



DATRICK DETER GRECO Institut d'Astrophysique de Paris Institut Lagrange de Paris

Home



DARK MATTERS 11-13 December 2017 Joe Silk's 75th Birthday



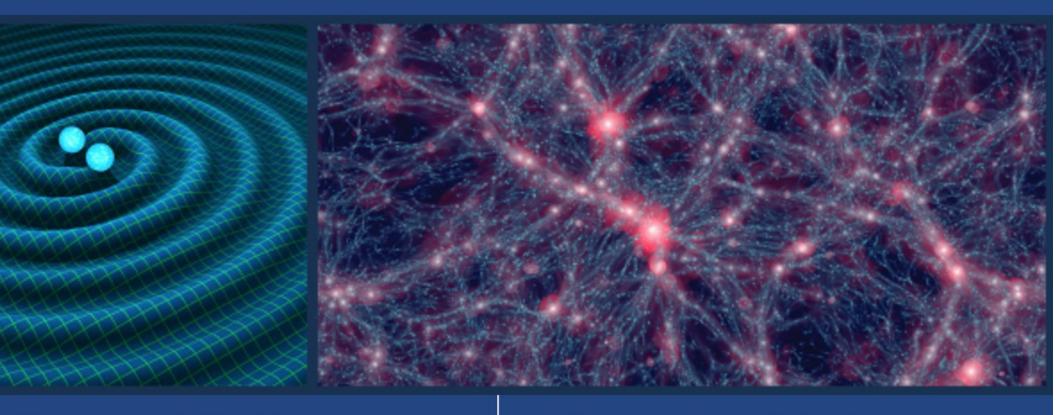
Participants

Monday 11th December

Evidences for inflation constraints on alternatives

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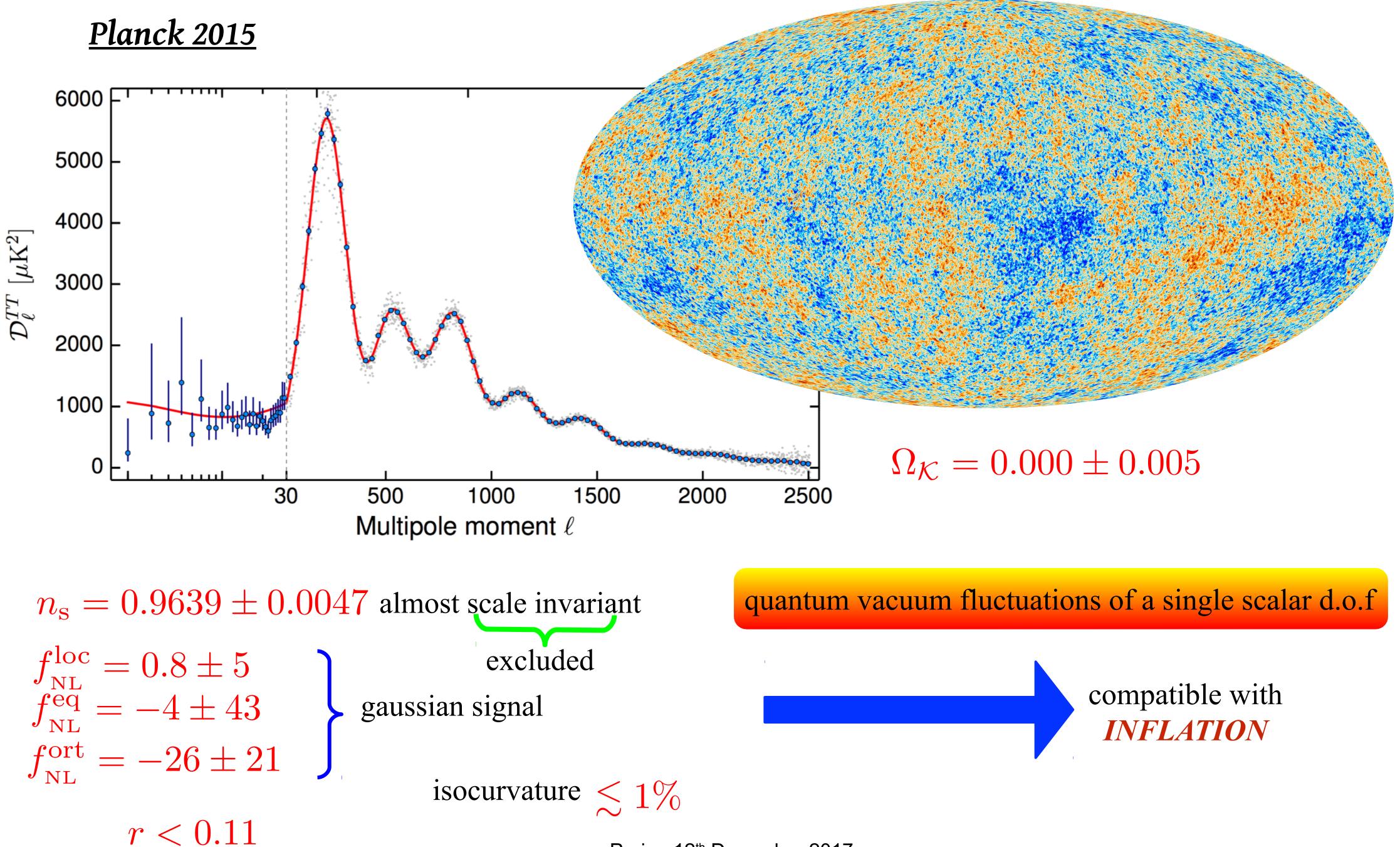




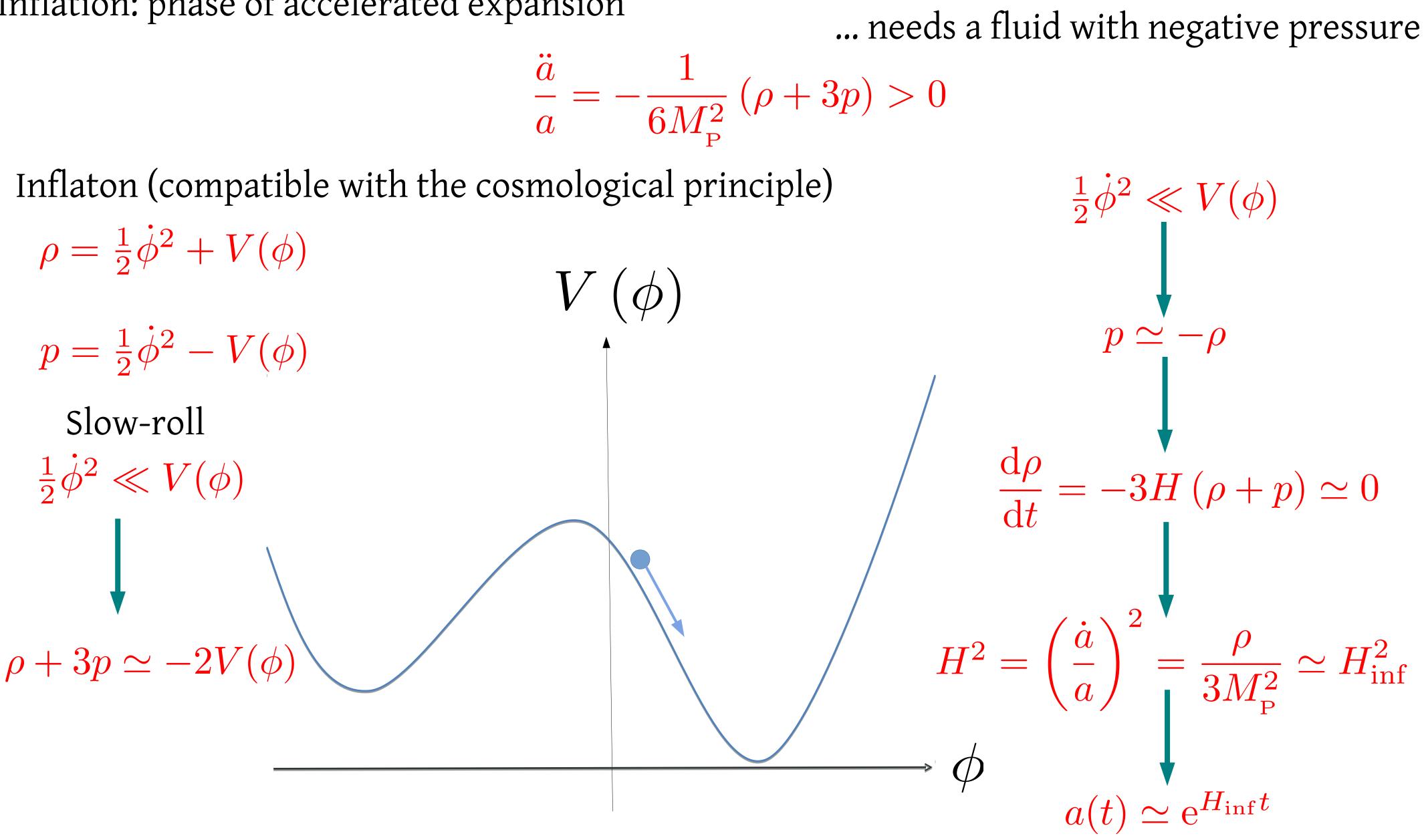
Tuesday 12th December

Wednesday 13th December

THANKS TO JÉDÔME MADTIN FOD HIS HELD



Inflation: phase of accelerated expansion



Standard questions and inflationary answers

Singularity Not solved... actually not addressed! **Horizon** $d_{\rm H} \equiv a(t) \int_{t_{\rm ini}}^{t} \frac{\mathrm{d}\tau}{a(\tau)}$ can be made as big as one wishes **Flatness** $\frac{\mathrm{d}}{\mathrm{d}t} |\Omega - 1| = -2\frac{\ddot{a}}{\dot{a}^3} \longrightarrow \ddot{a} > 0 \& \dot{a} > 0$

Homogeneity & Isotropy

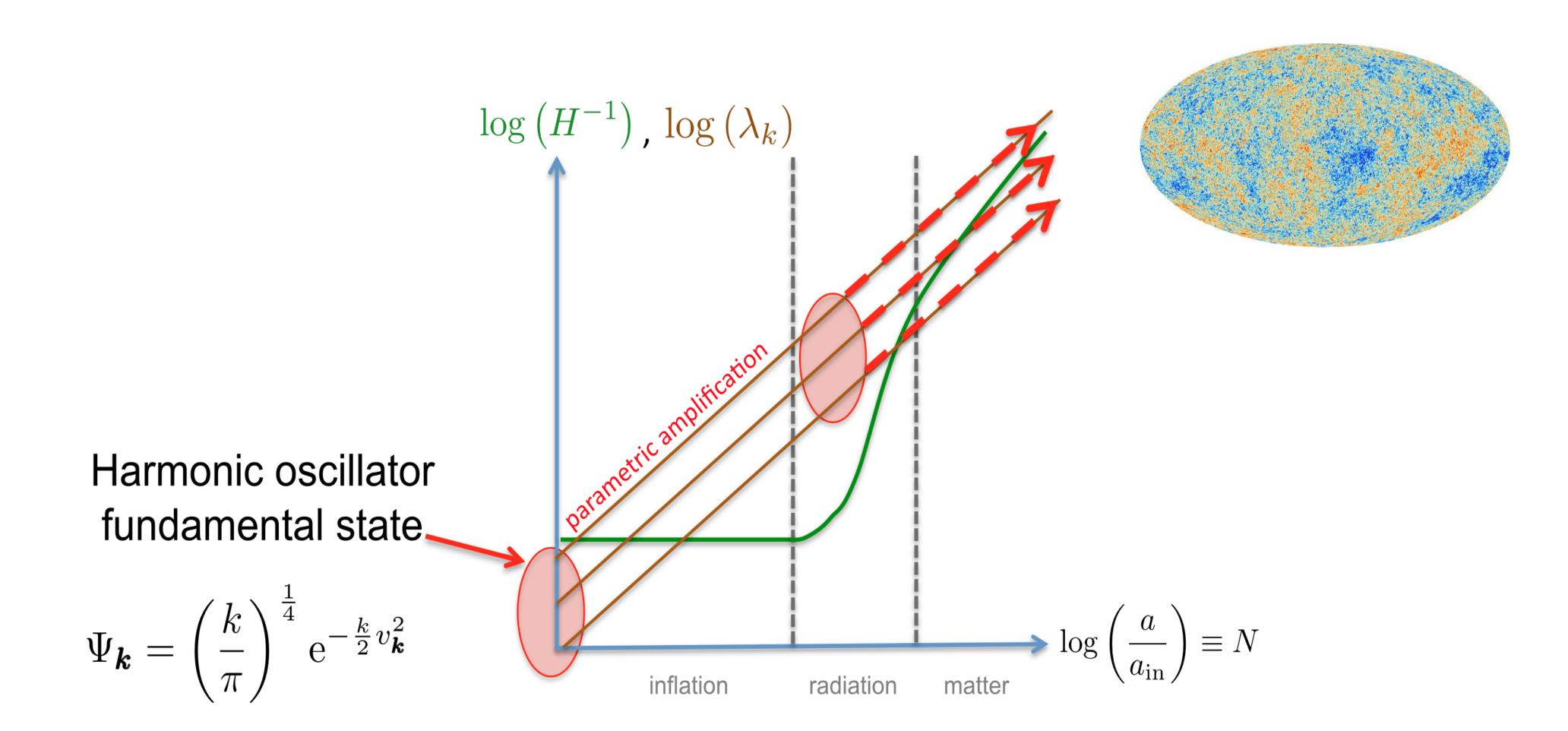
Initial Universe = very small patch vacuum state! Accelerated expansion drives the shear to zero...

