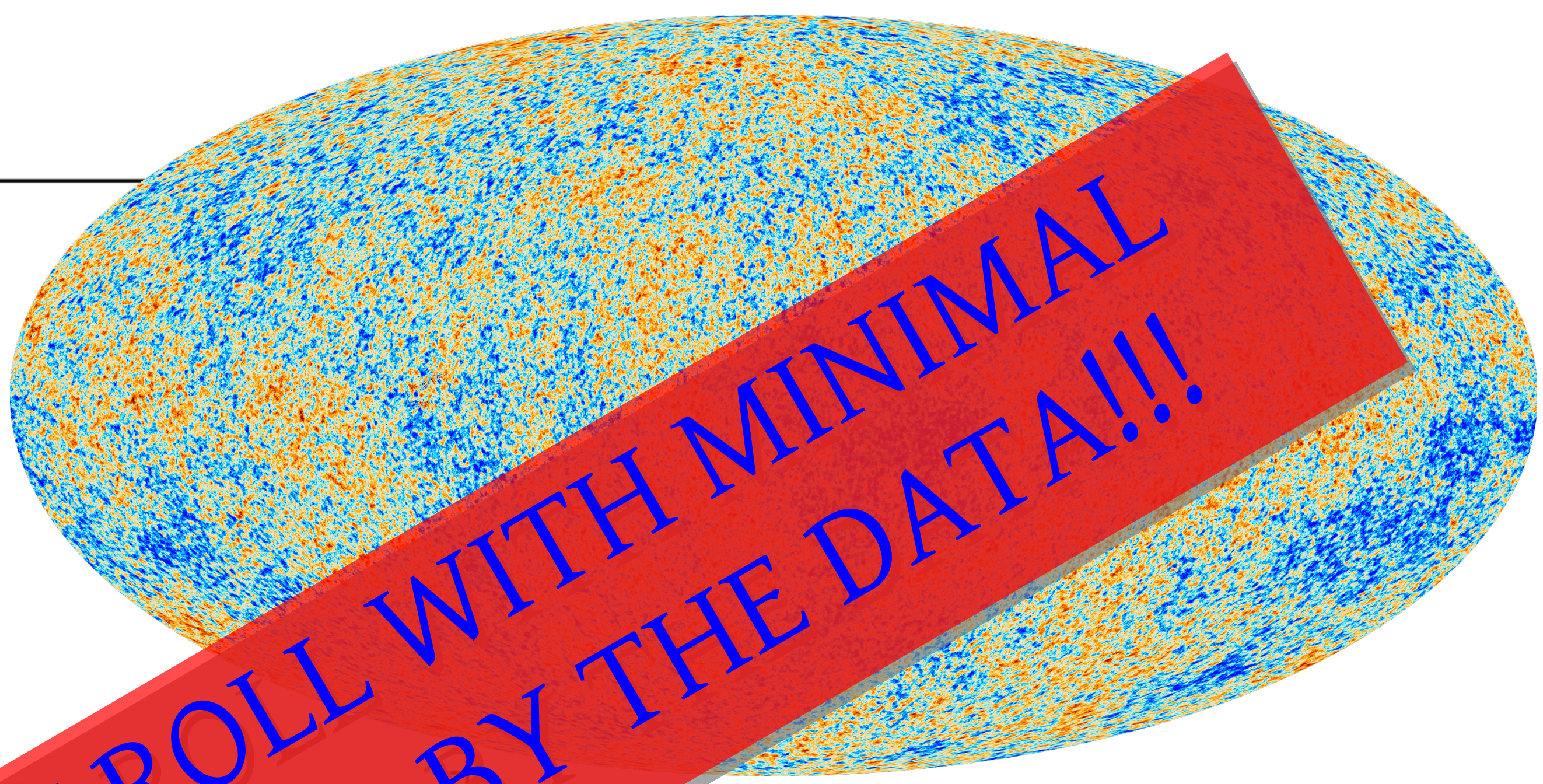
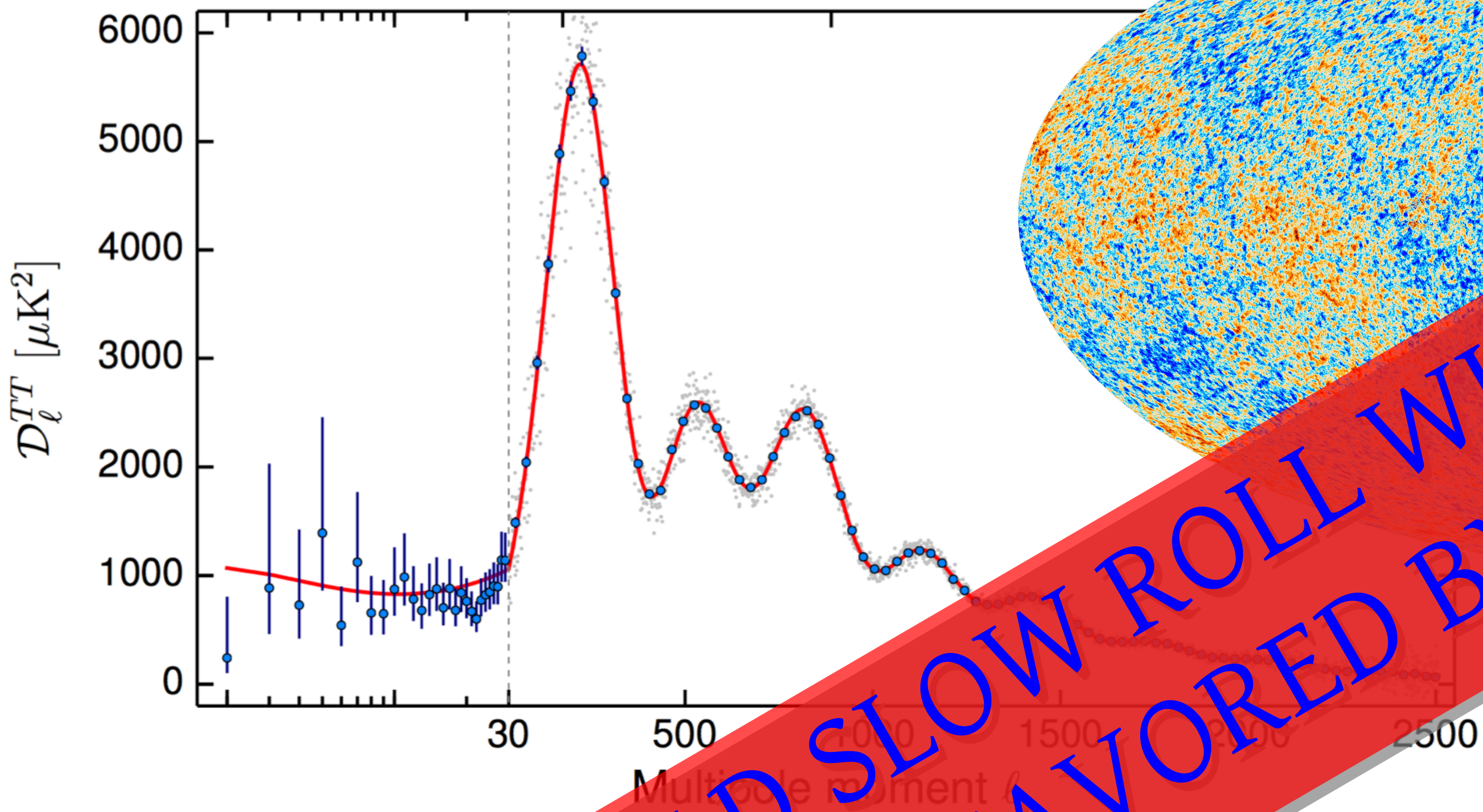


$\mathcal{O}(\epsilon_*)$ Unobservable at the moment

Comparing with data: pb = 100's of different models



Planck 2015



$$\Omega_{\kappa} = -0.040^{+0.038}_{-0.041}$$

$$n_s = 0.9639 \pm 0.0044$$

almost scale invariant

$$f_{loc}^{NL} = 0.8 \pm 1.5$$

excluded

$$f_{NL}^{eq} = -4 \pm 42$$

gaussian signal

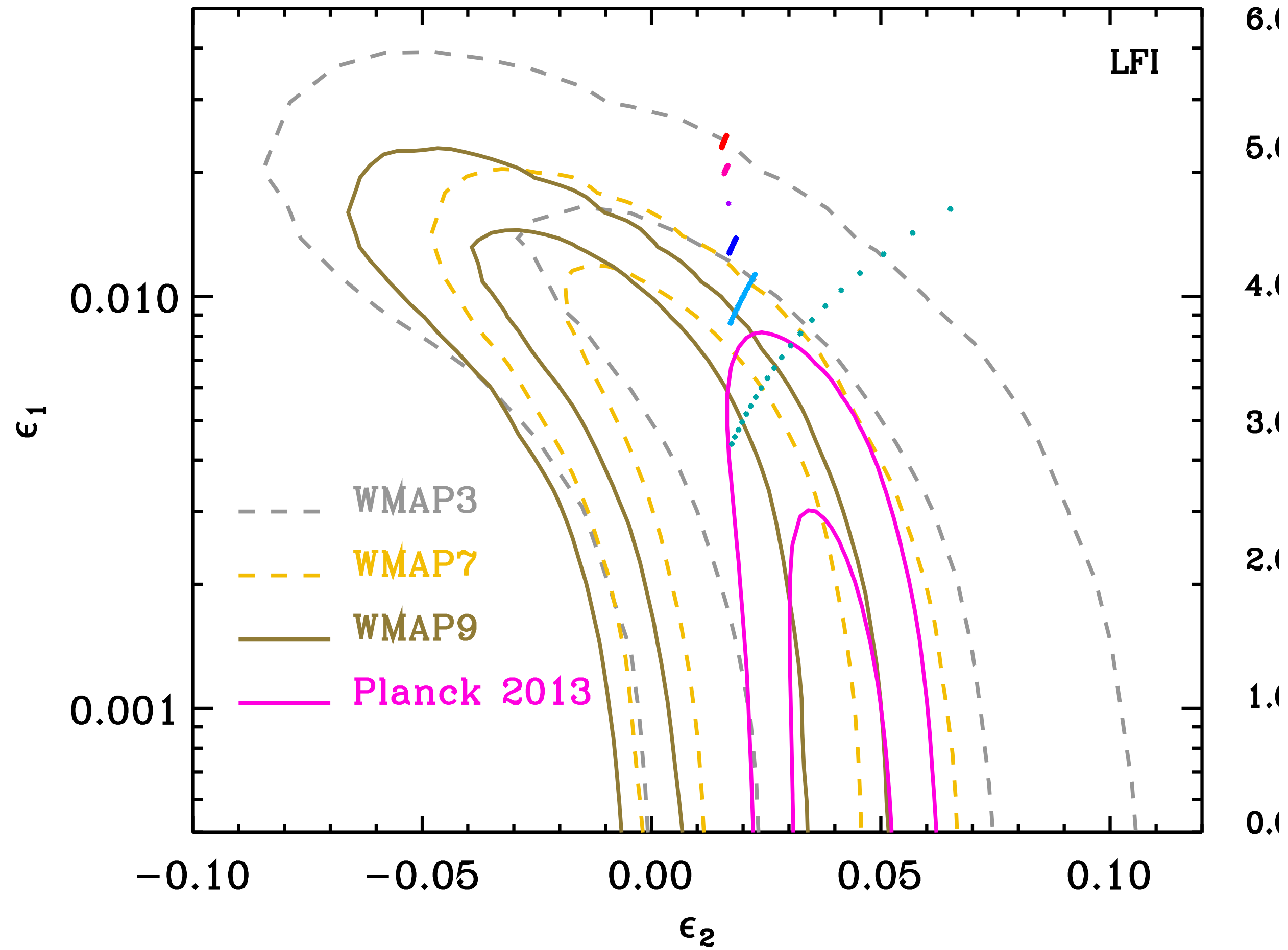
$$f_{NL}^{ort} = -20 \pm 21$$
$$r < 0.11$$

$$\text{isocurvature} \lesssim 1\%$$

quantum vacuum fluctuations of a single scalar d.o.f



compatible with *INFLATION*



CMB constraints on the slow-roll parameters

Planck 2015 constraints:

