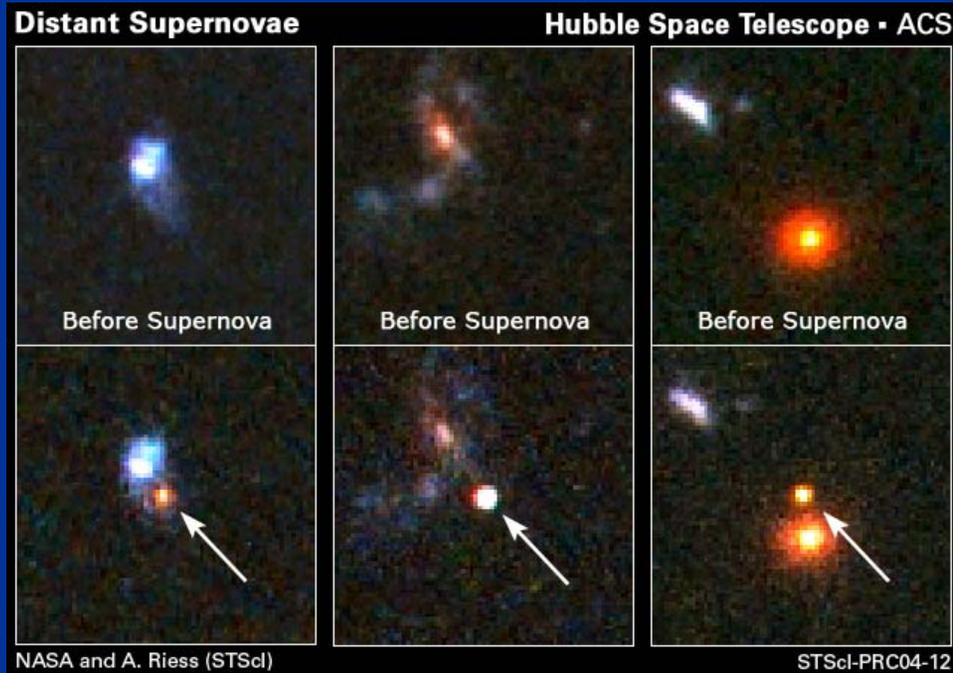


Learning to Love the Scatter in Type Ia SuperNovae

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with Scott Dodelson (Fermilab)
Phys. Rev. D. 74:063515,2006

Type Ia SuperNovae

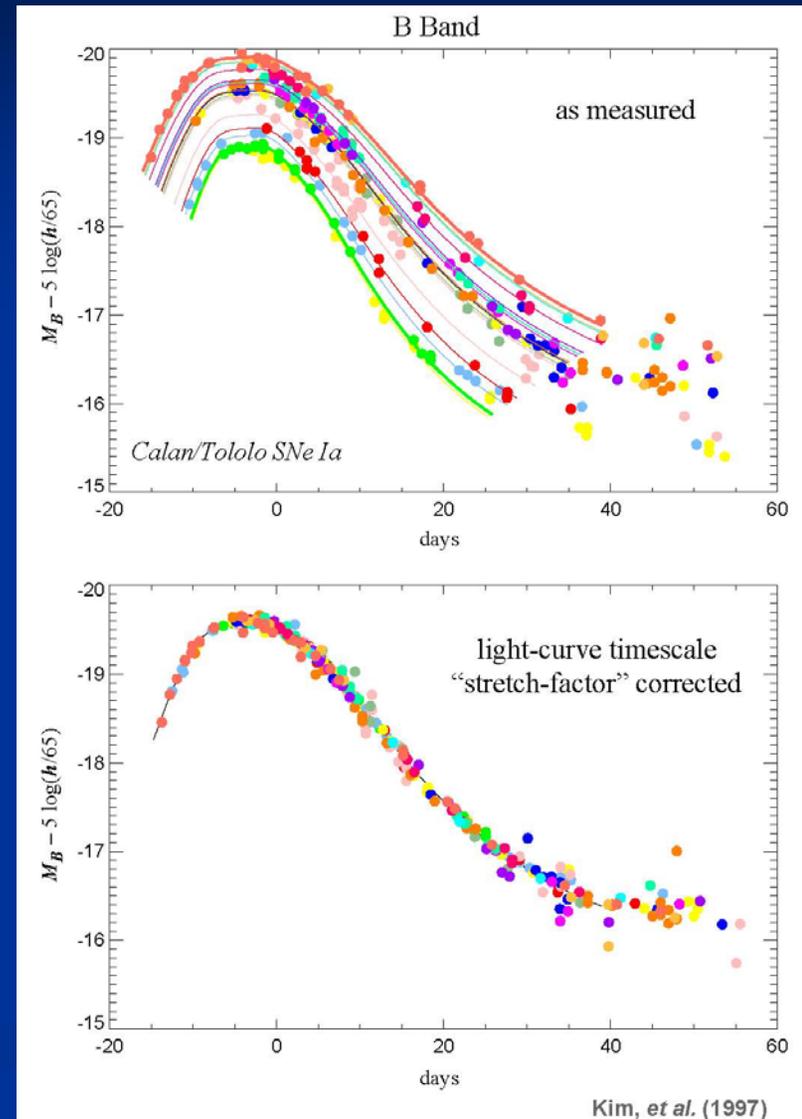
High redshift SNIa are detected through image subtraction