Perturbations Bonus of the theory: predictions!!! **Others** dark matter/energy, baryogenesis,...

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accelerated expansion (inflation)

Almost constant Hubble radius $R_{\rm H} = H_{\rm inf}^{-1}$



Everything depends on the slow roll parameters

$$\epsilon_1 \simeq \frac{1}{2M_{\rm P}^2} \left(\frac{V_{,\phi}}{V}\right)^2$$

$$\epsilon_2 \simeq \frac{2}{M_{\rm P}^2} \left[\left(\frac{V_{,\phi}}{V} \right)^2 - \frac{V_{,\phi\phi}}{V} \right]$$

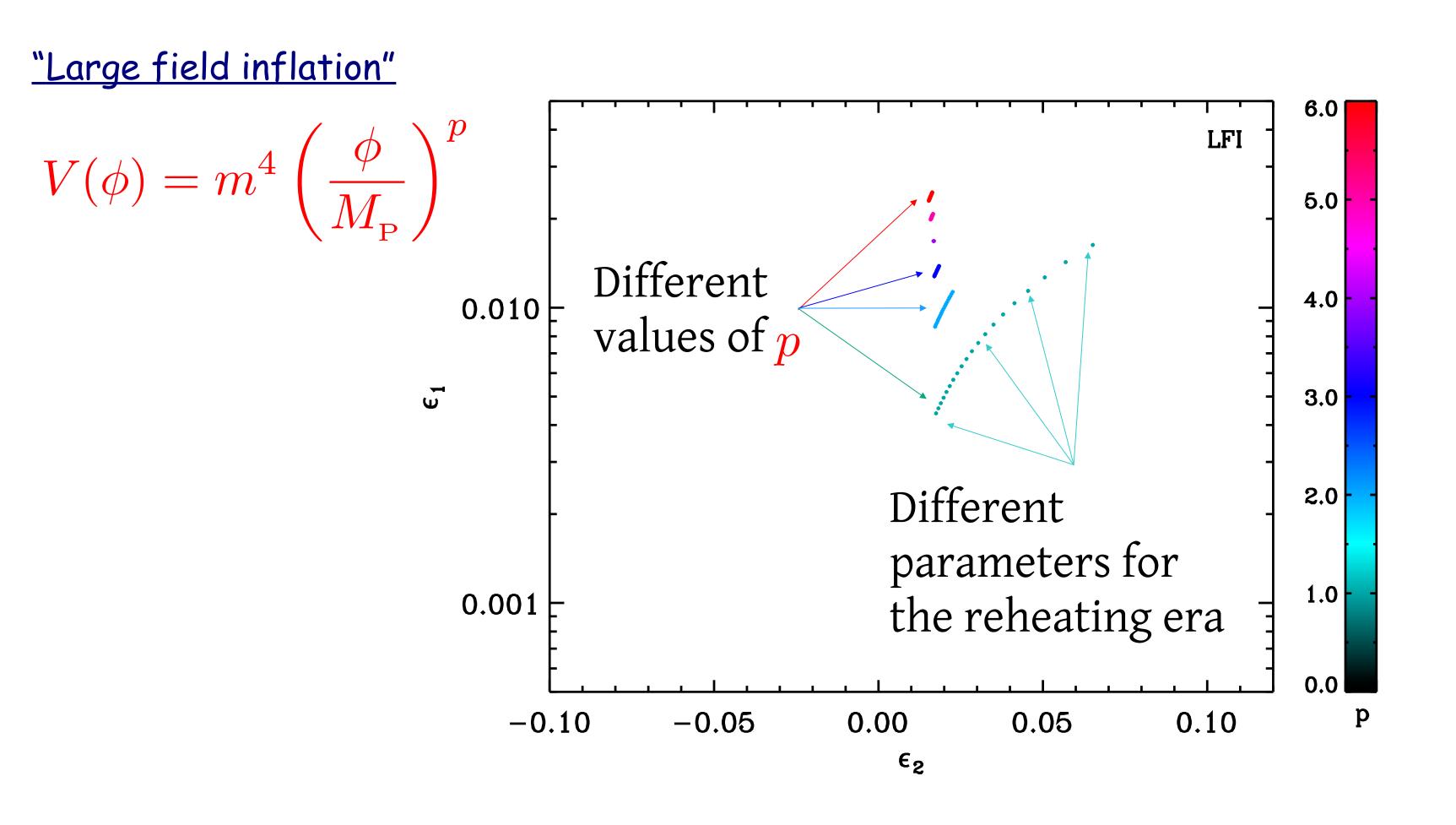
Evaluated at ϕ_{\star} (pivot scale leaves Hubble radius):

$$\mathcal{P}_{\zeta} \propto \frac{H_{\inf}^2}{M_{P}^2 \epsilon_{1\star}} \left[1 + \mathcal{O}\left(k, \epsilon_{\star}\right) \right]$$
 Spectral indices
$$\begin{cases} n_{s} = 1 + \frac{d \ln \mathcal{P}_{\zeta}}{d \ln k} = -2\epsilon_{1\star} - \epsilon_{2\star} \\ n_{T} = \frac{d \ln \mathcal{P}_{h}}{d \ln k} = -2\epsilon_{1\star} \end{cases}$$

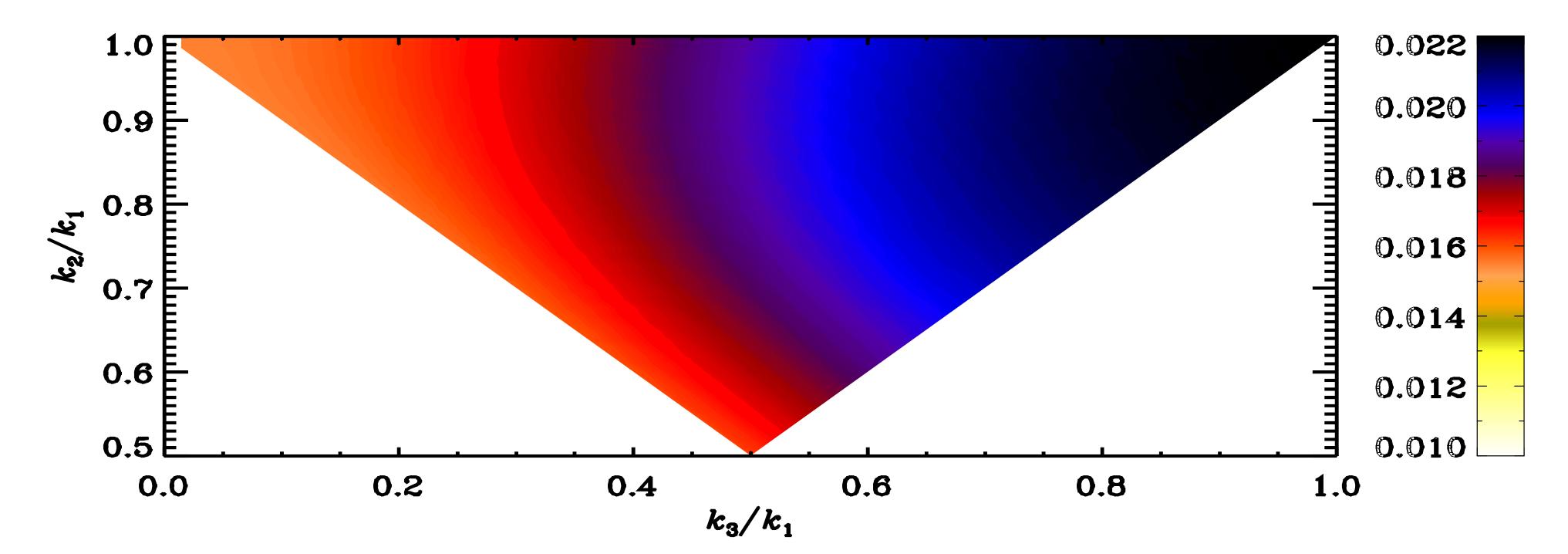
Consistency relation
$$r \equiv \frac{T}{S} = \frac{\mathcal{P}_{\zeta}}{\mathcal{P}_{h}} = 16\epsilon_{1\star} = -8n_{\mathrm{T}}$$

+ running

Various models, various predictions... (one example)