Slow-roll parameters

$$\epsilon_{1\star} < 0.0068$$

$$\epsilon_{2\star} = 0.029^{+0.008}_{-0.007}$$

Deviation from exact scale invariance

$$n_s = 0.9645 \pm 0.0049 < 1$$

Inflaton potential

$$M_{\text{P}} \frac{|V_{,\phi}|}{V} < 0.14$$

$$M_{\text{P}}^2 \frac{V_{,\phi\phi}}{V} = 0.01^{+0.005}_{-0.009}$$

No detection:

GW

$$r < 0.1$$

Running

$$\frac{dn_s}{d \ln k} = -0.0134 \pm 0.0009$$

Energy scale

$$H_{\text{inf}}^2 \sim \mathcal{P}_{\zeta} \epsilon_{1\star} \propto \rho_{\star}$$

$$\rho_{\star}^{\frac{1}{4}} < 2.2 \times 10^{16} \text{ GeV}$$

$$H_{\text{inf}} < 1.2 \times 10^{14} \text{ GeV}$$

Bayesian analysis \implies best models: J. Martin, C. Ringeval & V. Vennin, [arXiv:1303.3787](https://arxiv.org/abs/1303.3787)

Encyclopedia Inflationaris

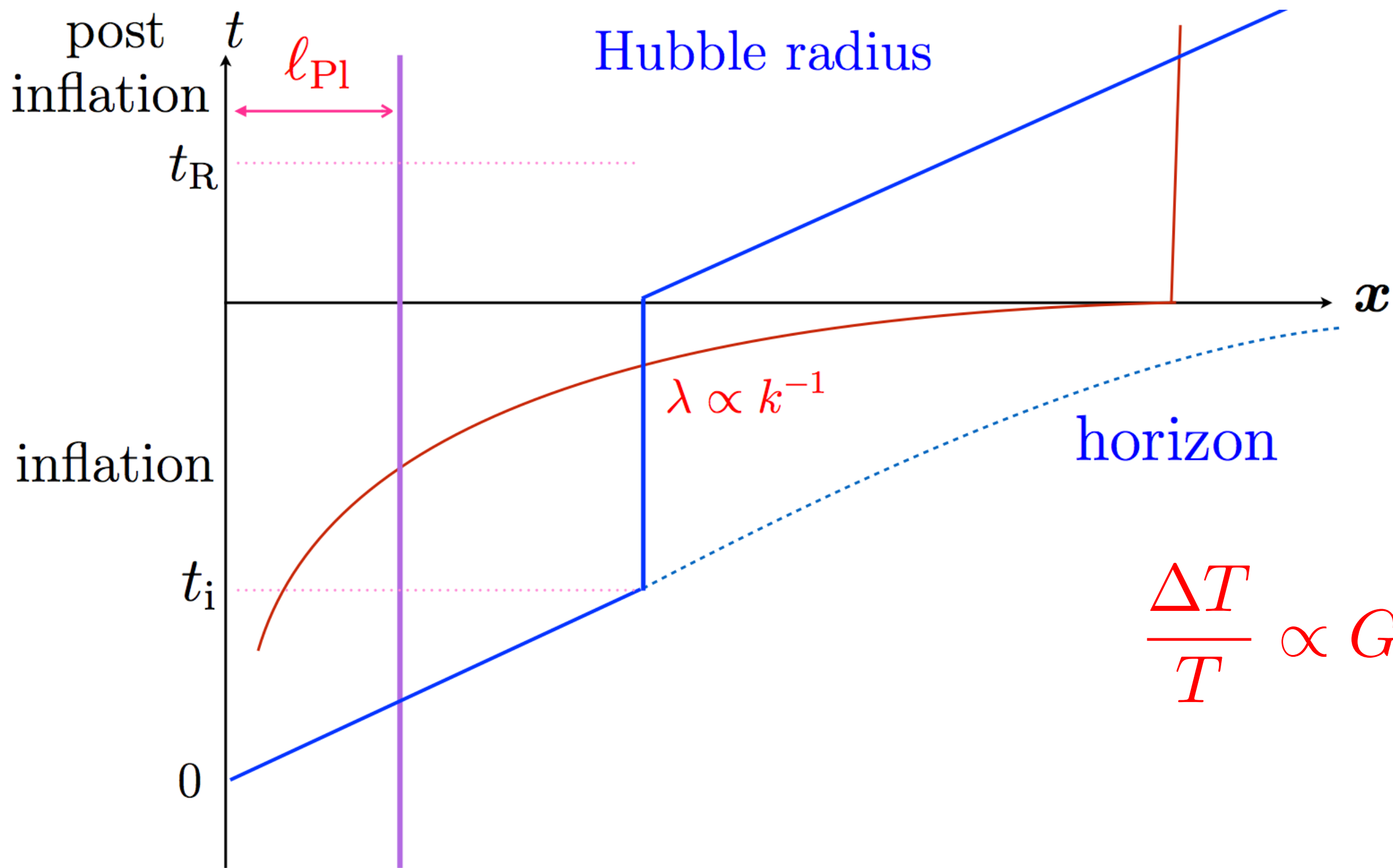
- Inflation:**
- ☺ solves cosmological puzzles
 - ☺ uses GR + scalar fields [(semi-)classical]
 - ☺ can be implemented in high energy theories
 - ☺ string implementation (brane inflation, ...)

 - ☺ makes falsifiable predictions...
 - ☺ ... consistent with all known observations

why bother with alternatives?

“When you have eliminated the impossible, whatever remains, however improbable, must be the truth.”

Sherlock Holmes (The sign of the four):



Singularity

$$\exists t_{\text{sing}} ; a(t_{\text{sing}}) = 0$$

Trans-Planckian

$$\exists t_{\text{sub}} ; \ell(t_{\text{sub}}) = \ell_0 \frac{a(t_{\text{sub}})}{a_0} \leq \ell_P$$

Hierarchy (amplitude)?

$$\frac{V(\phi)}{\Delta\phi^4} \leq 10^{-12}$$

Classical GR?

$$\frac{\Delta T}{T} \propto G_N E_{\text{inf}}^2 \sim \left(\frac{E_{\text{inf}}}{M_P} \right)^2 \rightarrow E_{\text{inf}} \sim 10^{-3} M_P$$

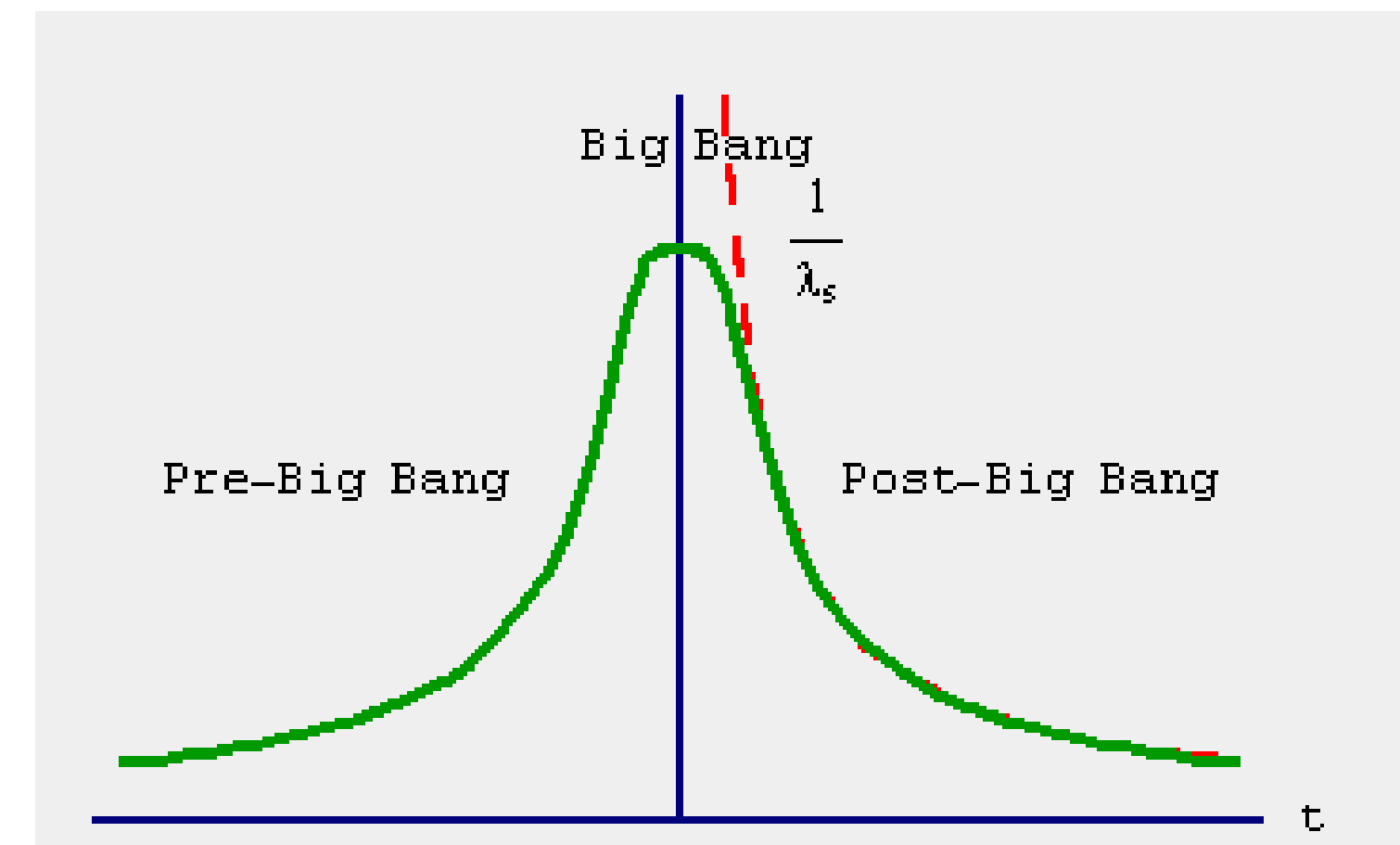
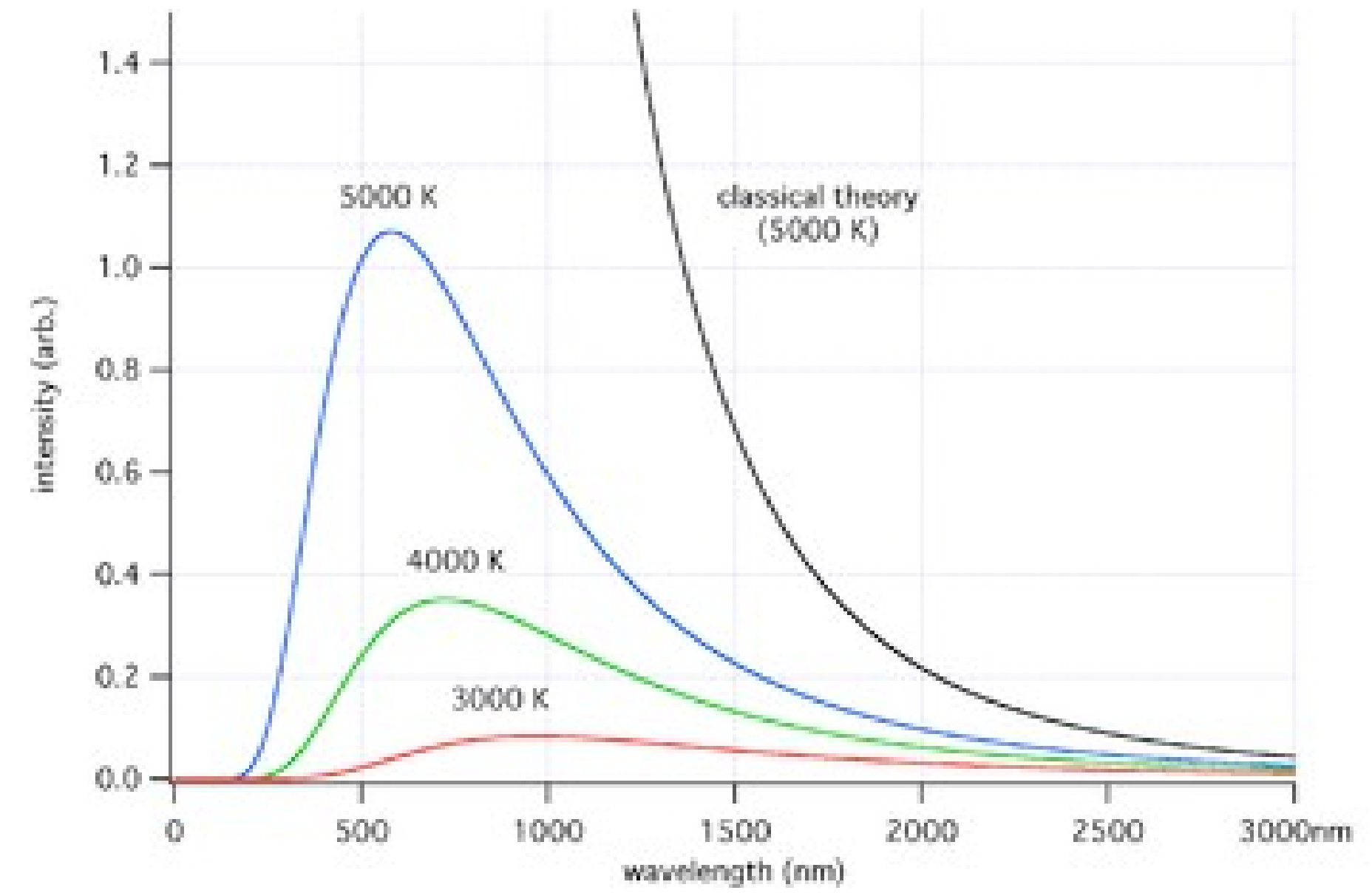
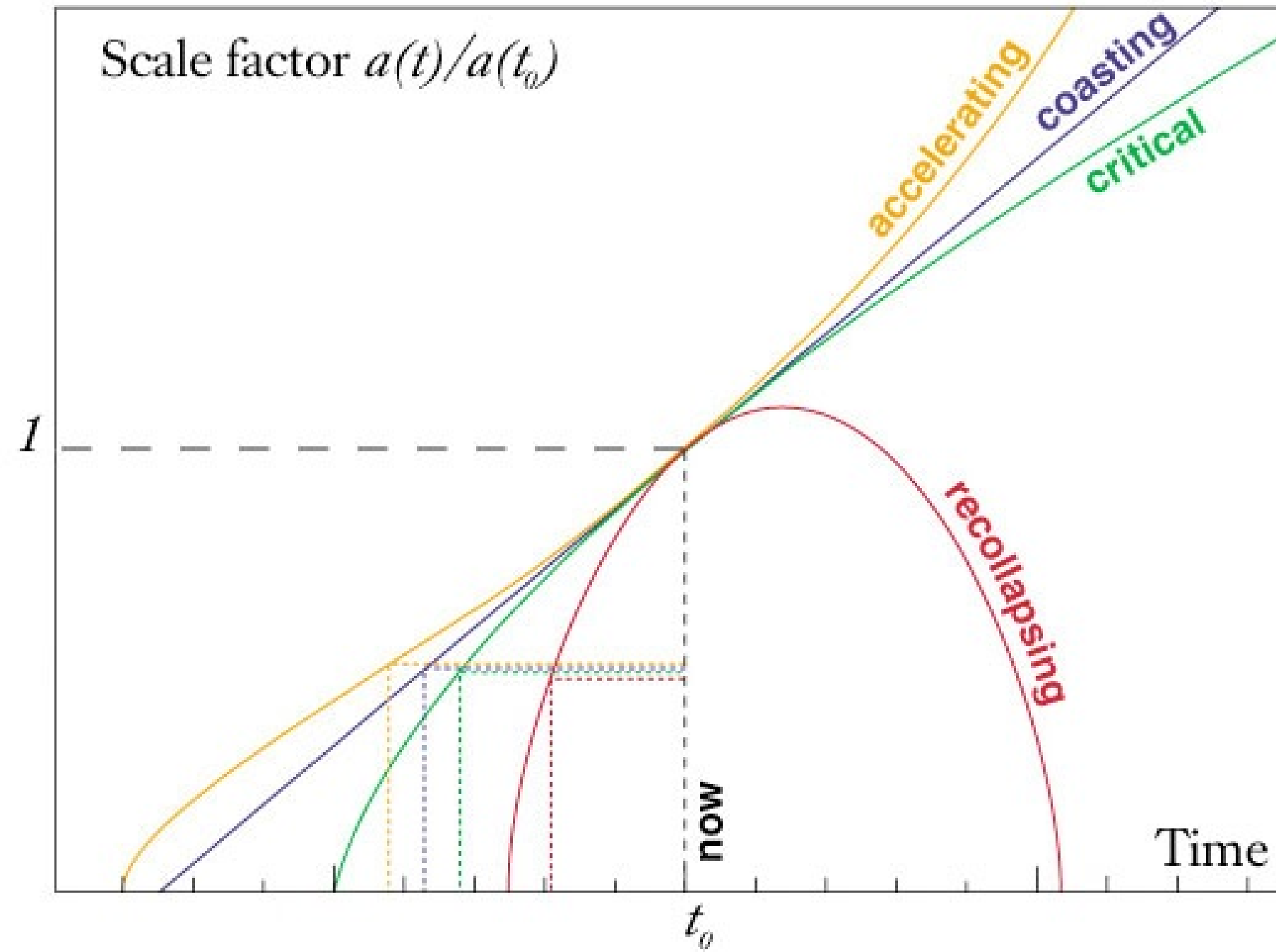
η problem & Lyth bound