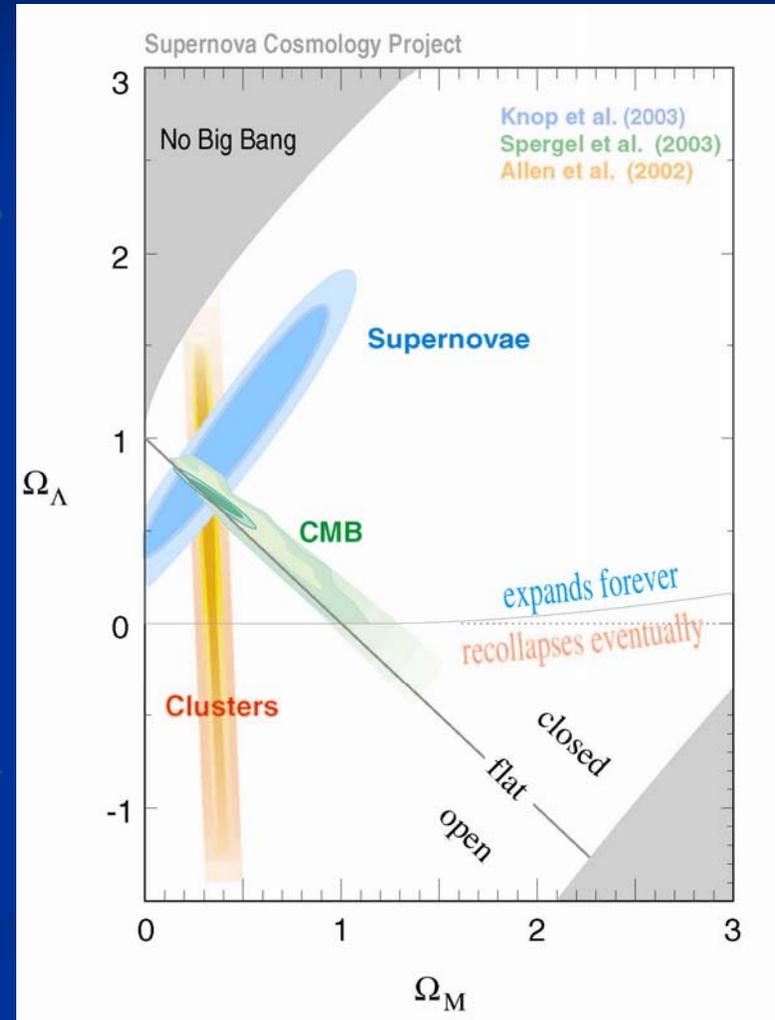
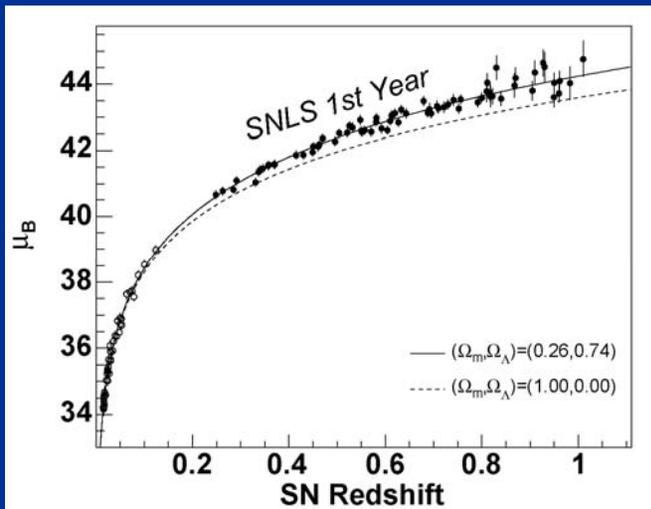
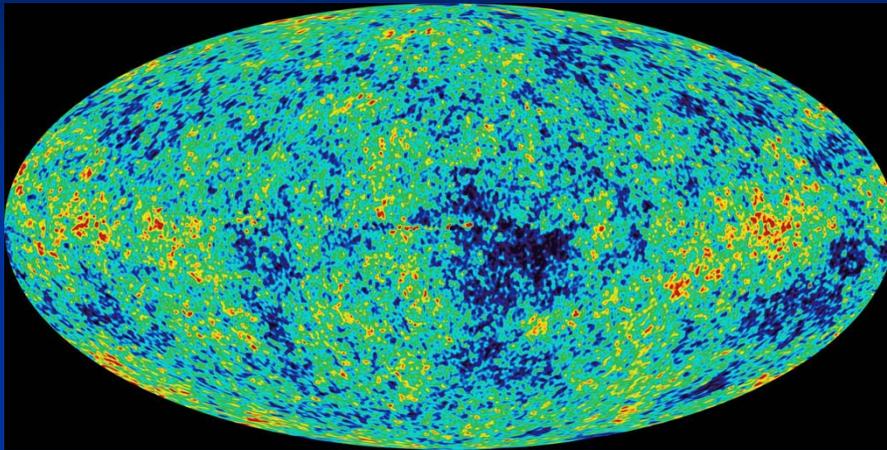
SN1994D imaged with HST.
High-Z SN Search Team