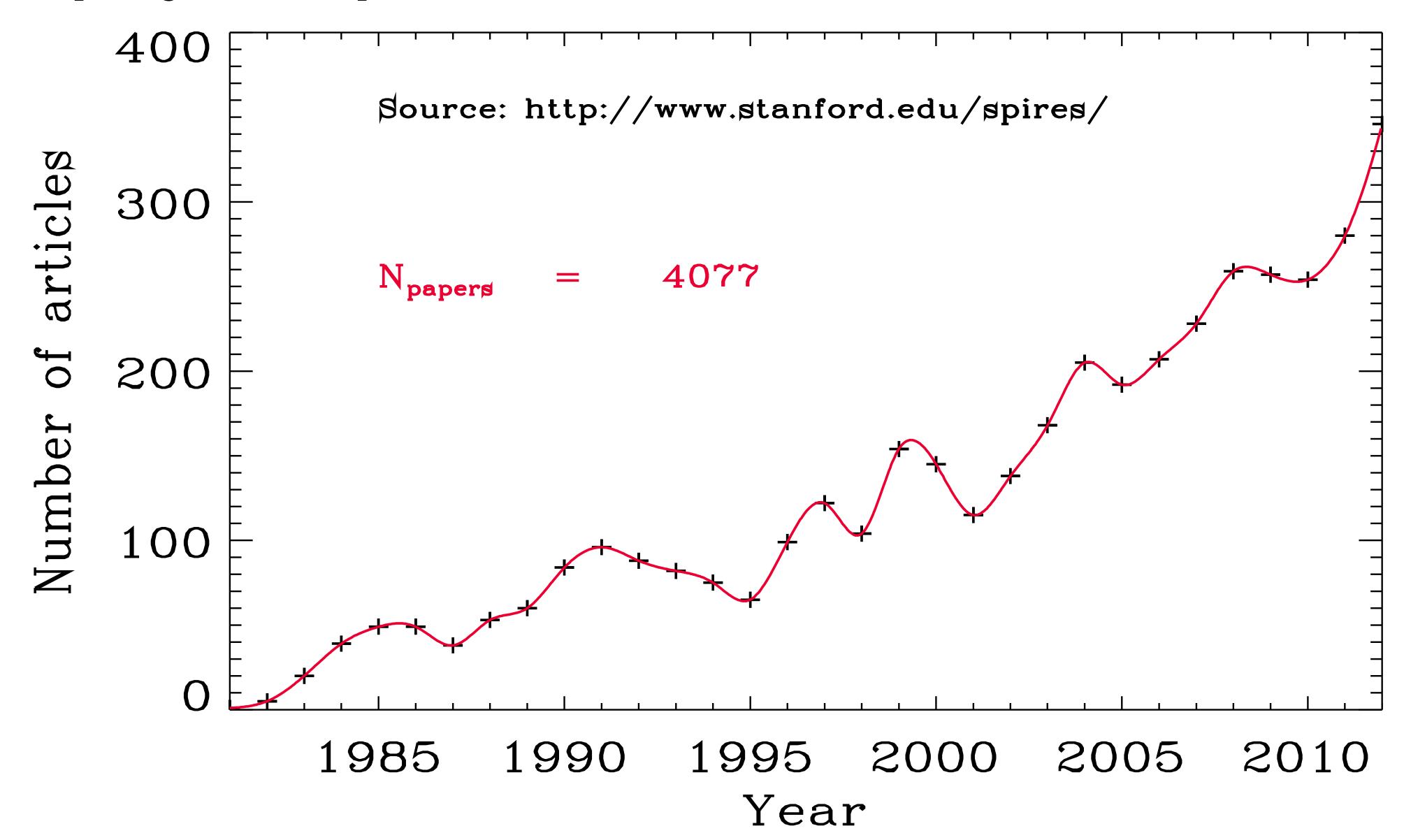


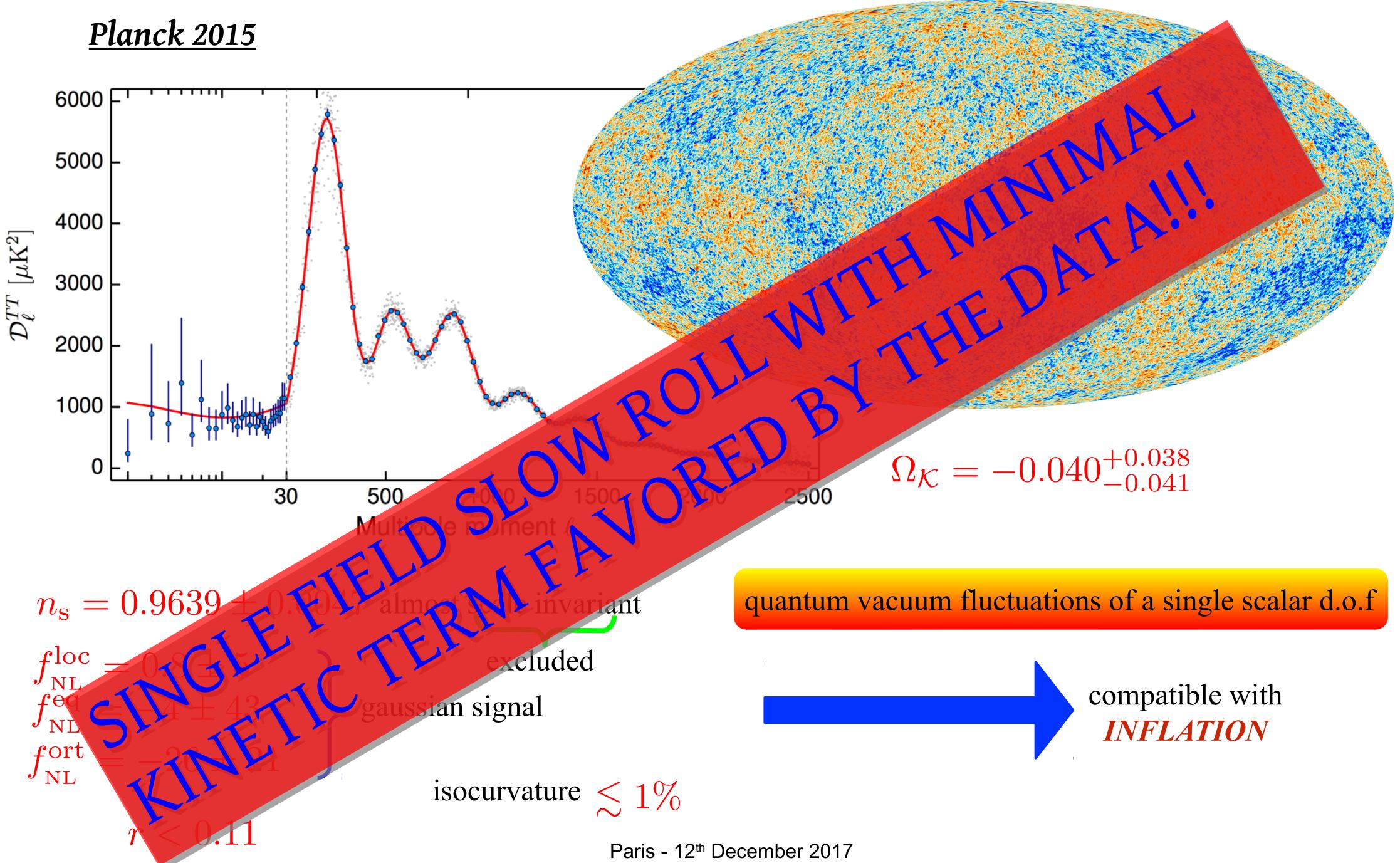
3 point correlation function...

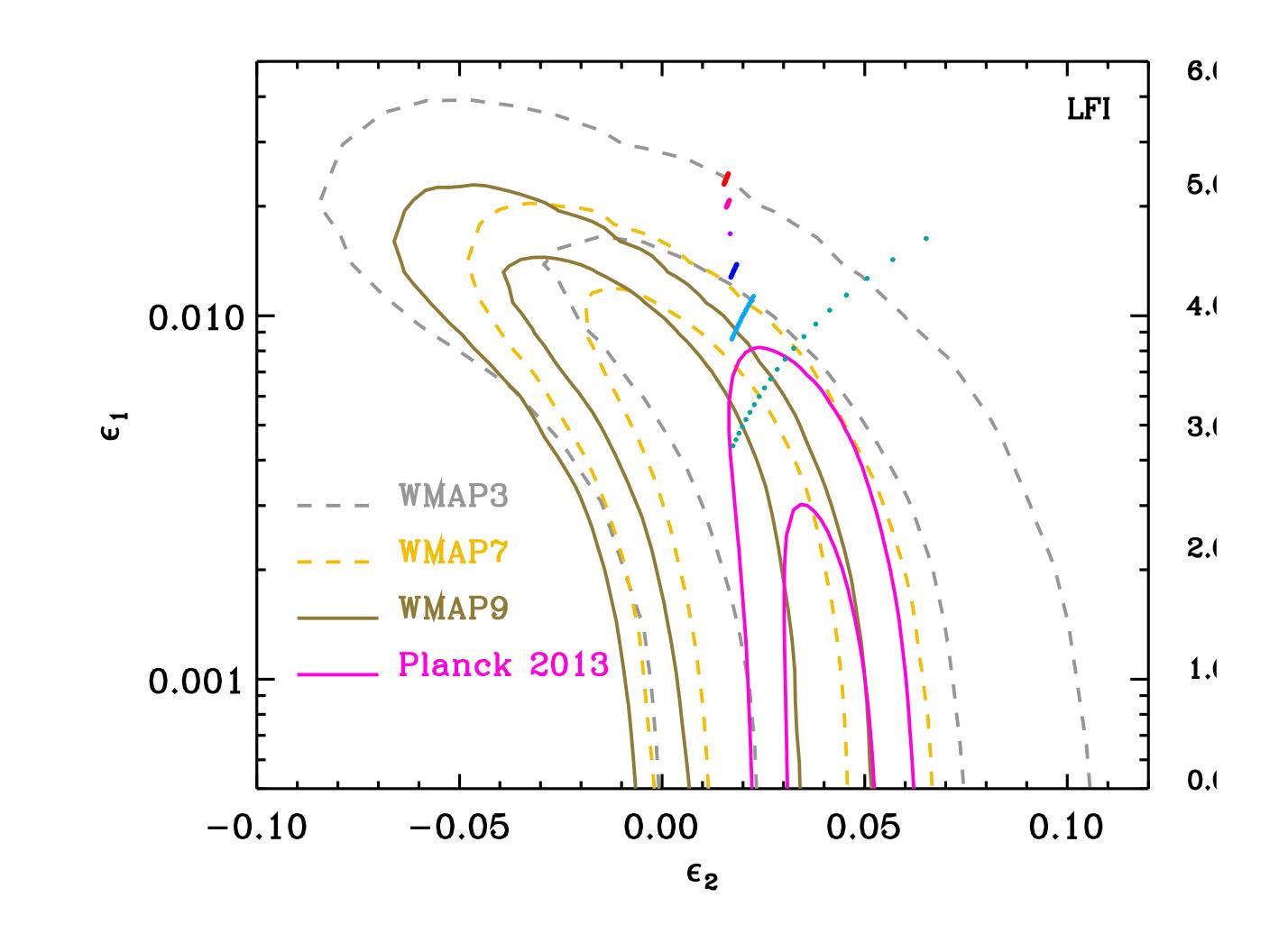


 $\left\langle \zeta_{\boldsymbol{k}_{1}}\zeta_{\boldsymbol{k}_{2}}\zeta_{\boldsymbol{k}_{3}}\right\rangle \propto f_{_{\mathrm{NL}}}\left(\boldsymbol{k}_{1},\boldsymbol{k}_{2},\boldsymbol{k}_{3}\right)\frac{\delta\left(\boldsymbol{k}_{1}+\boldsymbol{k}_{2}+\boldsymbol{k}_{3}\right)}{k_{1}^{3}k_{2}^{3}k_{3}^{3}}\left[k_{1}^{3}\mathcal{P}_{\zeta}\left(\boldsymbol{k}_{2}\right)\mathcal{P}_{\zeta}\left(\boldsymbol{k}_{3}\right)+2\text{ perms.}\right]$ Unobservable at the moment $\mathcal{O}(\epsilon_{\star})$

Comparing with data: pb = 100's of different models







CMB constraints on the slow-roll parameters

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Planck 2015 constraints:

Slow-roll para

Deviation from exact scale invariance $n_{\rm s} = 0.9645 \pm 0.0049 < 1$

Inflaton pote

No detection:

GW r < 0.1

Running $\frac{\mathrm{d}n_{\mathrm{s}}}{\mathrm{d}\ln k} = -0.0134 \pm 0.009$ Energy scale $H_{\rm inf}^2 \sim \mathcal{P}_{\zeta} \epsilon_{1\star}$

Bayesian analysis \implies best models: J. Martin, C. Ringeval & V. Vennin, arXiv:1303.3787

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ameters

$$\begin{aligned}
\epsilon_{1\star} < 0.0068 \\
\epsilon_{2\star} = 0.029^{+0.008}_{-0.007}
\end{aligned}$$
ntial

$$M_{\rm P} \frac{|V,\phi|}{V} < 0.14 \\
M_{\rm P} \frac{V,\phi\phi}{V} = 0.01^{+0.005}_{-0.009}
\end{aligned}$$

$$\epsilon_{2\star} \propto \rho_{\star} \qquad \rho_{\star} \qquad \rho_{\star} \sim \rho_{\star}$$

Encyclopedia Inflationaris

• solves cosmological puzzles

- uses GR + scalar fields [(semi-)classical]
- can be implemented in high energy theories • string implementation (brane inflation, ...)

• makes falsifiable predictions...

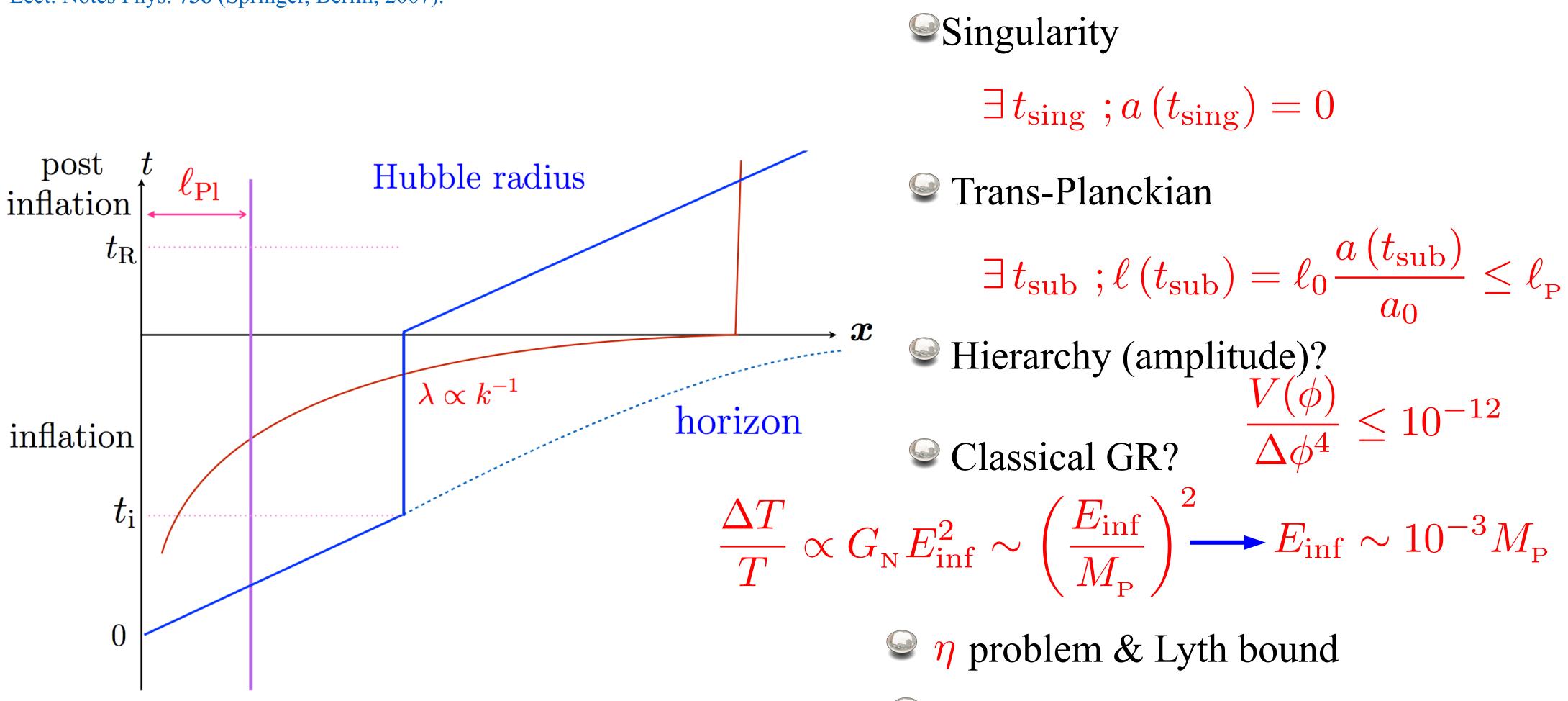
• ... consistent with all known observations

why bother with alternatives?