Initial condition & entropy

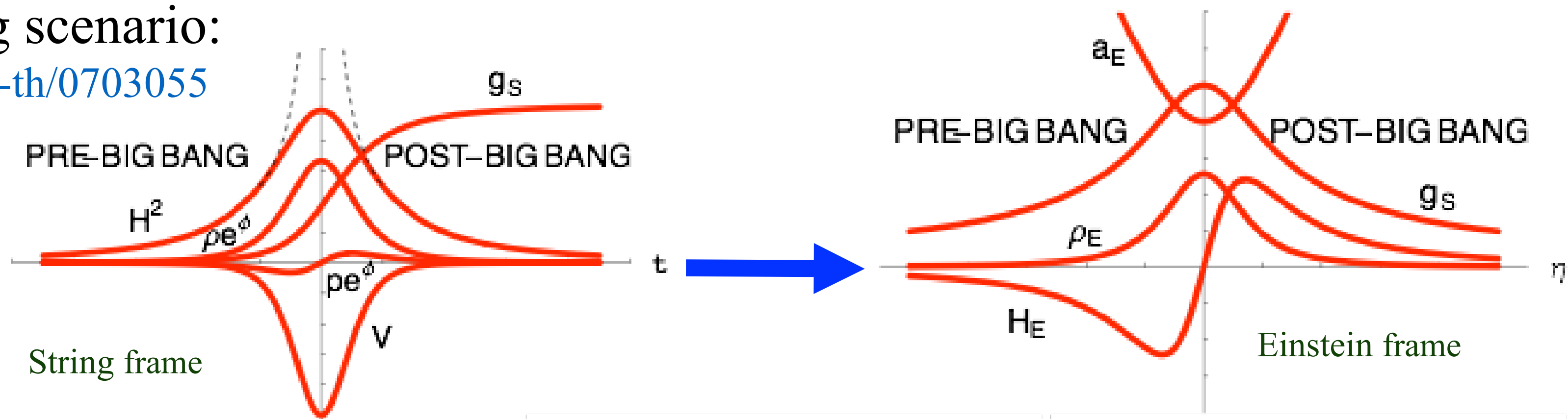
Eternal inflation & measure (anthropic)

Singularity problem

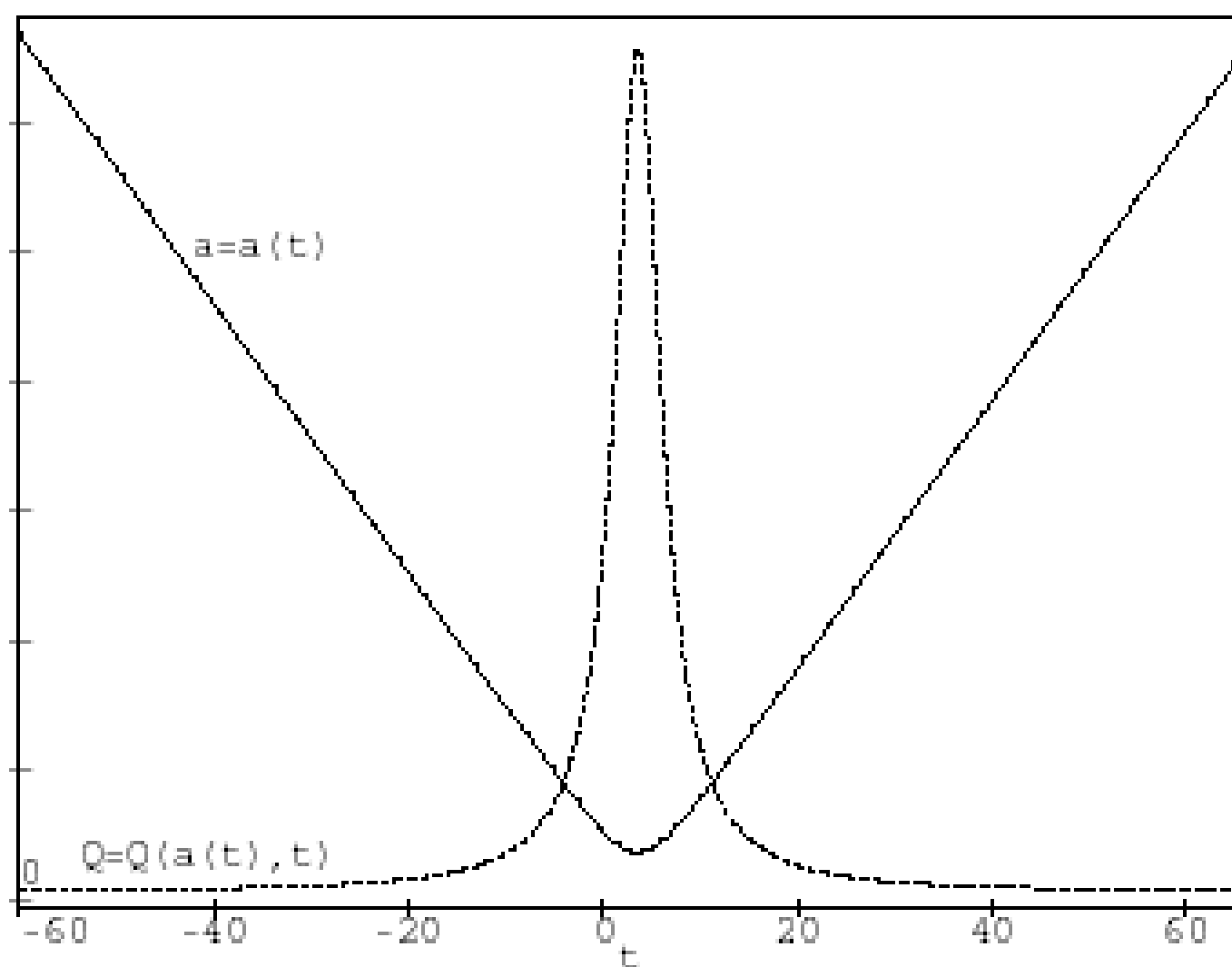
Quantum effect?