Type Ia SuperNovae as Standard Candles

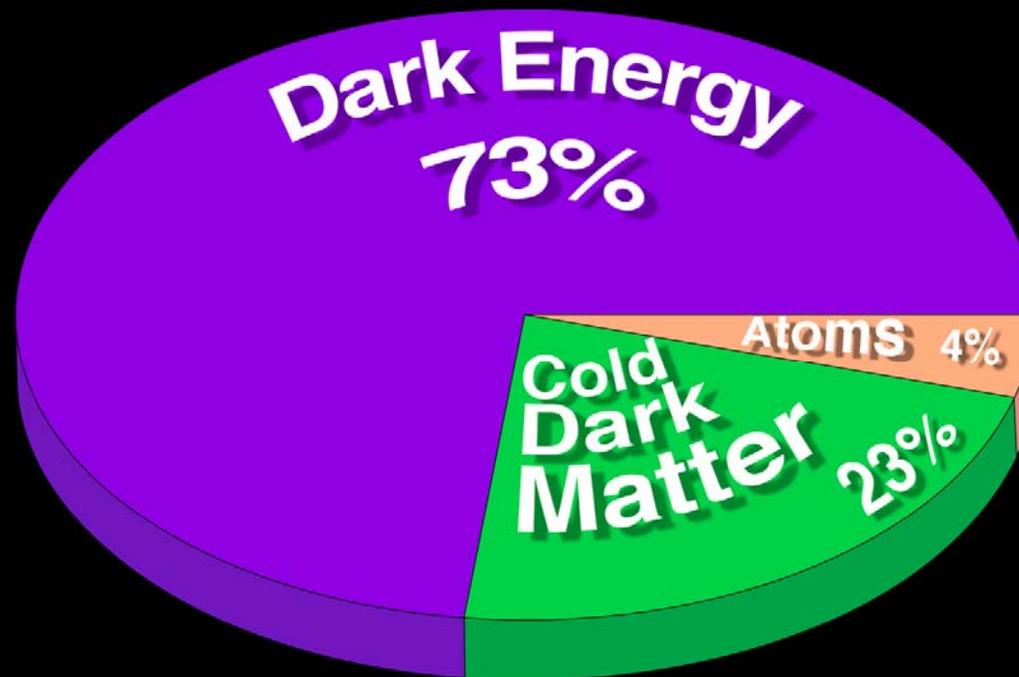
- Type Ia SuperNovae have different – but self-similar – light curves.
- They are standardizable candles...
- ... once the “stretch factor” of their light curves is taken into account.