"When you have eliminated the impossible, whatever remains, however improbable, must be the truth." Sherlock Holmes (The sign of the four):

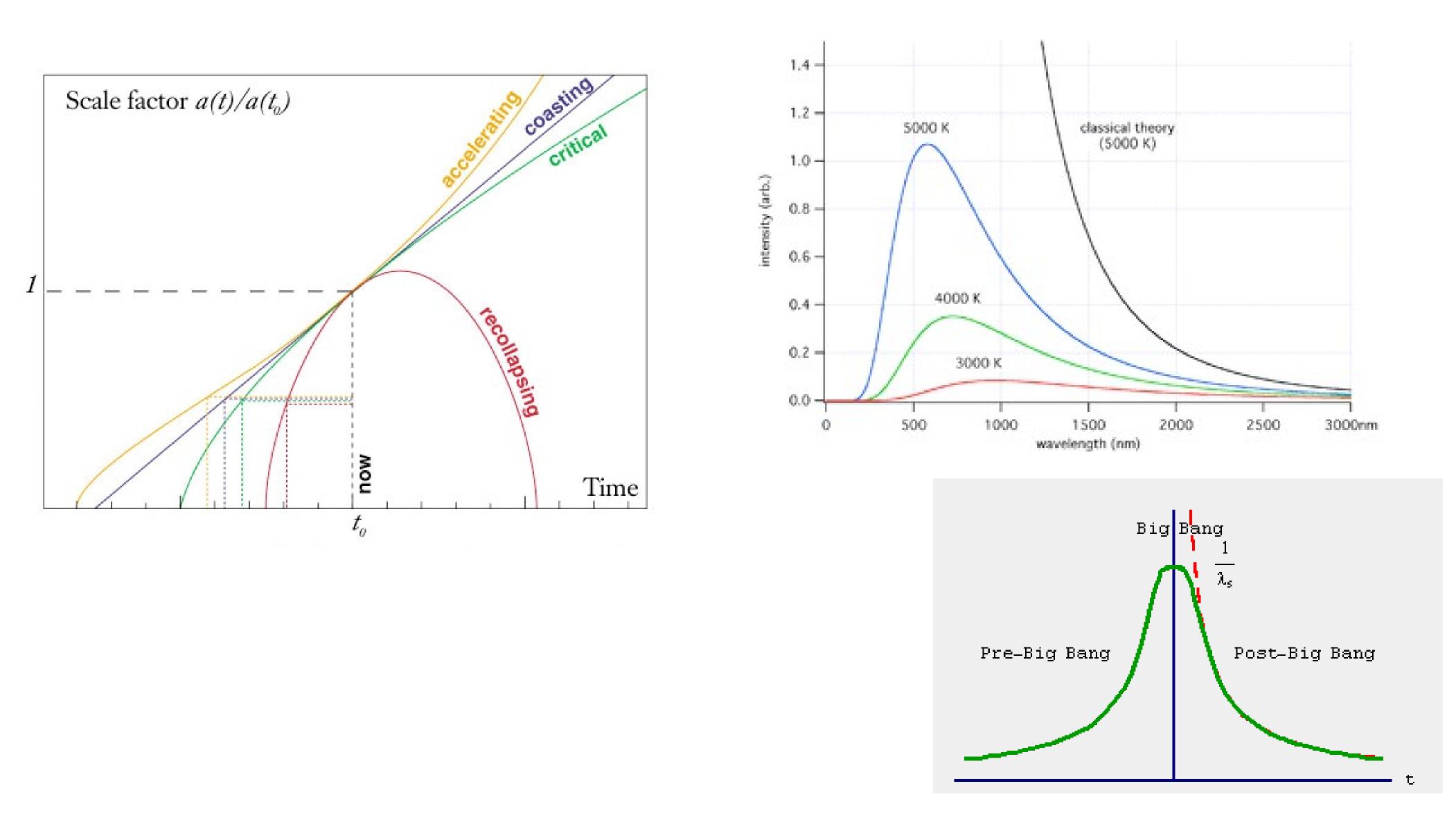
Inflation:

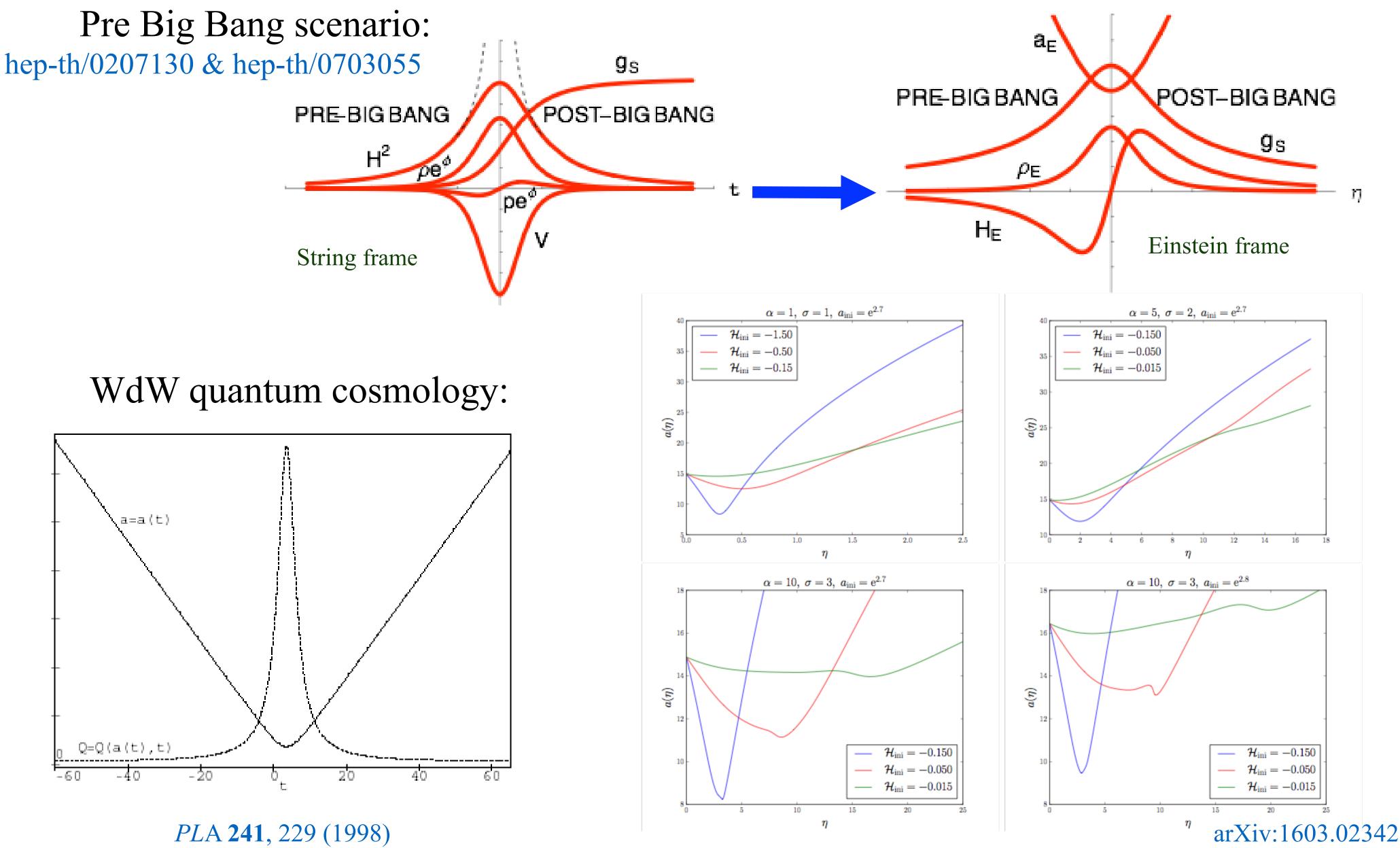
From R. Brandenberger, *in* M. Lemoine, J. Martin & PP (Eds.), "Inflationary cosmology", Lect. Notes Phys. **738** (Springer, Berlin, 2007).



- Initial condition & entropy
- Eternal inflation & measure (anthropic)

Singularity problem Quantum effect?





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A brief history of bouncing cosmology

R. C. Tolman, "On the Theoretical Requirements for a Periodic Behaviour of the Universe", PRD 38, 1758 (1931) G. Lemaître, "L'Univers en expansion", Ann. Soc. Sci. Bruxelles (1933)

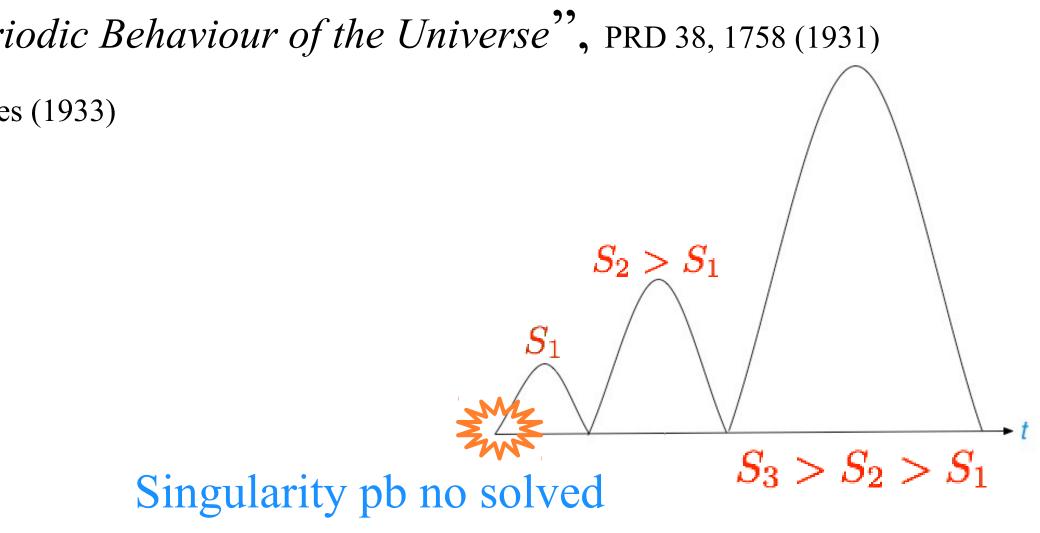
A. A. Starobinsky, "On one non-singular isotropic cosmological model", Sov. Astron. Lett. 4, 82 (1978) V. N. Melnikov, S.V. Orlov, Phys. Lett. A 70, 263 (1979).

R. Durrer & J. Laukerman, "The oscillating Universe: an alternative to inflation", Class. Quantum Grav. 13, 1069 (1996)

• • •

Many new ideas, models...