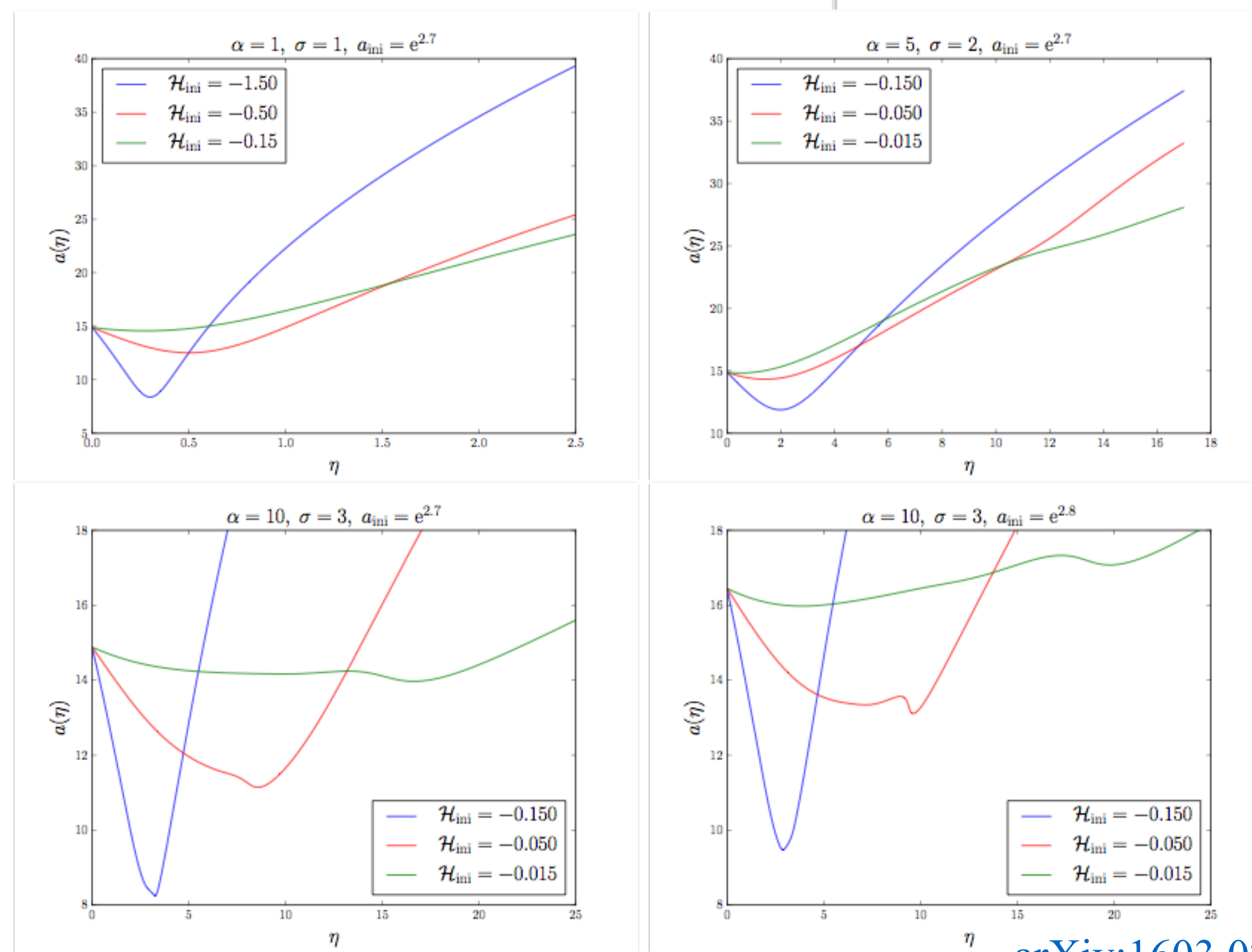
Pre Big Bang scenario:
[hep-th/0207130](http://arxiv.org/abs/hep-th/0207130) & [hep-th/0703055](http://arxiv.org/abs/hep-th/0703055)



WdW quantum cosmology:



PLA **241**, 229 (1998)



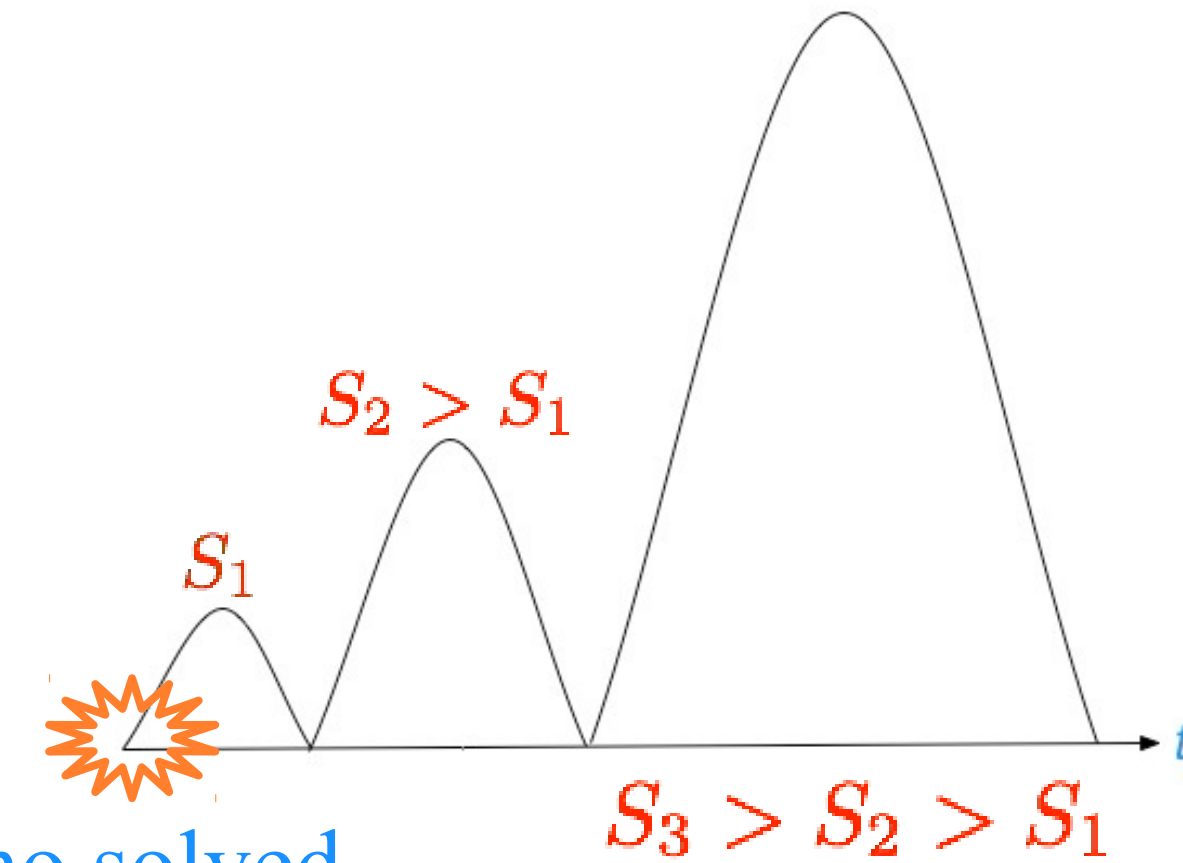
[arXiv:1603.02342](http://arxiv.org/abs/1603.02342)

A brief history of bouncing cosmology

→ R. C. Tolman, “*On the Theoretical Requirements for a Periodic Behaviour of the Universe*”, PRD 38, 1758 (1931)

→ G. Lemaître, “*L’Univers en expansion*”, Ann. Soc. Sci. Bruxelles (1933)

...



...

Singularity pb no solved

→ A. A. Starobinsky, “*On one non-singular isotropic cosmological model*”, Sov. Astron. Lett. 4, 82 (1978)

→ V. N. Melnikov, S.V. Orlov, Phys. Lett. A 70, 263 (1979).

→ R. Durrer & J. Laukerman, “*The oscillating Universe: an alternative to inflation*”, Class. Quantum Grav. 13, 1069 (1996)

→ Many new ideas, models...

→ M. Novello & S.E. Perez Bergliaffa, “*Bouncing cosmologies*”, Phys. Rep. 463, 127 (2008)

→ D. Battefeld & PP, “*A Critical Review of Classical Bouncing Cosmologies*”, Phys. Rep. 571, 1 (2015)

→ R. Brandenberger & PP, “*Bouncing cosmologies: Progress and problems*”, Found. Phys. (2017)

Model listing:

Quantum gravity

LQG & LQC

Canonical quantum gravity (WdW)

String theory

Non relativistic quantum gravity

Model listing:

Quantum gravity

LQG & LQC

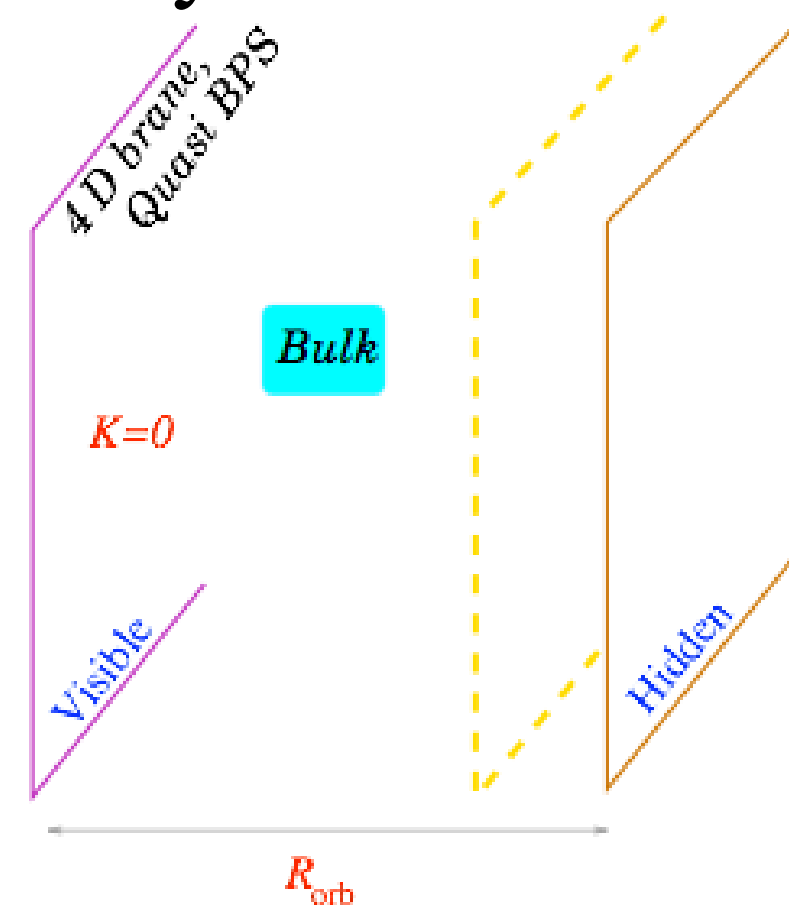
Non relativistic quantum gravity

Canonical quantum gravity (WdW)