Cosmology and the Concordance Model



Cosmic Pie



Gravitational Lensing of SNIa

- The (lensing-induced) scatter in SNIa luminosity distance is usually seen as a nuisance
- Advantages of using SNIa to do lensing:
 - Information that comes *for free!*
The information about gravitational lensing is buried in the scatter of $d_L(z)$.
It can be mined out of SNIa searches data at no extra (experimental) cost.
 - SNIa are *point sources*.
Gravitational lensing will increase or decrease their luminosity but no prior knowledge of their shape is required: they're point-like!
- Disadvantages:
 - There are very few of them

Distance Modulus

- Measured distance modulus μ of a SNIa is made of three contributions

$$\mu = \mu_0 + \delta\mu_{\text{int}} + \delta\mu_{\text{cos}}$$

where

- μ_0 is the background distance modulus
- $\delta\mu_{\text{int}}$ is the intrinsic scatter in SNIa luminosity
- $\delta\mu_{\text{cos}}$ is the scatter in SNIa luminosity due to *weak lensing*
- Goal: can we mine the information hidden in the SNIa scatter?

SN Ia Scatter due to Gravitational Lensing $\delta\mu_{\text{cos}}$

$$\mu = \mu_0 + \delta\mu_{\text{int}} + \delta\mu_{\text{cos}}$$

- The variance of $\delta\mu_{\text{cos}}$ can be related to the variance of the lensing convergence \mathbf{K}