M. Novello & S.E. Perez Bergliaffa, "Bouncing cosmologies", Phys. Rep. 463, 127 (2008) ----- D. Battefeld & PP, "A Critical Review of Classical Bouncing Cosmologies", Phys. Rep. 571, 1 (2015) R. Brandenberger & PP, "Bouncing cosmologies: Progress and problems", Found. Phys. (2017) Paris - 12th December 2017



Model listing:

Quantum gravity

LQG & LQC Canonical quantum gravity (WdW) String theory

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Non relativistic quantum gravity

Model listing:

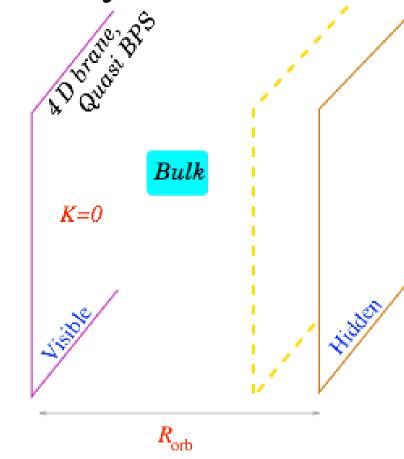
Quantum gravity

Ekpyrotic & cyclic Branes

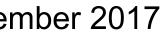
LQG & LQC

Canonical quantum gravity (WdW)

String theory

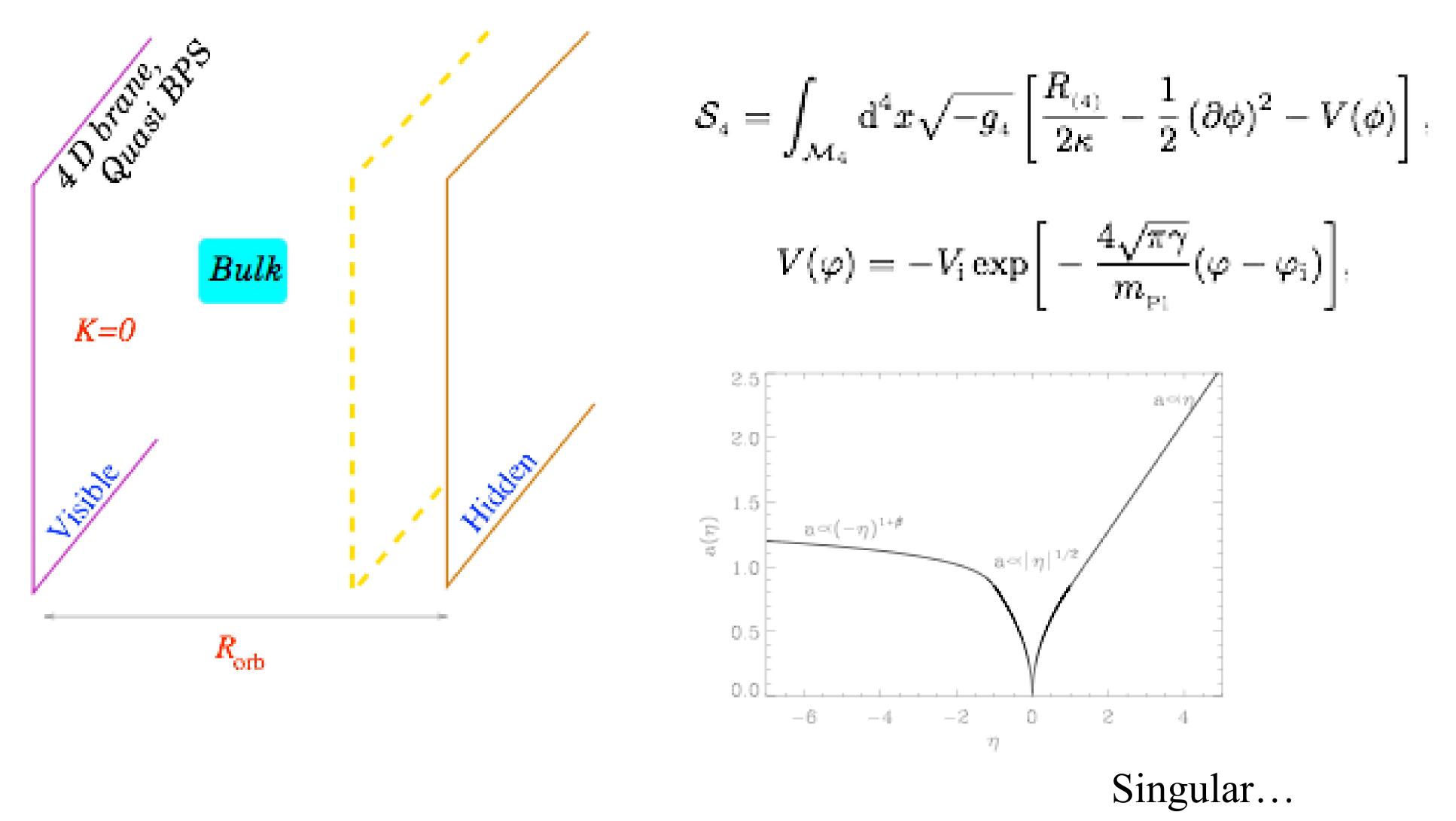


Non relativistic quantum gravity

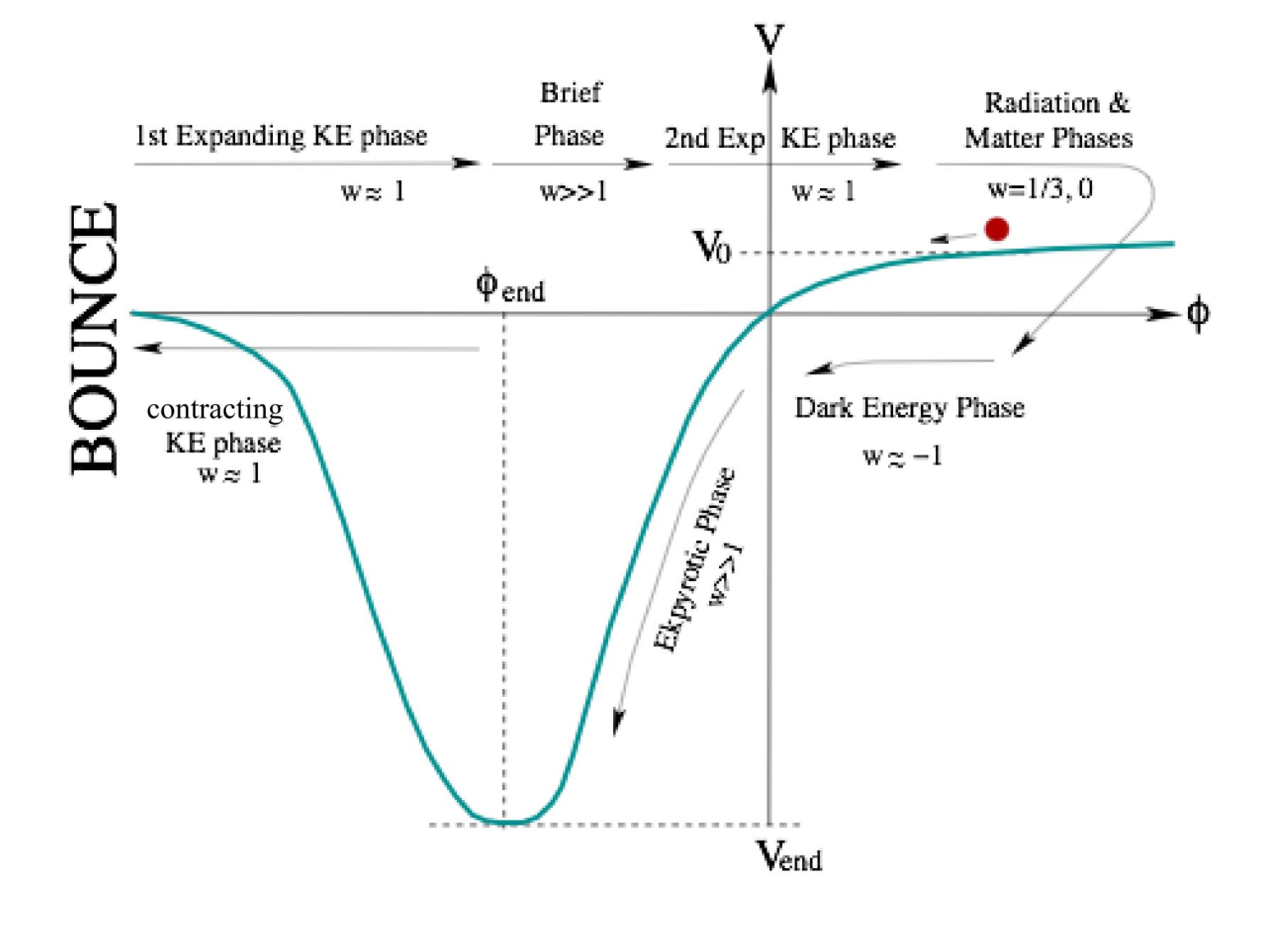


Ekpyrotic scenario:

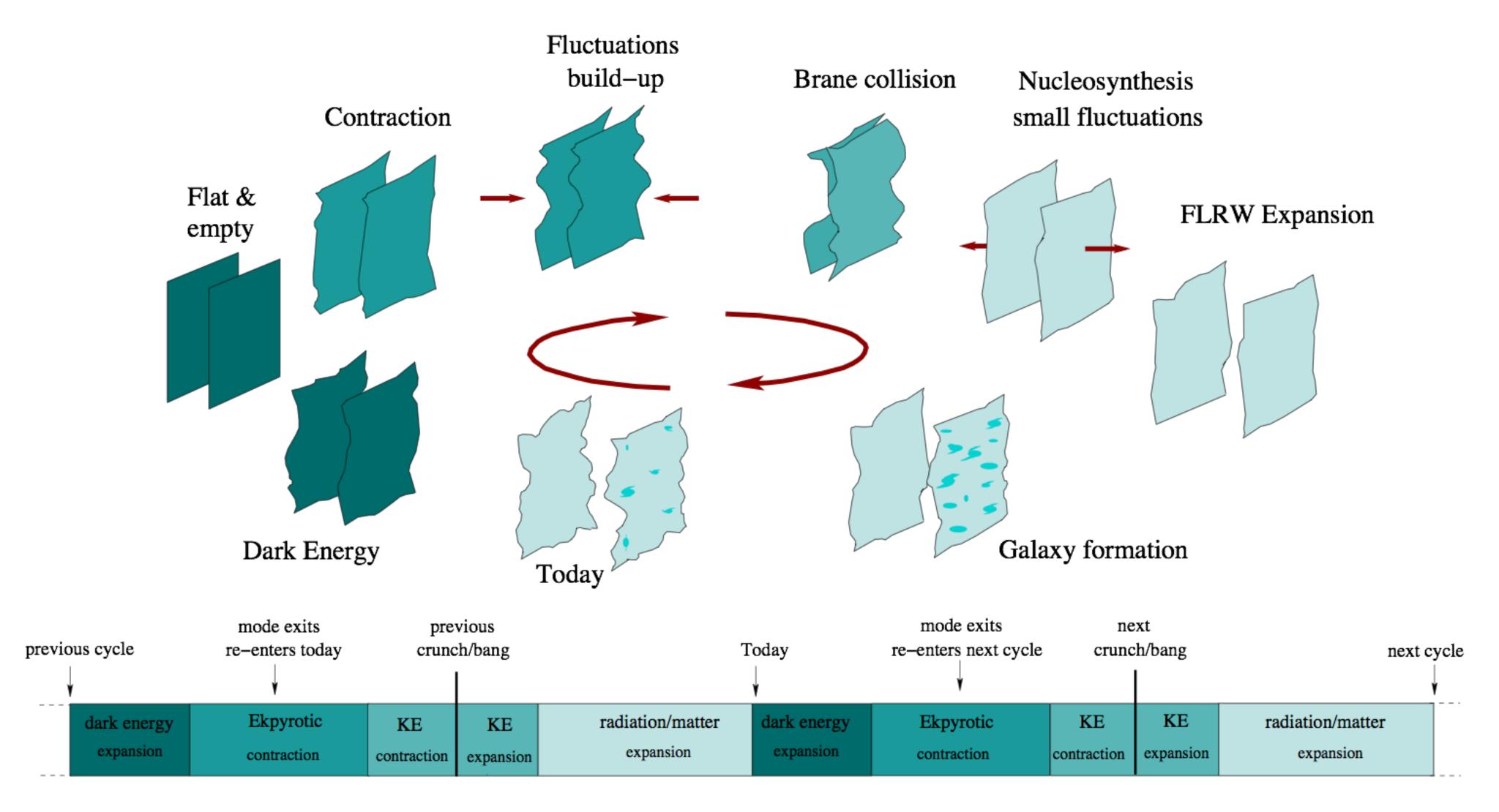




$$imes \int_{\mathcal{M}_5} \mathrm{d}^5 x \sqrt{-g_5} \left[R_{\scriptscriptstyle (5)} - rac{1}{2} \left(\partial arphi
ight)^2 - rac{3}{2} rac{\mathrm{e}^{2 arphi} \mathcal{F}^2}{5 \; !}
ight],$$