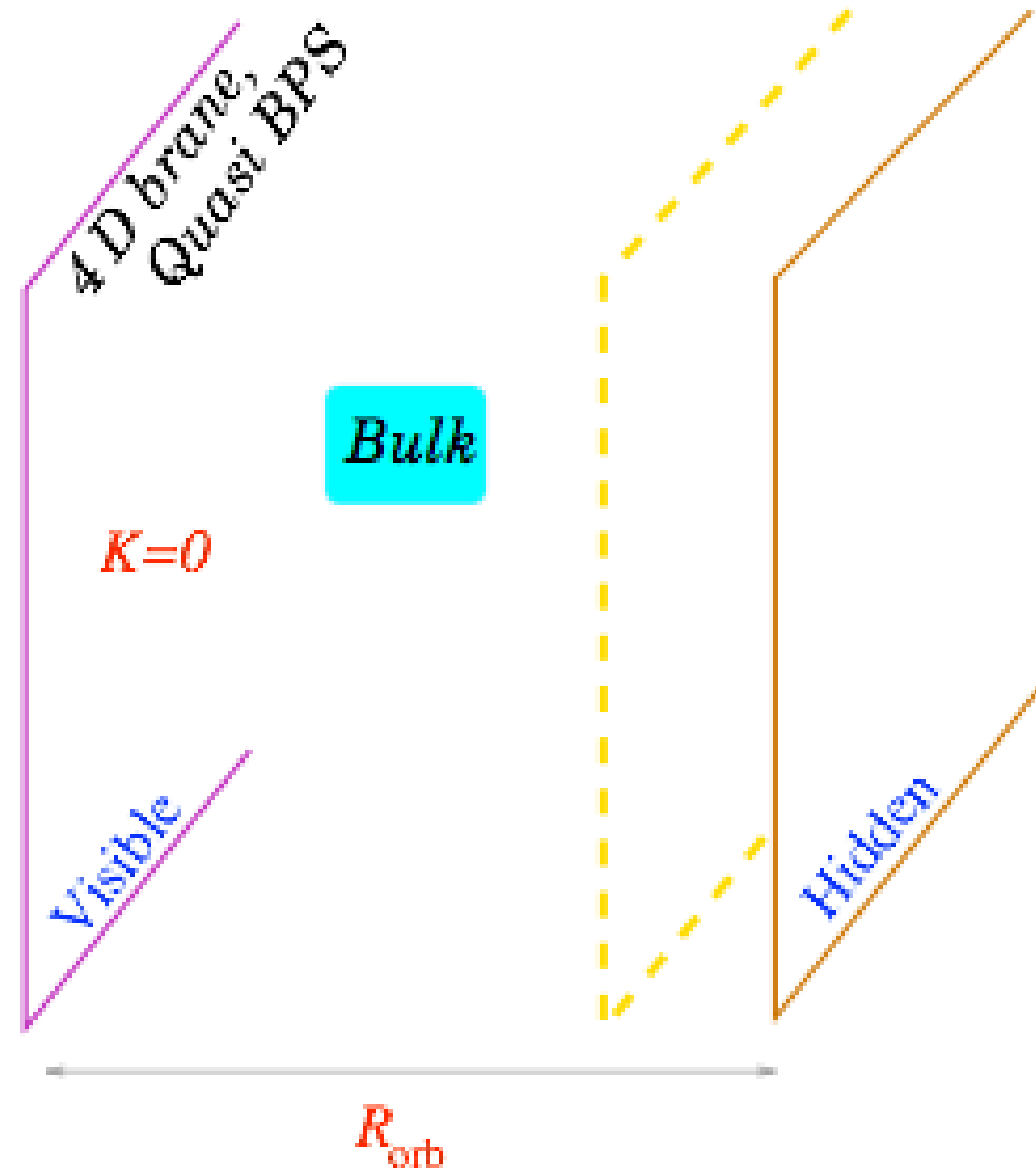
Ekpyrotic & cyclic

String theory

Branes



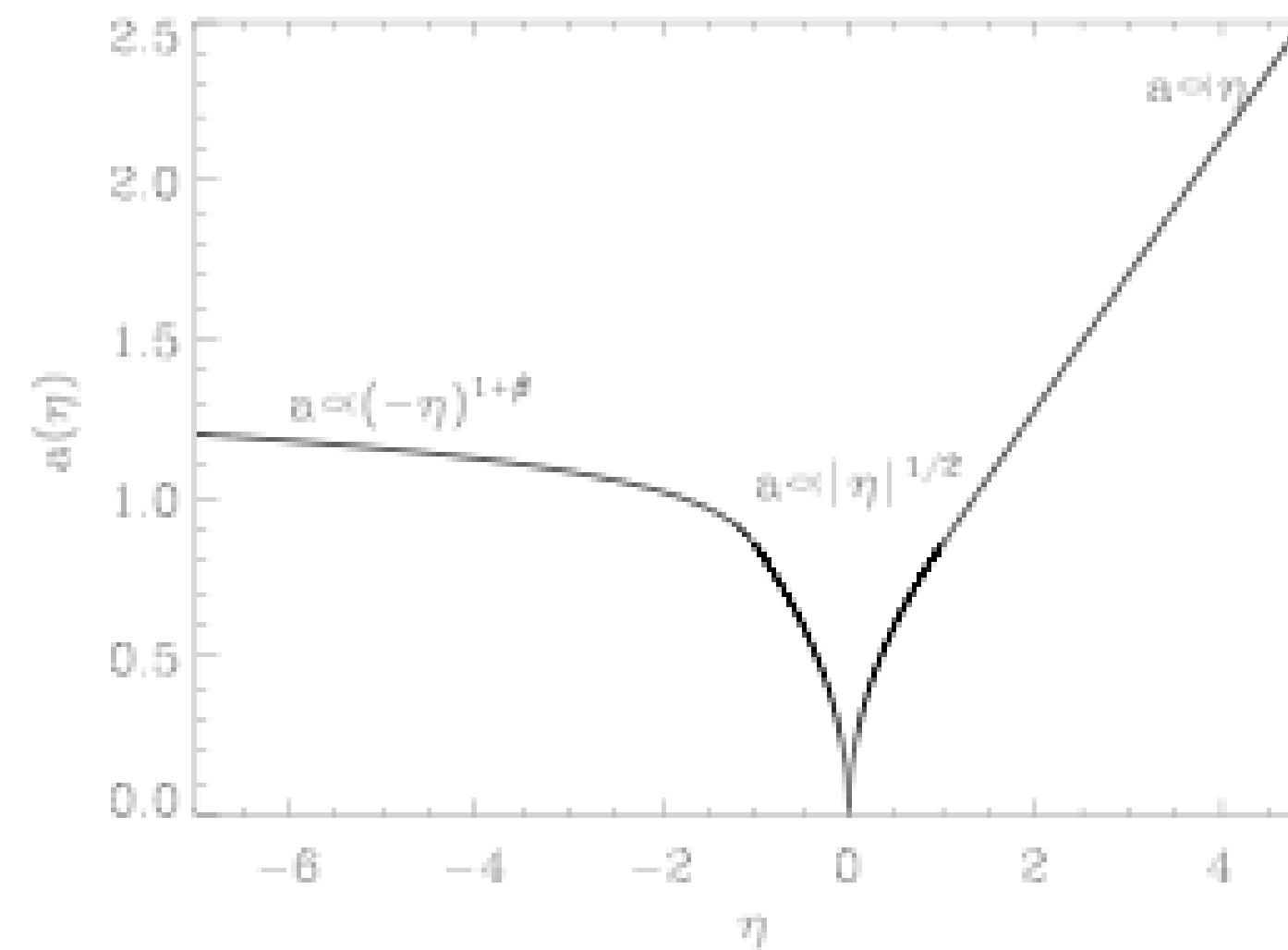
Ekpyrotic scenario:



$$\mathcal{S}_5 \propto \int_{\mathcal{M}_5} d^5x \sqrt{-g_5} \left[R_{(5)} - \frac{1}{2} (\partial\varphi)^2 - \frac{3}{2} \frac{e^{2\varphi} \mathcal{F}^2}{5!} \right],$$

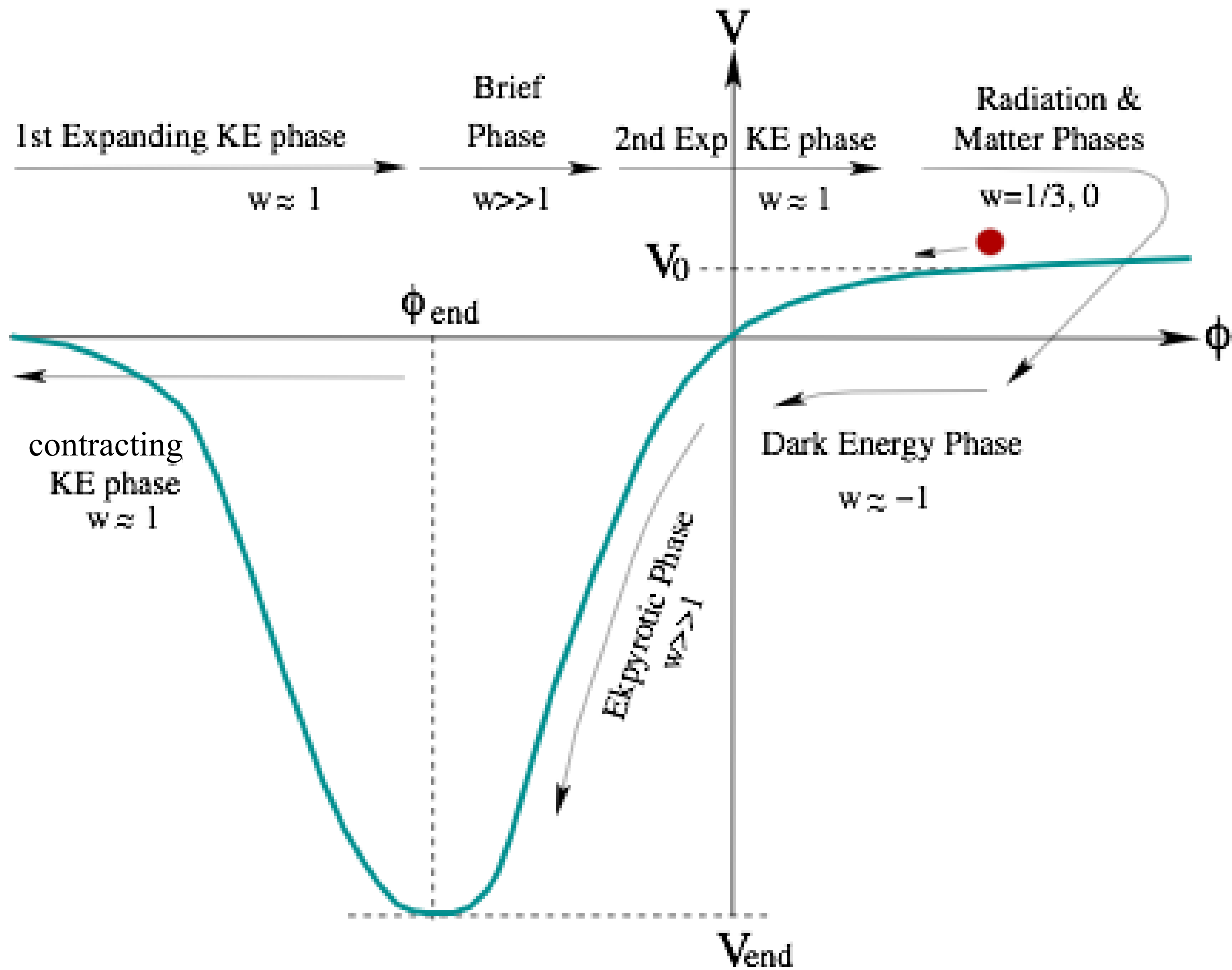
$$\mathcal{S}_4 = \int_{\mathcal{M}_4} d^4x \sqrt{-g_4} \left[\frac{R_{(4)}}{2\kappa} - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right],$$

$$V(\phi) = -V_i \exp \left[-\frac{4\sqrt{\pi\gamma}}{m_{Pl}} (\phi - \phi_i) \right],$$

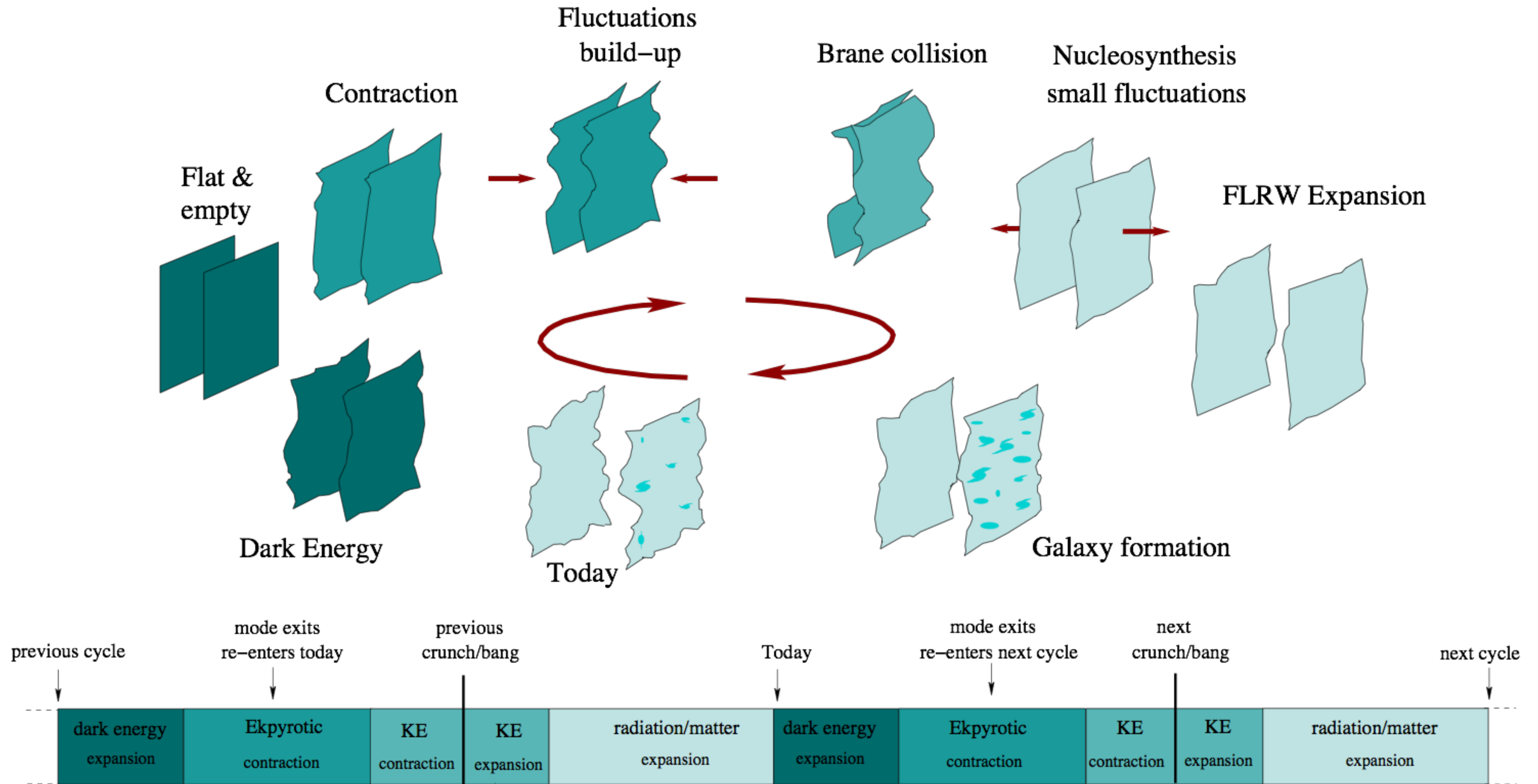


Singular...