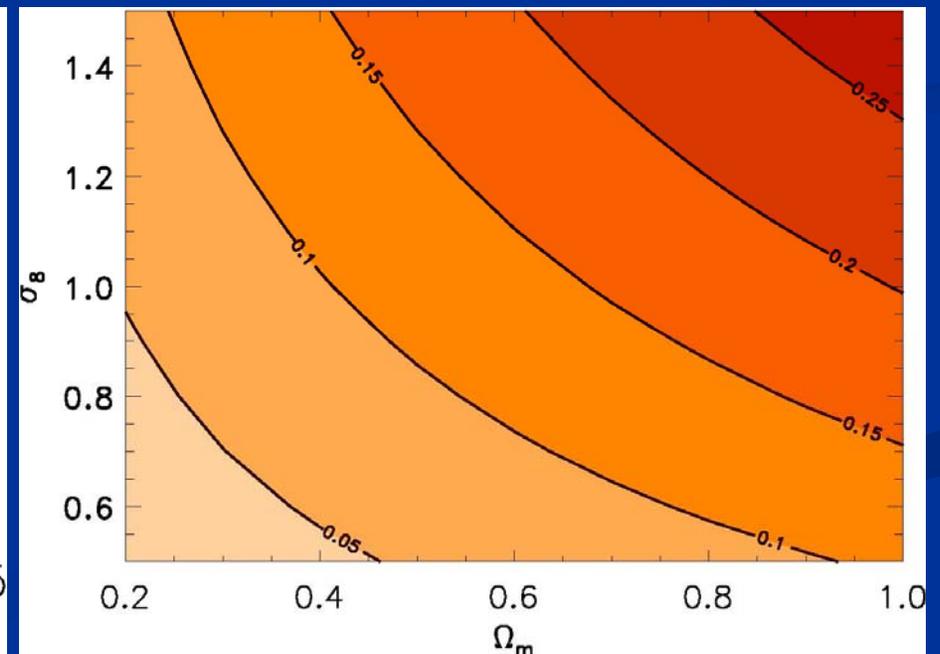
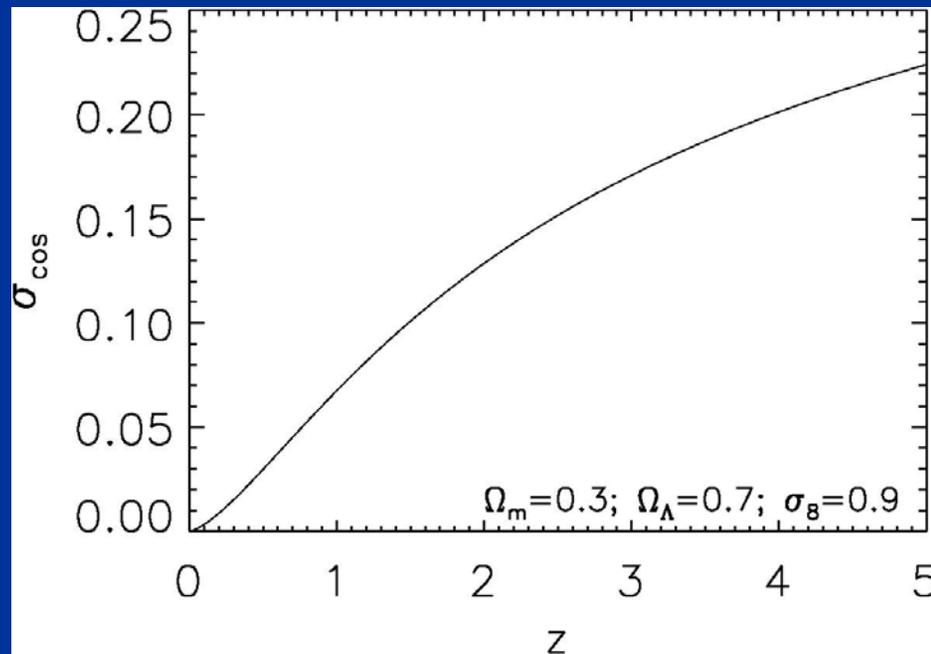
$$\sigma_{\text{cos}}^2 = \left[\frac{5}{\ln(10)} \right]^2 \langle \mathbf{K}^2 \rangle = \frac{225\pi\Omega_m^2 H_0^4}{4[\ln(10)]^2} \int_0^{\chi_s} d\chi [1+z(\chi)]^2 \frac{\chi^2(\chi_s - \chi)^2}{\chi_s^2} \int_0^\infty \frac{dk}{k^2} \Delta^2[k, z(\chi)]$$

Frieman [1996], Hamana and Futamase [1999]

- σ_{cos}^2 depends on the power spectrum $\Delta^2(\mathbf{k}, z)$, which in turn depends on Ω_m and σ_8 .
- Assume Λ CDM cosmology and use the algorithm of Smith et al. [2002] to evaluate the dimensionless power spectrum $\Delta^2(\mathbf{k}, z)$

Calculation of σ_{cos}

- Given a procedure to evaluate $\Delta^2(k,z)$, σ_{cos} can be evaluated numerically and used to infer the value of the cosmological parameters.
- Intuitively makes sense: more matter or more clumpiness lead to more gravitational lensing and to more scatter.