Cyclic extension



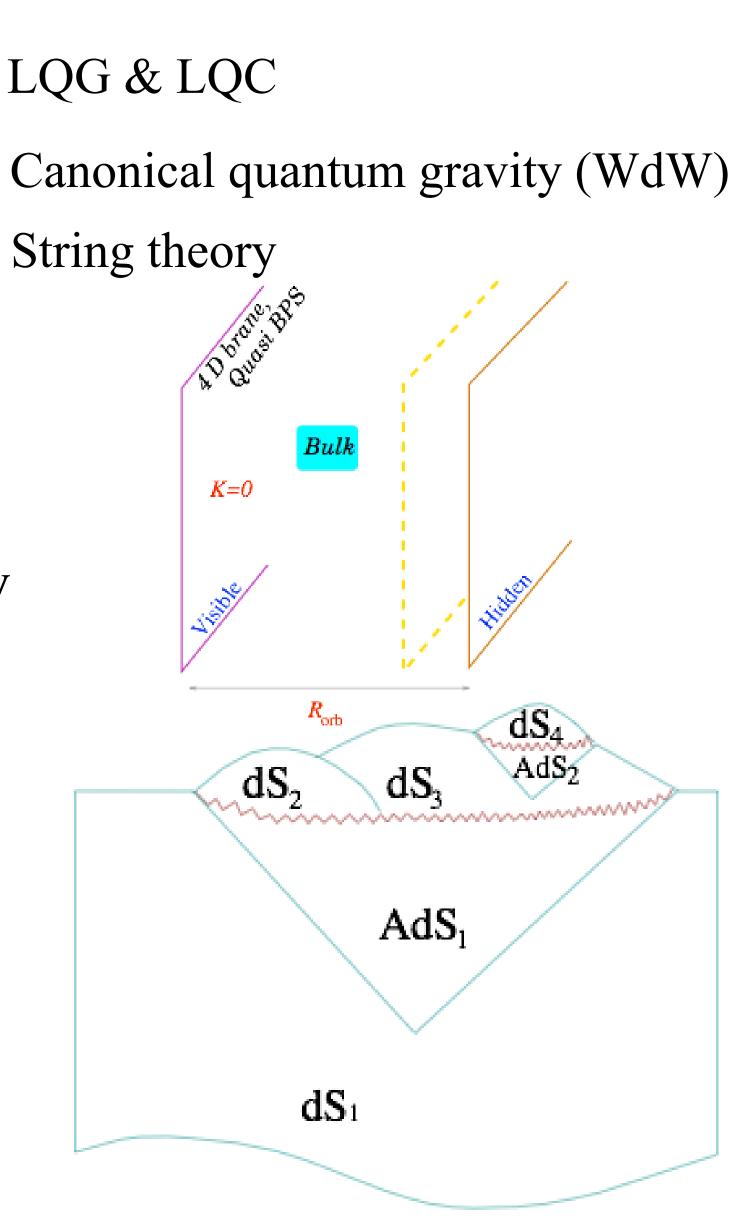
Model listing:

Quantum gravity

Ekpyrotic & cyclic Branes

String gas cosmology Antigravity Galileon Massive gravity

Multiverse models Strings & AdS/CFT



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Non relativistic quantum gravity WdW)

> Horava-Lifshitz Lee-Wick & Quintom F(R), f(T), Gauss-Bonnet

Mimetic matter Non-linear electromagnetic action Spinors & torsion

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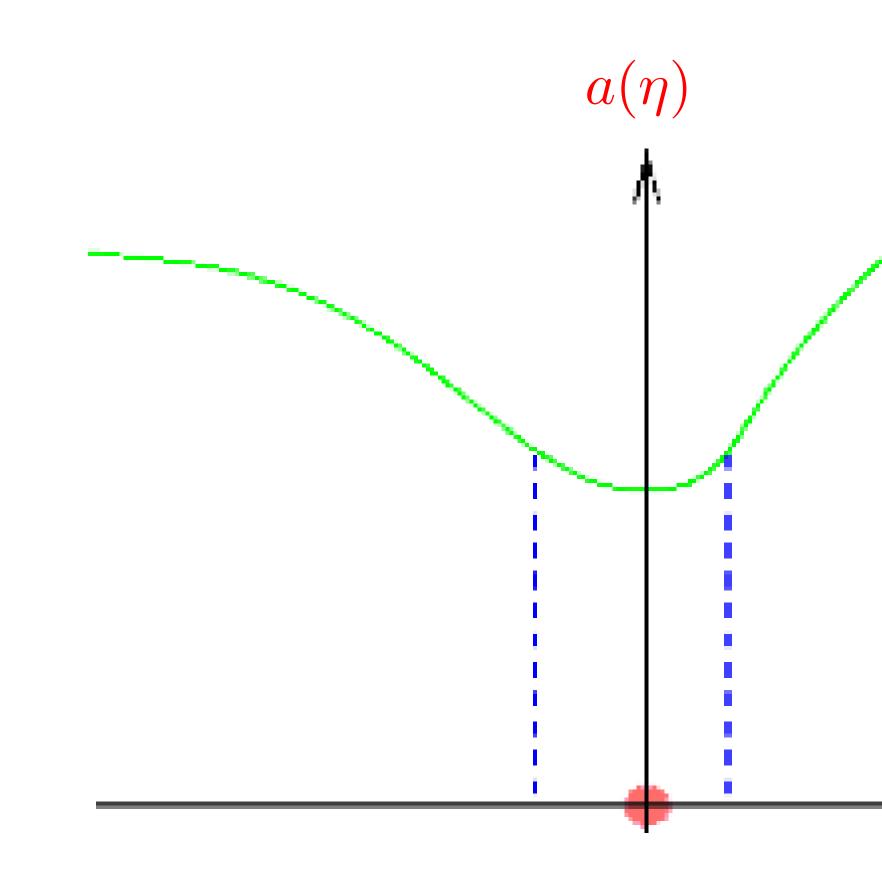
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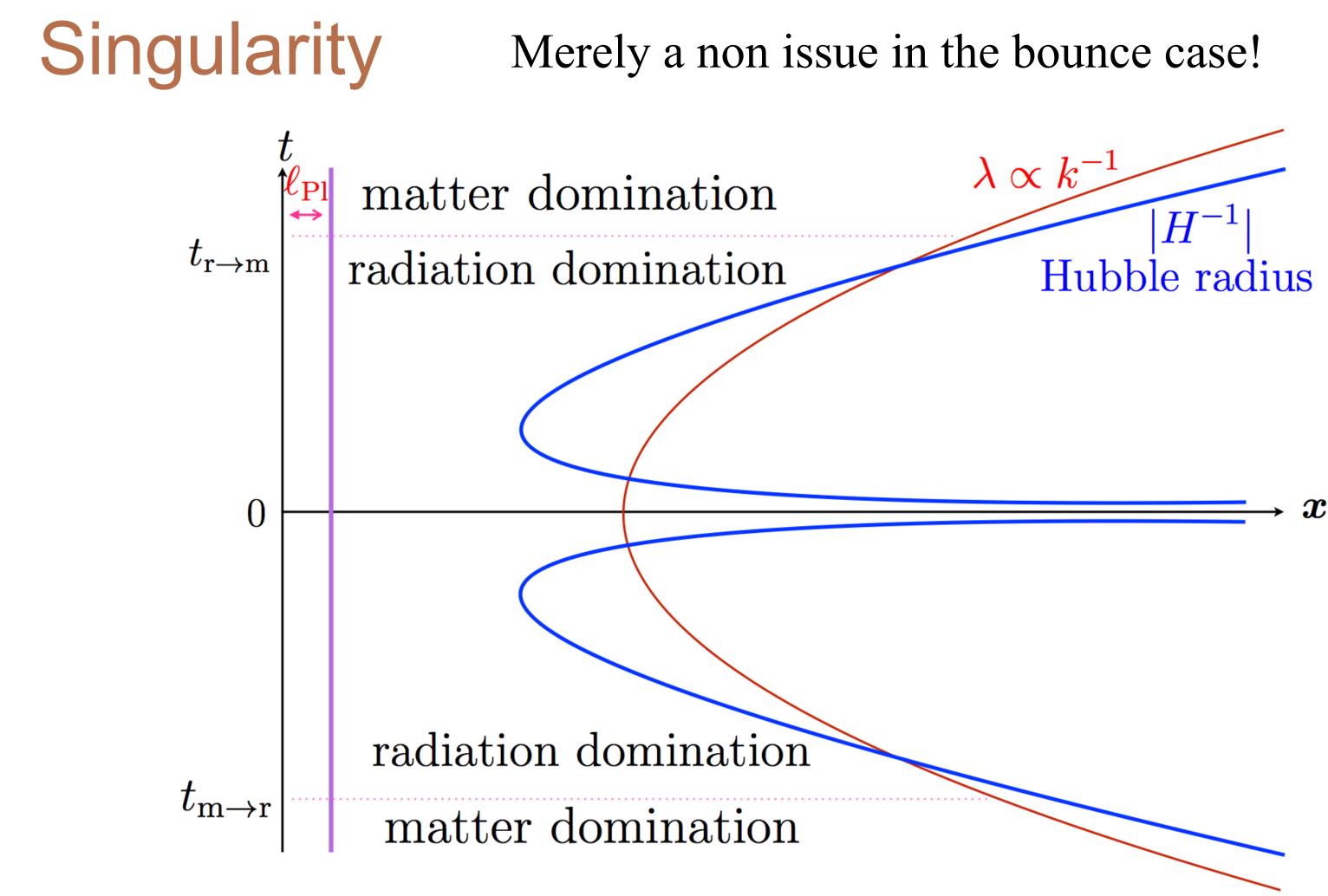
Standard questions and bouncing answers

Singularity Merely a non issue in the bounce case!



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 $\rightarrow \eta$

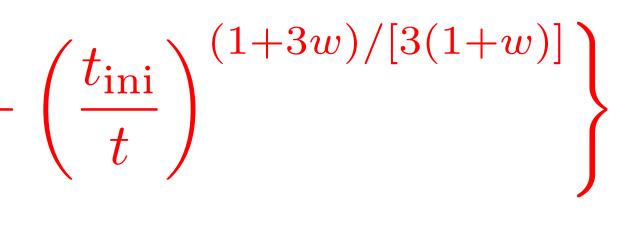


Singularity Merely a non issue in the bounce case! **Horizon** $d_{\rm H} \equiv a(t) \int_{t_{\rm ini}}^{t} \frac{\mathrm{d}\tau}{a(\tau)}$ can be made divergent easily if $t_{\rm ini} \to -\infty$