BOUNCE



Cyclic extension



Model listing:

Quantum gravity

LQG & LQC

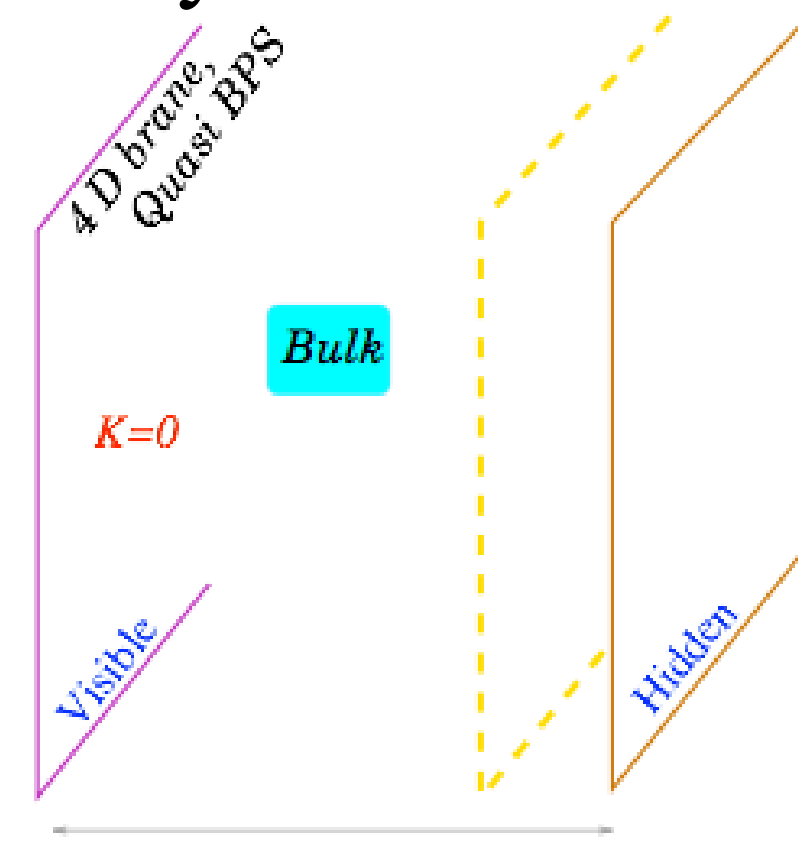
Non relativistic quantum gravity

Canonical quantum gravity (WdW)

Ekpyrotic & cyclic

String theory

Branes



Horava-Lifshitz

String gas cosmology

Lee-Wick & Quintom

Antigravity

$F(R)$, $f(T)$, Gauss-Bonnet

Galileon

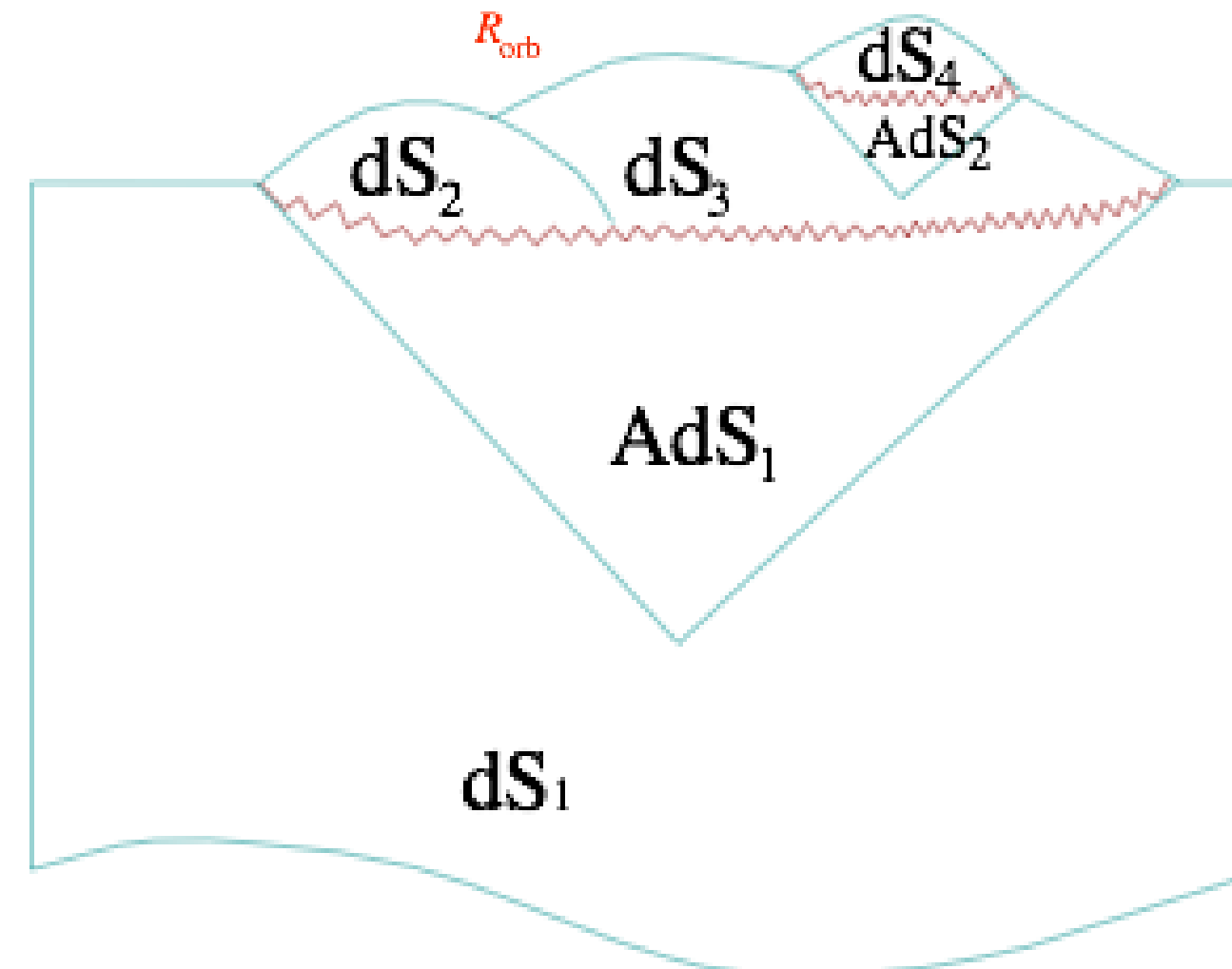
Massive gravity

Mimetic matter

Multiverse models

Non-linear electromagnetic action

Strings & AdS/CFT



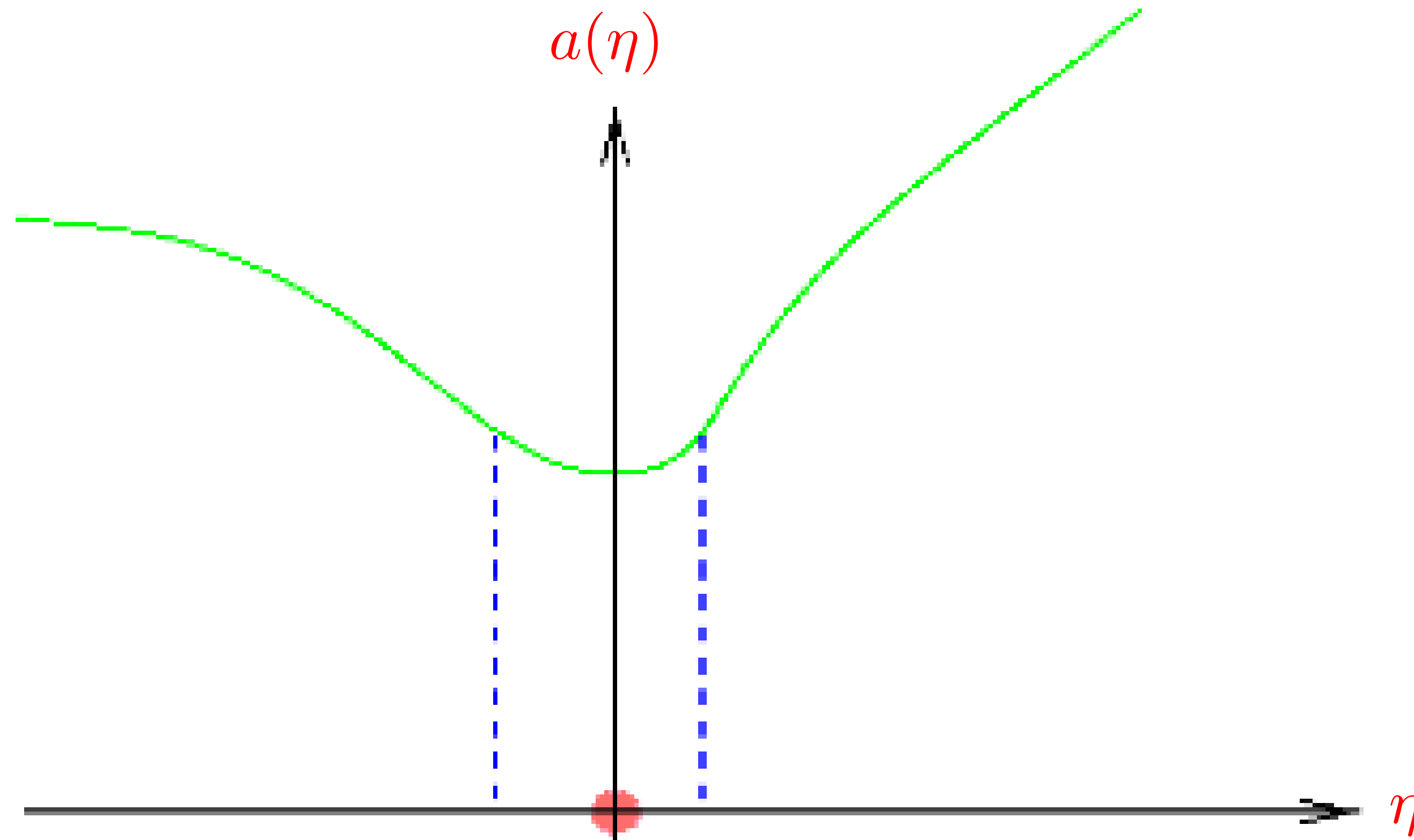
Spinors & torsion

Standard questions and bouncing answers

arXiv:0809.2022

Singularity

Merely a non issue in the bounce case!