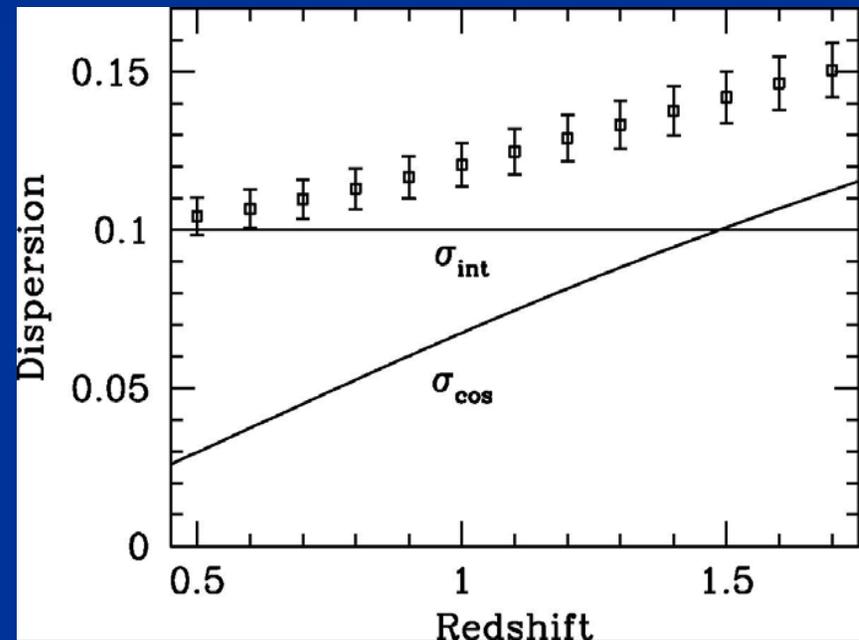


From σ_{cos} to the likelihood

- Recall that

$$\mu = \mu_0 + \delta\mu_{\text{int}} + \delta\mu_{\text{cos}}$$

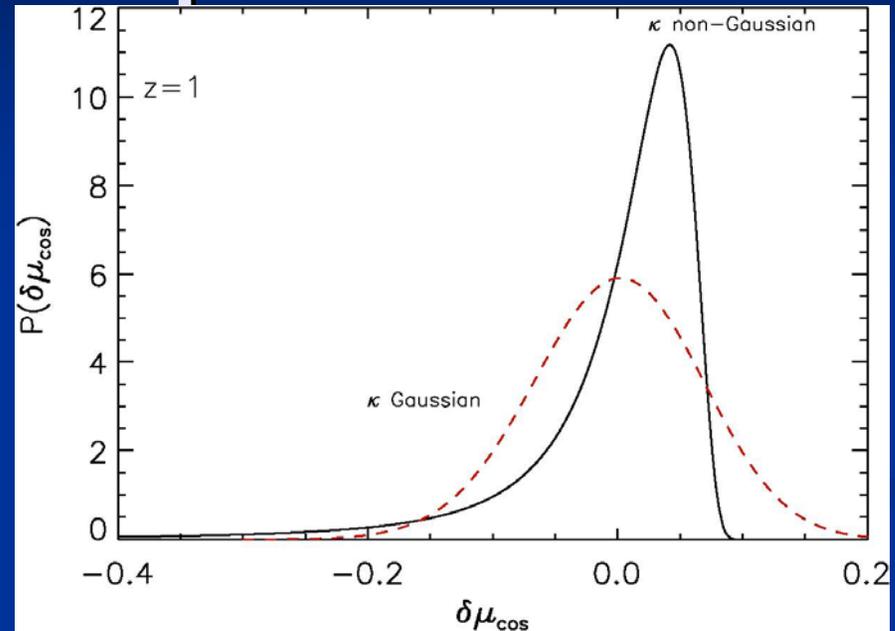
- Information about σ_{cos} is not straightforward to extract from the total scatter because of the *unknown* intrinsic scatter component σ_{int} .
- The best way to proceed is to perform a MonteCarlo simulation of a SNAP-like experiment and then a **likelihood analysis**.