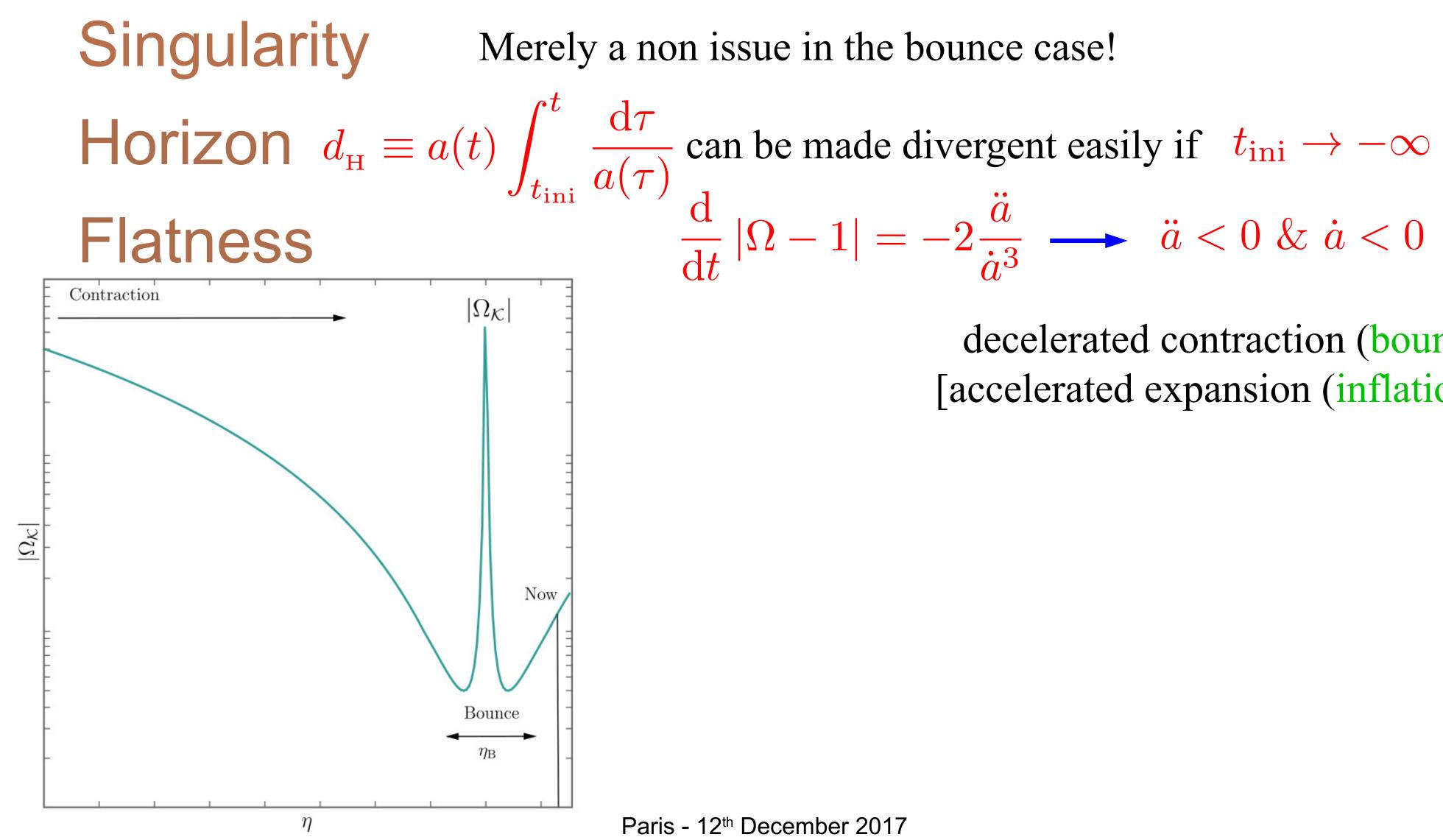
ple:
$$\begin{aligned} d_{\rm H}(t) &= \frac{3(1+w)}{1+3w} |t| \begin{cases} 1-w \\ w &> -\frac{1}{3} \end{cases} \\ w &= t_{\rm ini} \rightarrow -\infty \end{cases} \end{aligned}$$

Examp

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 $=\infty$



decelerated contraction (bounce) [accelerated expansion (inflation)]

Singularity Merely a non issue in the bounce case! **Horizon** $d_{\rm H} \equiv a(t) \int_{t_{\rm ini}}^{t} \frac{\mathrm{d}\tau}{a(\tau)}$ can be made divergent easily if $t_{\rm ini} \to -\infty$ **Flatness** $\frac{\mathrm{d}}{\mathrm{d}t} |\Omega - 1| = -2\frac{\ddot{a}}{\dot{a}^3} \longrightarrow \ddot{a} < 0 \& \dot{a} < 0$

Homogeneity Large & flat Universe + low initial density + diffusion

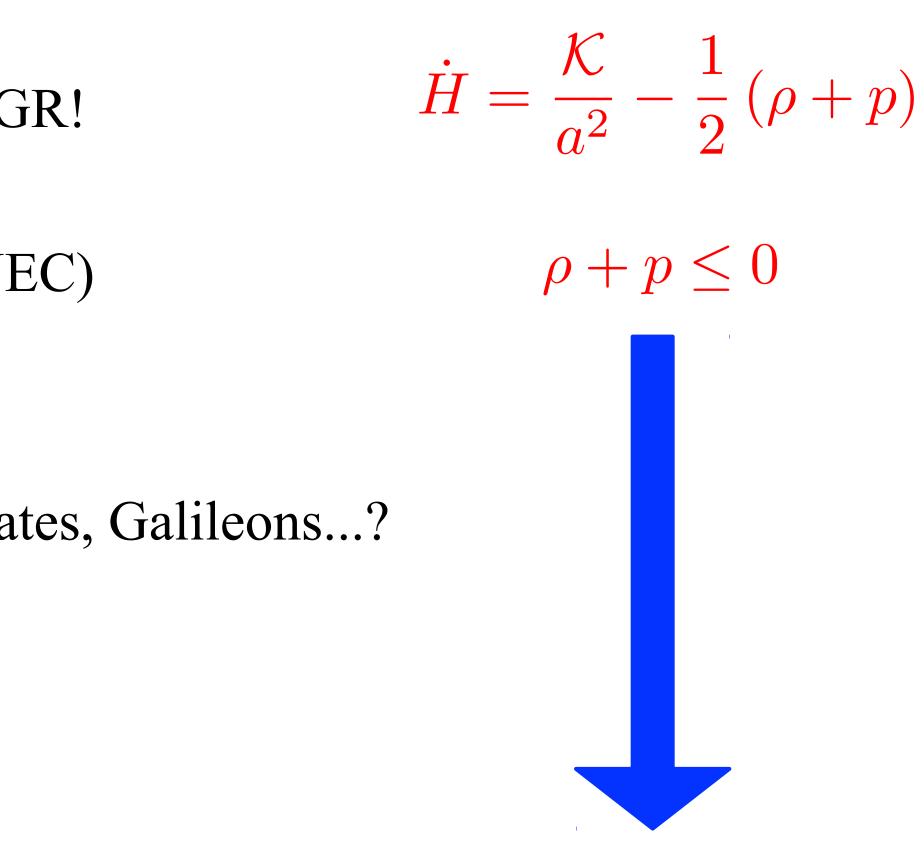
 $t_{\text{dissipation}}(\lambda) \ll t_{\text{Hubble}}$ enough time to dissipate any wavelength

- accelerated expansion (inflation) or decelerated contraction (bounce)
 - \implies quantum vacuum fluctuations...

Implementing a bounce = problem with GR!

Null Energy Condition (NEC)

Many working examples: *K*-bounce, Ghost condensates, Galileons...?



Instabilities

Singularity Merely a non issue in the bounce case! **Horizon** $d_{\rm H} \equiv a(t) \int_{t_{\rm ini}}^{t} \frac{\mathrm{d}\tau}{a(\tau)}$ can be made divergent easily if $t_{\rm ini} \to -\infty$ **Flatness** $\frac{\mathrm{d}}{\mathrm{d}t} |\Omega - 1| = -2\frac{\ddot{a}}{\dot{a}^3} \longrightarrow \ddot{a} < 0 \& \dot{a} < 0$ accelerated expansion (inflation) or decelerated contraction (bounce) **Homogeneity** Large & flat Universe + low initial density + diffusion $t_{\text{dissipation}}(\lambda) \ll t_{\text{Hubble}}$ enough time to dissipate any wavelength \implies quantum vacuum fluctuations... Isotropy Potentially problematic: model dependent

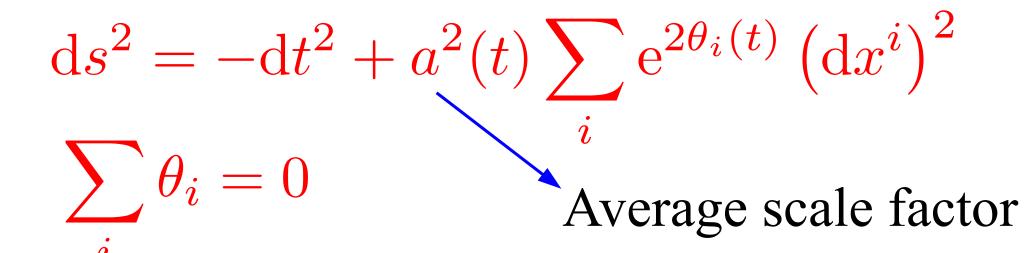
The problem with contraction: BKL/shear instability Ricci flat Bianchi I anisotropic universe

Friedman equations

$$H^{2} = \frac{\rho}{3M_{P}^{2}} + \left(\frac{1}{6}\sum_{i}\dot{\theta}_{i}^{2}\right)$$

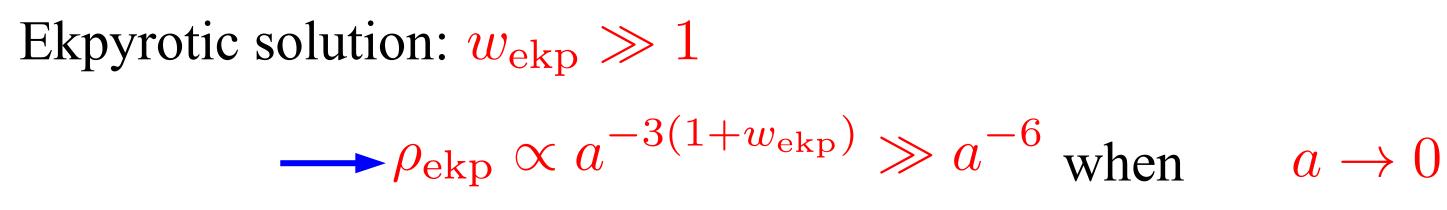
$$\ddot{\theta}_{i} + 3H(\dot{\theta}_{i} + 3H(\dot{\theta}_{$$