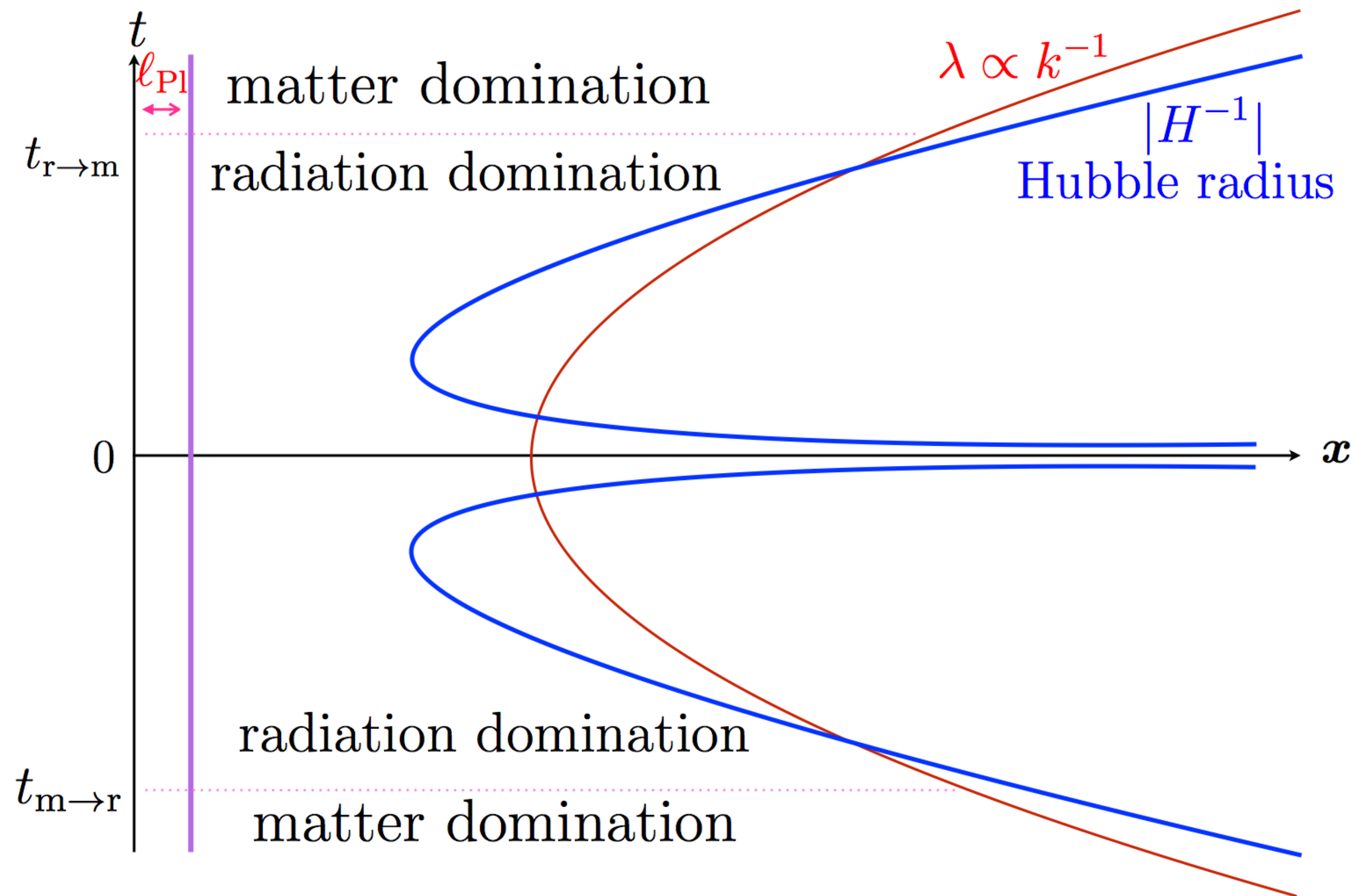


Standard questions and bouncing answers

arXiv:0809.2022

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Standard questions and bouncing answers

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Singularity

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Horizon $d_H \equiv a(t) \int_{t_{\text{ini}}}^t \frac{d\tau}{a(\tau)}$ can be made divergent easily if $t_{\text{ini}} \rightarrow -\infty$

Example:

$$d_H(t) = \frac{3(1+w)}{1+3w} |t| \left\{ 1 - \left(\frac{t_{\text{ini}}}{t} \right)^{(1+3w)/[3(1+w)]} \right\}$$
$$w > -\frac{1}{3} \implies \lim_{t_{\text{ini}} \rightarrow -\infty} d_H = \infty$$

Standard questions and bouncing answers

arXiv:0809.2022

Singularity

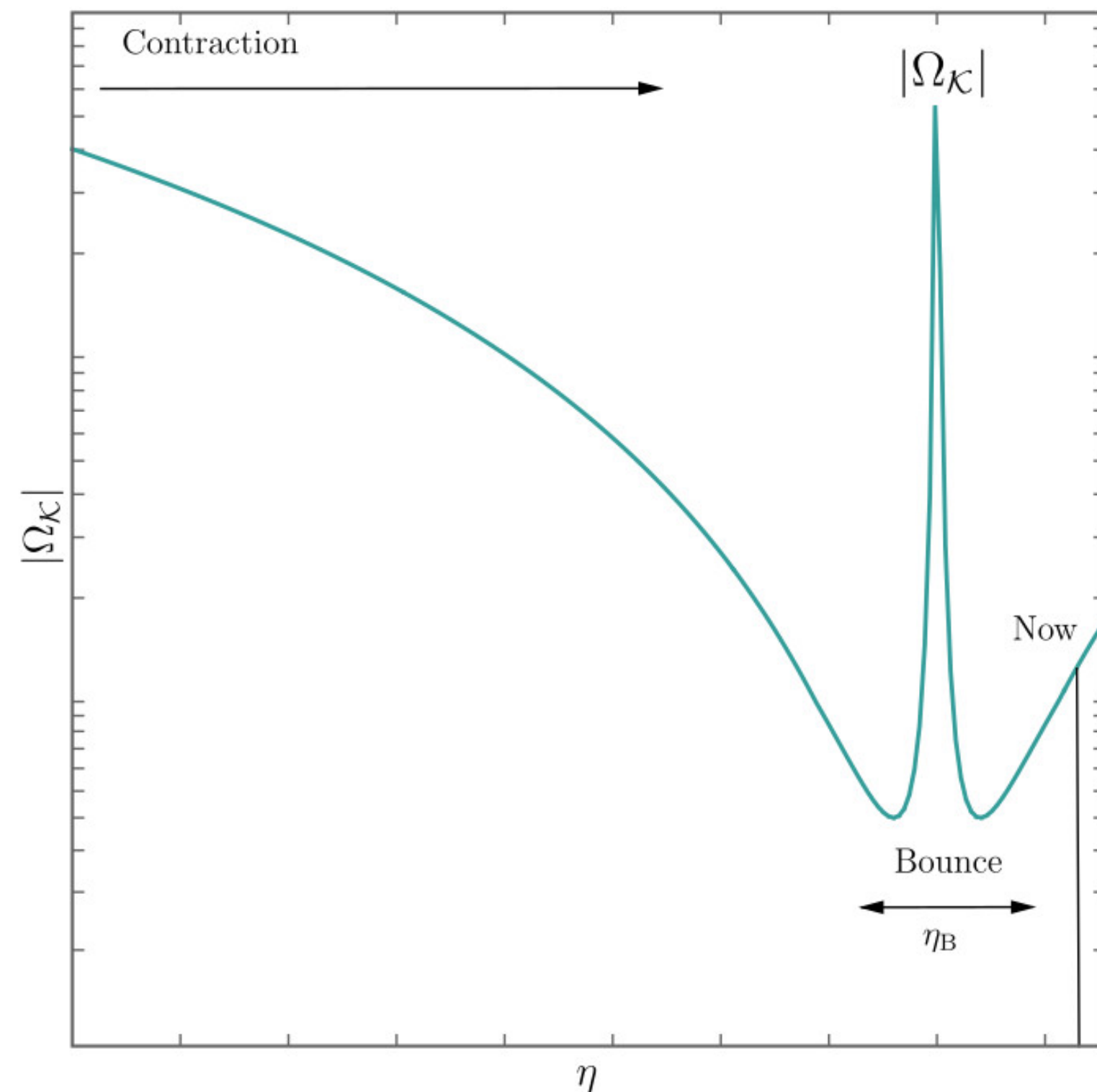
Merely a non issue in the bounce case!

Horizon $d_H \equiv a(t) \int_{t_{\text{ini}}}^t \frac{d\tau}{a(\tau)}$ can be made divergent easily if $t_{\text{ini}} \rightarrow -\infty$

Flatness

$$\frac{d}{dt} |\Omega - 1| = -2 \frac{\ddot{a}}{\dot{a}^3} \longrightarrow \ddot{a} < 0 \ \& \ \dot{a} < 0$$

decelerated contraction (**bounce**)
[accelerated expansion (**inflation**)]



Standard questions and bouncing answers

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accelerated expansion (**inflation**) or decelerated contraction (**bounce**)

Homogeneity

Large & flat Universe + low initial density + diffusion

$t_{\text{dissipation}}(\lambda) \ll t_{\text{Hubble}}$ enough time to dissipate any wavelength
 \implies quantum vacuum fluctuations...

Implementing a bounce = problem with GR!

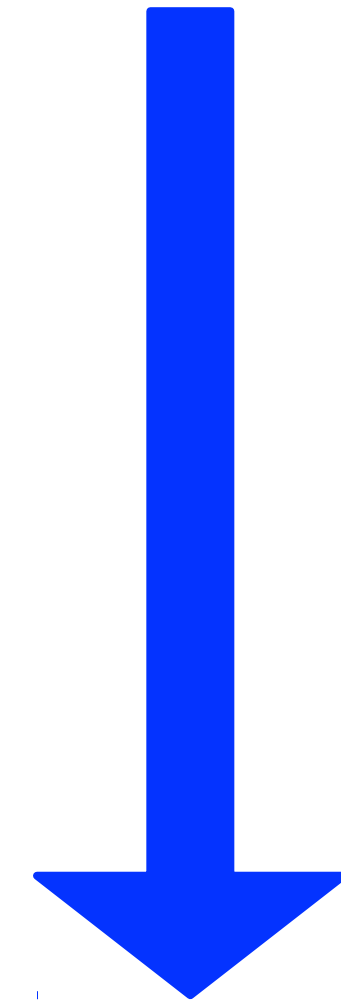
$$\dot{H} = \frac{\mathcal{K}}{a^2} - \frac{1}{2}(\rho + p)$$

Null Energy Condition (NEC)

$$\rho + p \leq 0$$

Many working examples:

K-bounce, Ghost condensates, Galileons...?



Instabilities

Standard questions and bouncing answers

arXiv:0809.2022

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Isotropy

Potentially problematic: model dependent

The problem with contraction: BKL/shear instability

Ricci flat Bianchi I anisotropic universe

$$ds^2 = -dt^2 + a^2(t) \sum_i e^{2\theta_i(t)} (dx^i)^2$$

$$\sum_i \theta_i = 0$$

Average scale factor

$$H_i \equiv \frac{1}{ae^{\theta_i}} \frac{d}{dt} (ae^{\theta_i}) = H + \dot{\theta}_i$$

$$\equiv \frac{\dot{a}}{a} \text{ Mean Hubble parameter}$$

Friedman equations

$$H^2 = \frac{\rho}{3M_P^2} + \frac{1}{6} \sum_i \dot{\theta}_i^2$$

$$\dot{H} = -\frac{\rho + p}{2M_P^2} - \frac{1}{2} \sum_i \dot{\theta}_i^2$$

$$\ddot{\theta}_i + 3H\dot{\theta}_i = 0 \implies \dot{\theta}_i \propto a^{-3}$$

$$\rho_{\text{shear}} \propto a^{-6}$$

Ekpyrotic solution: $w_{\text{ekp}} \gg 1$

$$\longrightarrow \rho_{\text{ekp}} \propto a^{-3(1+w_{\text{ekp}})} \gg a^{-6} \text{ when } a \rightarrow 0$$

Problem: regular bounce $\xrightarrow{\text{NEC}}$ phase with $w_{\text{bounce}} < -1$

So finally...

$$\rho_{\text{shear}} = \frac{M_{\text{P}}^2}{2} \sum_i \dot{\theta}_i^2 \propto a^{-6} \gg \rho_{\text{fluid}}$$

\longrightarrow back to the singularity!