Montecarlo Simulation: the need for pdfs

$$\mu = \mu_0 + \delta\mu_{\text{int}} + \delta\mu_{\text{cos}}$$

- Need to know the full pdf of $\delta\mu_{\text{cos}}$ (and not just its variance) for any given SNIa.
 - **Simplest thing:** assuming that $\delta\mu_{\text{cos}}$ is gaussian distributed.
This is not ok because the actual pdf is known to be skewed.
Wambsgness et al. [1997], Holz and Wald [1998]
 - **Best option:** using the results of Valageas [1999,2000] and Wang, Holz and Munshi [2002], the correct expression for the distribution for $\delta\mu_{\text{cos}}$ can actually be derived.

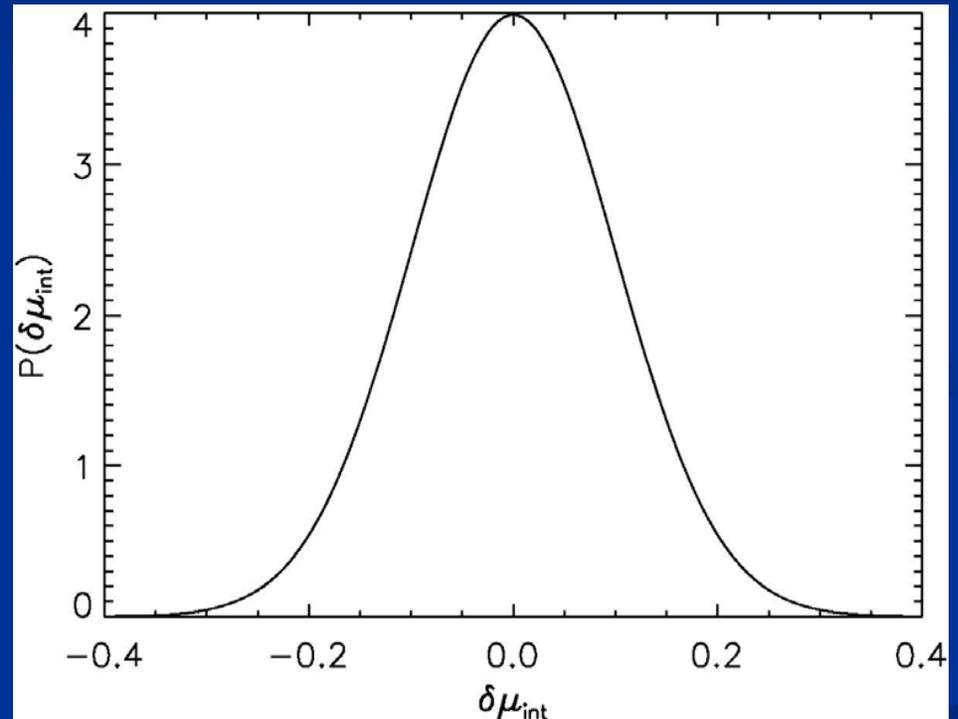


The skewness of the distribution for $\delta\mu_{\text{cos}}$:
Most of the Universe is empty and therefore
⇒ Most SNIa are slightly demagnified and only a tiny fraction is highly magnified
⇒ Most SNIa will appear fainter (and farther) and only a few will appear brighter (and closer)
⇒ Most SNIa will have their modulus slightly increased and only a few will have it decreased.

SNIa Intrinsic Scatter $\delta\mu_{\text{int}}$

$$\mu = \mu_0 + \delta\mu_{\text{int}} + \delta\mu_{\text{cos}}$$

- **Assumption:** $\delta\mu_{\text{int}}$ is Gaussian distributed with dispersion $\sigma_{\text{int}}=0.1$ and $\langle\delta\mu_{\text{int}}\rangle=0$
- **Assumption:** $\delta\mu_{\text{int}}$ is redshift *independent* (no evolution)