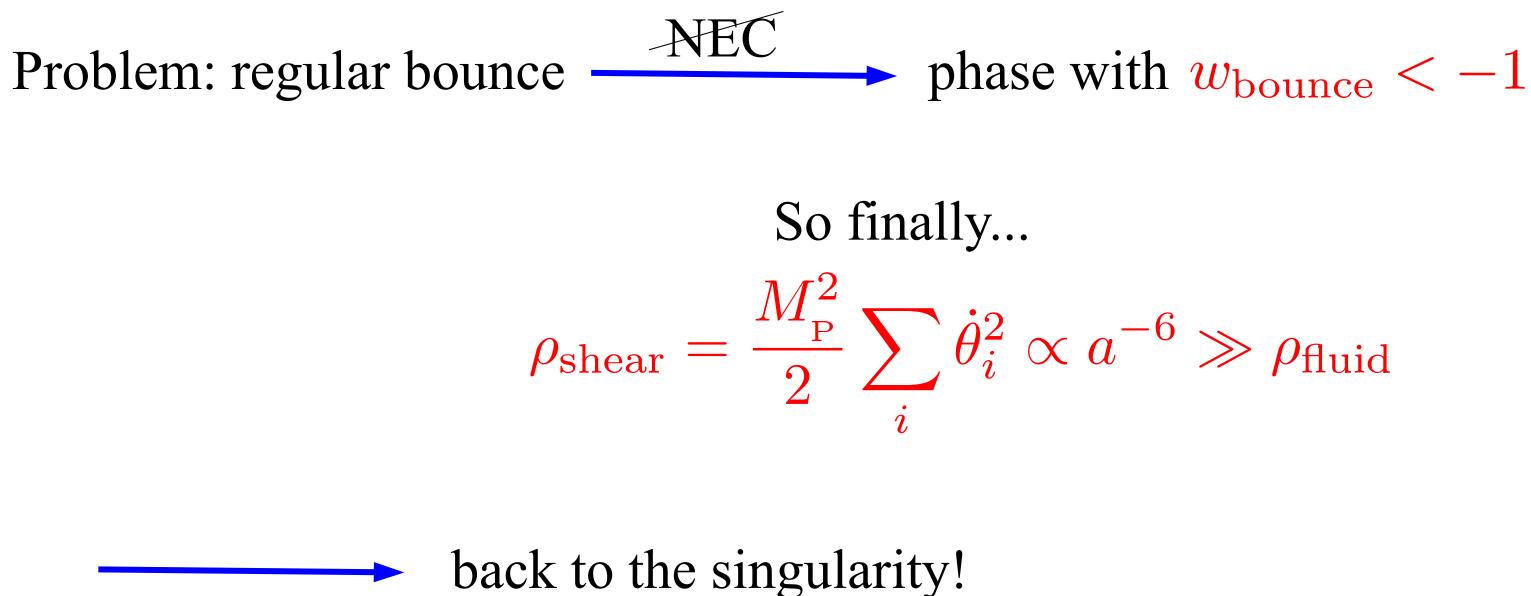
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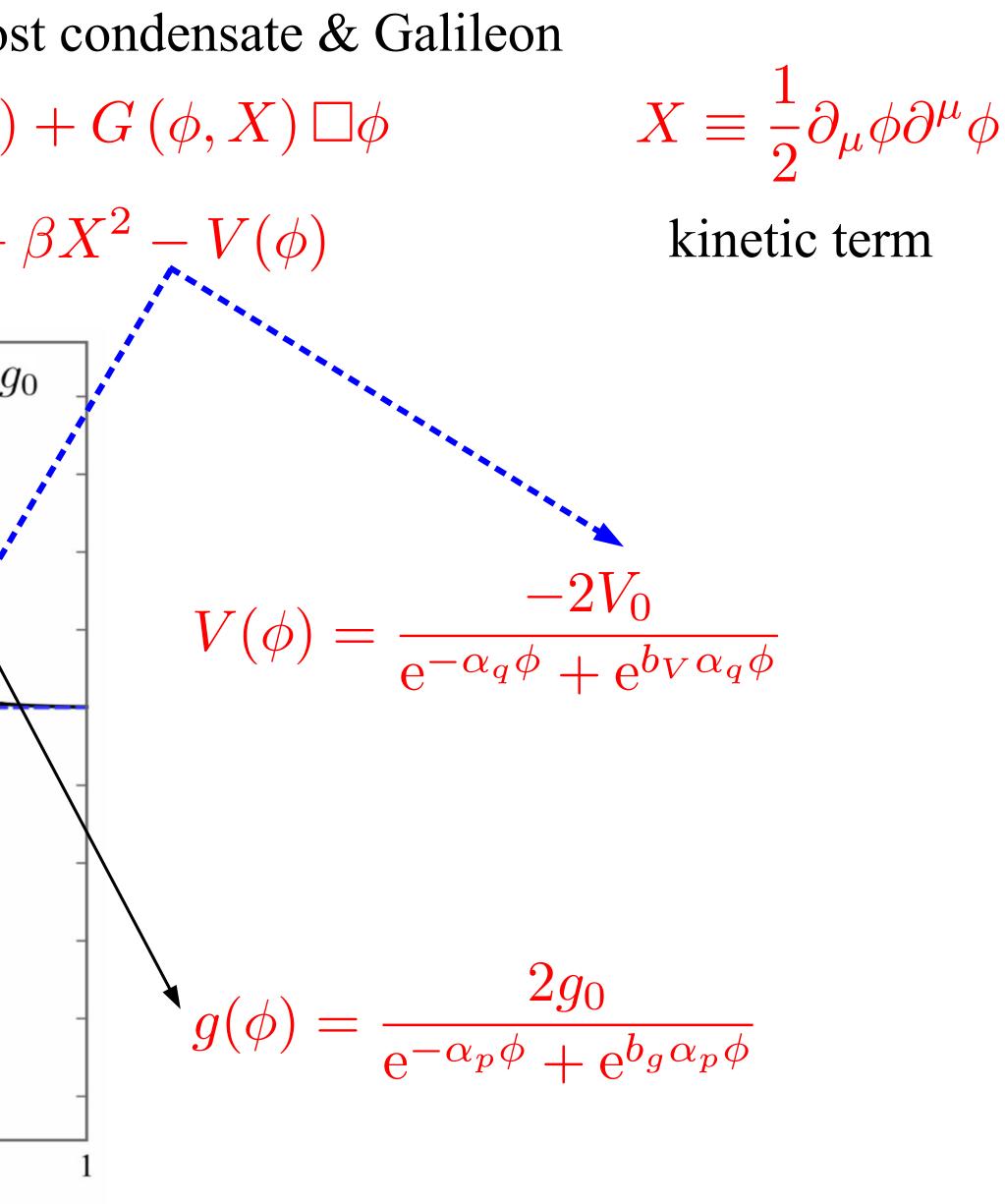
n Hubble parameter

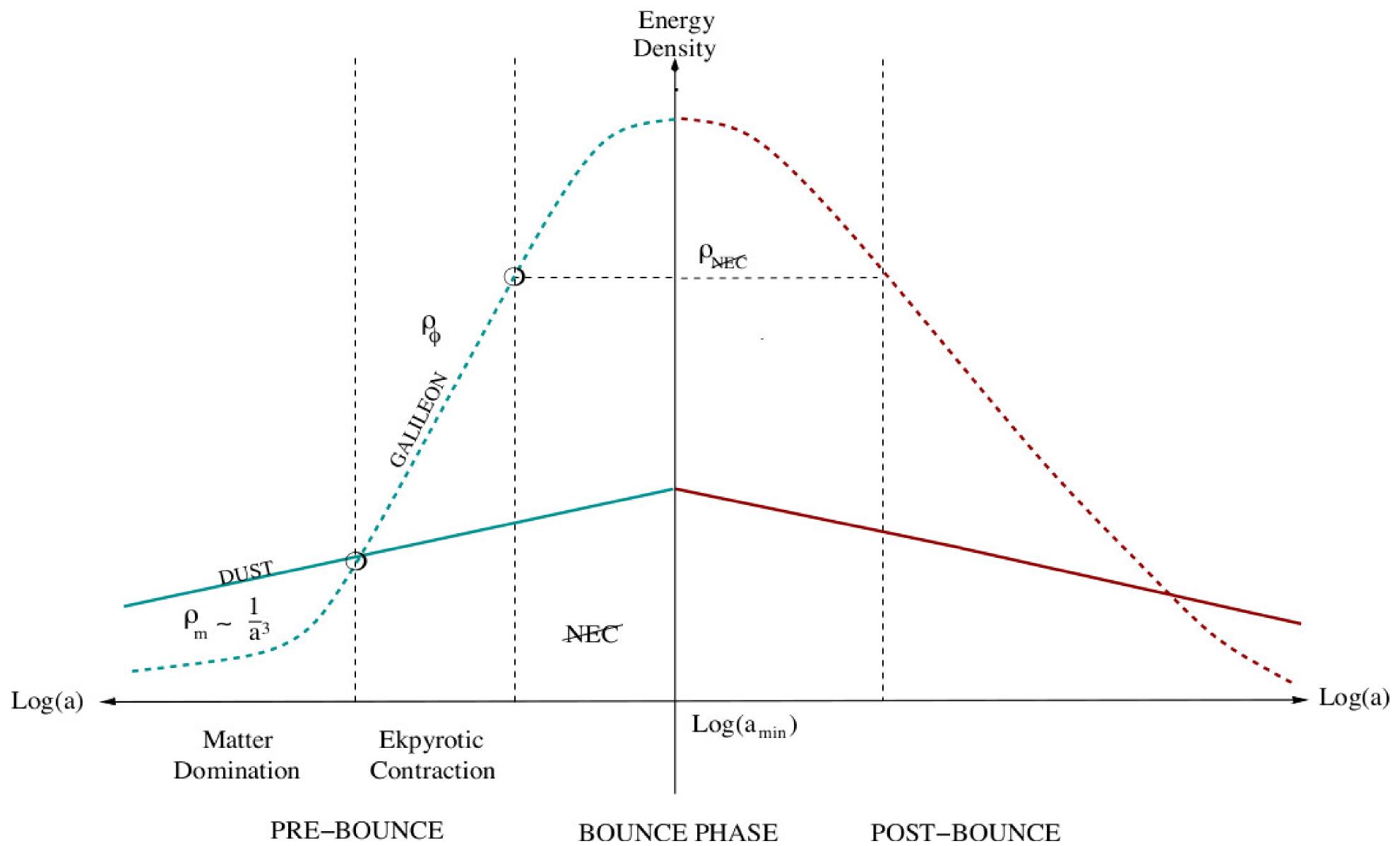
 $\dot{\theta}_i = 0 \implies \dot{\theta}_i \propto a^{-3}$ $ho_{
m shear} \propto a^{-6}$





A nonsingular bounce model: ghost condensate & Galileon $\mathcal{L}\left[\phi\left(x\right)\right] = K\left(\phi, X\right) + G\left(\phi, X\right) \Box\phi$ Specific choices: $K(\phi, X) = M_{\rm P}^2 \left[1 - g(\phi)\right] X + \beta X^2 - V(\phi)$ $G(X) = \gamma X$ $g(\phi)/g_0$ 1.00 $g(\phi)$ 0.75 0.50 10-28 $V(\phi)$ 0.25 0.00-0.25 -0.50 -0.75 -1.00 $V(\phi)/V_0$ -1.25 -1 0 Φ

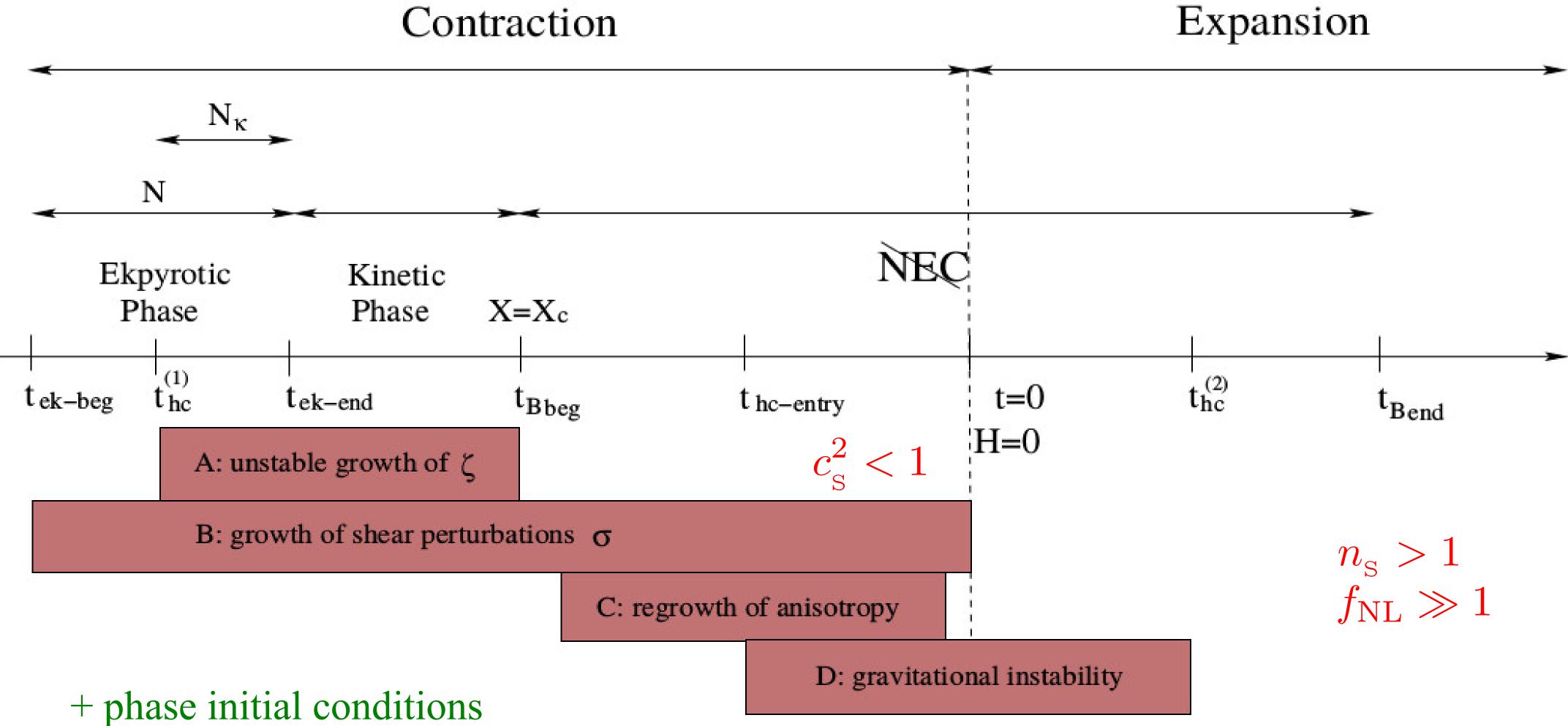




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POST-BOUNCE

Summary of possible problems



Singularity Merely a non issue in the bounce case! **Horizon** $d_{\rm H} \equiv a(t) \int_{t_{\rm ini}}^{t} \frac{\mathrm{d}\tau}{a(\tau)}$ can be made divergent easily if $t_{\rm ini} \to -\infty$ **Flatness** $\frac{\mathrm{d}}{\mathrm{d}t} |\Omega - 1| = -2\frac{\ddot{a}}{\dot{a}^3} \longrightarrow \ddot{a} < 0 \& \dot{a} < 0$ accelerated expansion (inflation) or decelerated contraction (bounce) **Homogeneity** Large & flat Universe + low initial density + diffusion $t_{\text{dissipation}}(\lambda) \ll t_{\text{Hubble}}$ enough time to dissipate any wavelength \implies quantum vacuum fluctuations... Isotropy Potentially problematic: model dependent **Others** dark matter/energy, baryogenesis, ...