A nonsingular bounce model: ghost condensate & Galileon

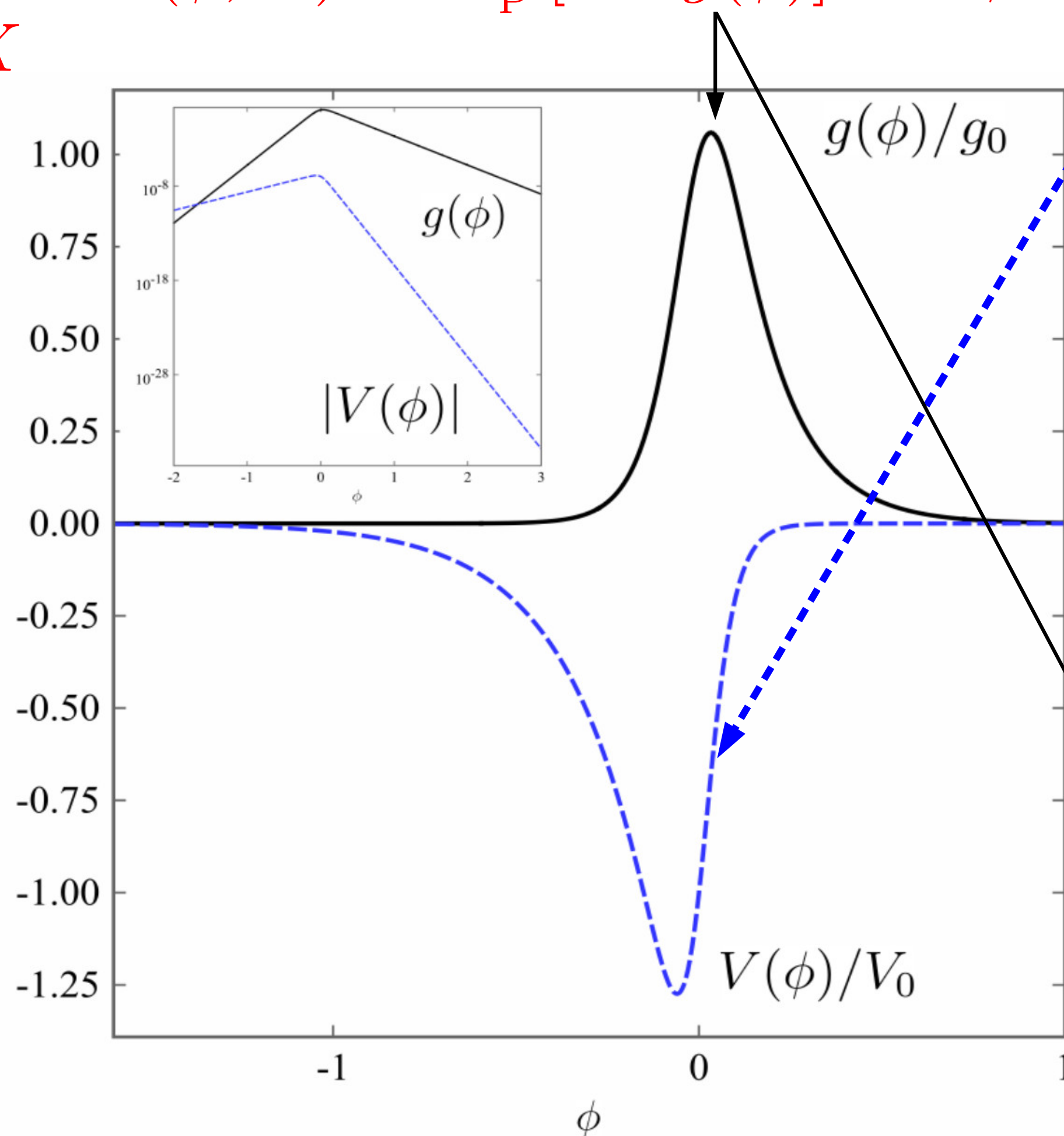
$$\mathcal{L}[\phi(x)] = K(\phi, X) + G(\phi, X) \square\phi$$

$$X \equiv \frac{1}{2} \partial_\mu \phi \partial^\mu \phi$$

kinetic term

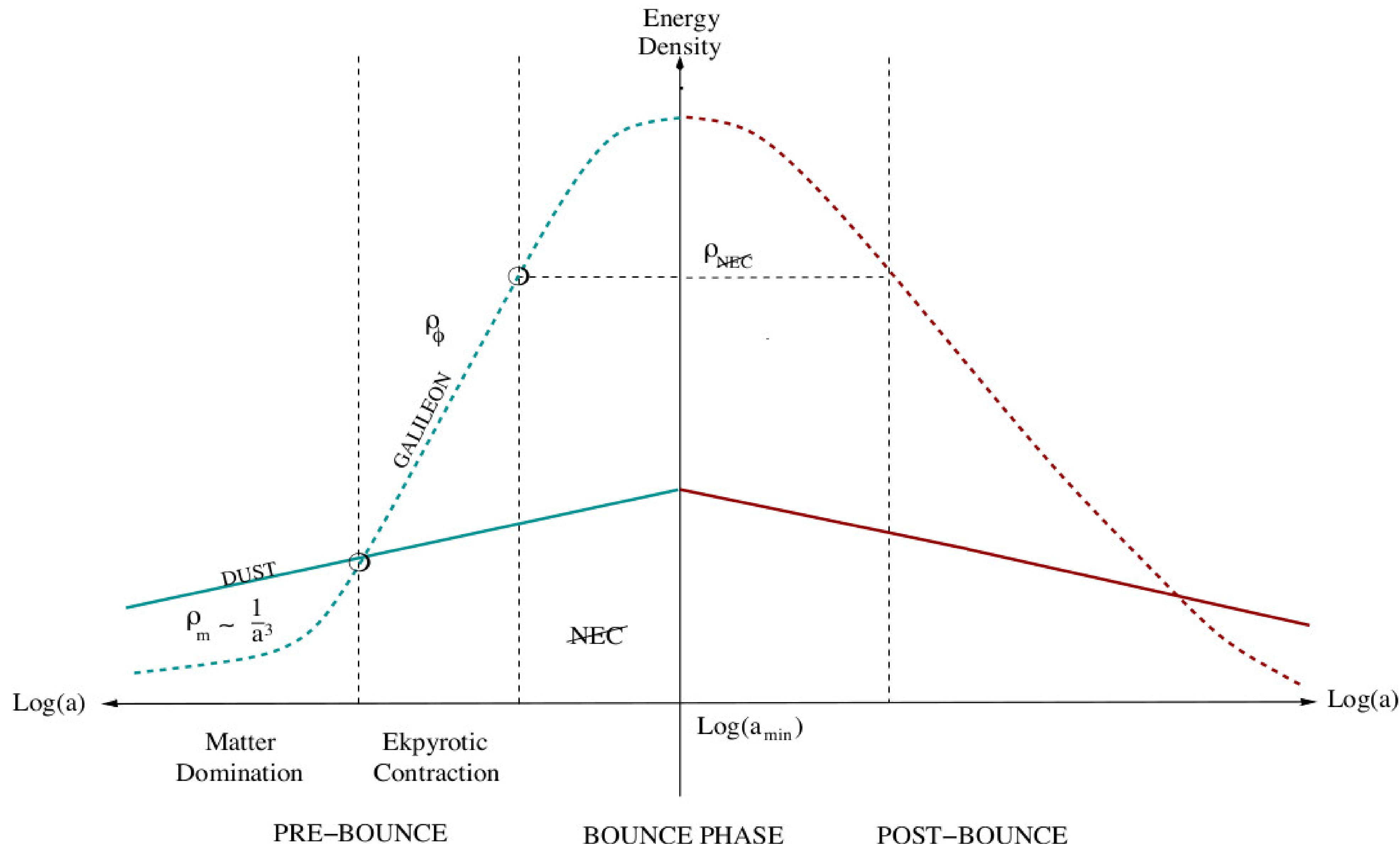
Specific choices: $K(\phi, X) = M_{\text{P}}^2 [1 - g(\phi)] X + \beta X^2 - V(\phi)$

$$G(X) = \gamma X$$

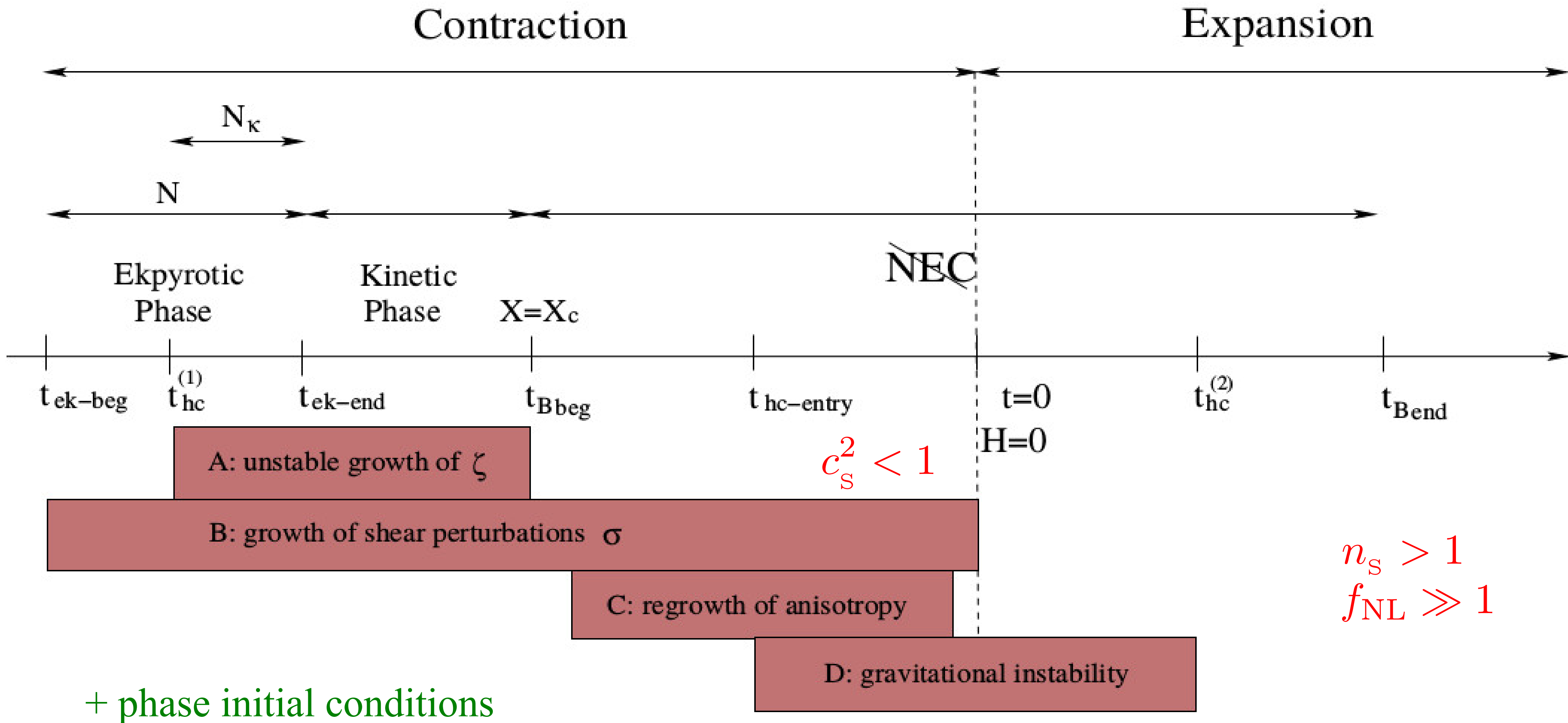


$$V(\phi) = \frac{-2V_0}{e^{-\alpha_q \phi} + e^{b_v \alpha_q \phi}}$$

$$g(\phi) = \frac{2g_0}{e^{-\alpha_p \phi} + e^{b_g \alpha_p \phi}}$$



Summary of possible problems



Standard questions and bouncing answers

arXiv:0809.2022

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dark matter/energy, baryogenesis, ...