Background Distance Modulus μ_0

$$\mu = \mu_0 + \delta\mu_{\text{int}} + \delta\mu_{\text{cos}}$$

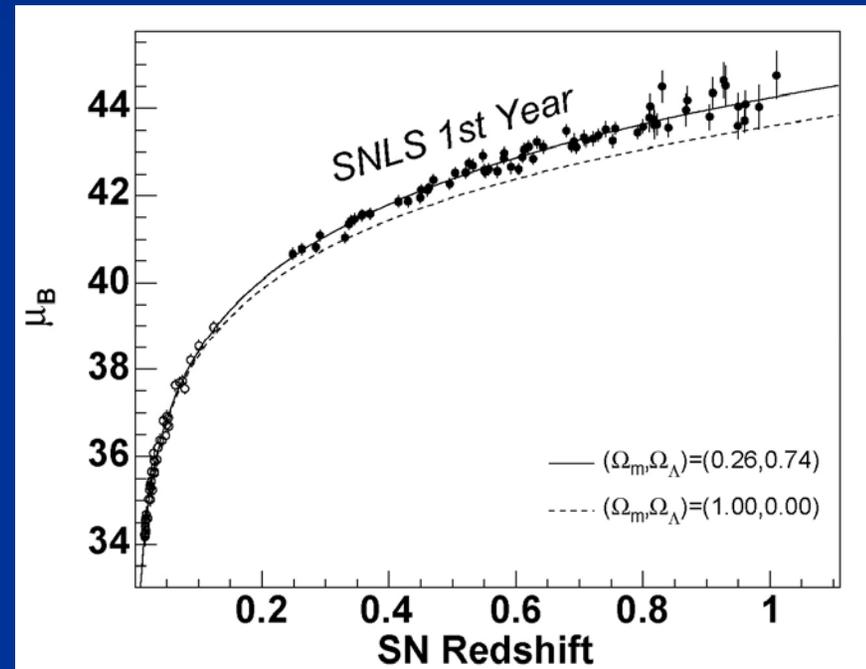
- μ_0 depends on $(\Omega_m, \Omega_\Lambda)$ through

$$\mu_0 = 5 \text{Log} \left[\frac{d_L(z)}{10 \text{ pc}} \right]$$

$$d_L(z) = (1+z)c \int_0^z \frac{dz'}{H(z')}$$

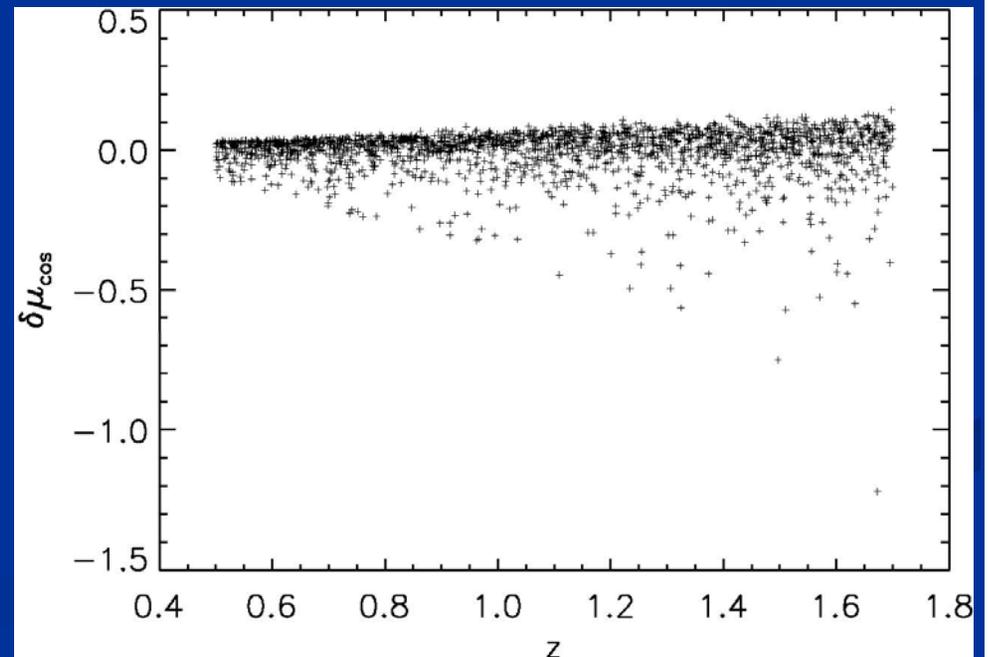
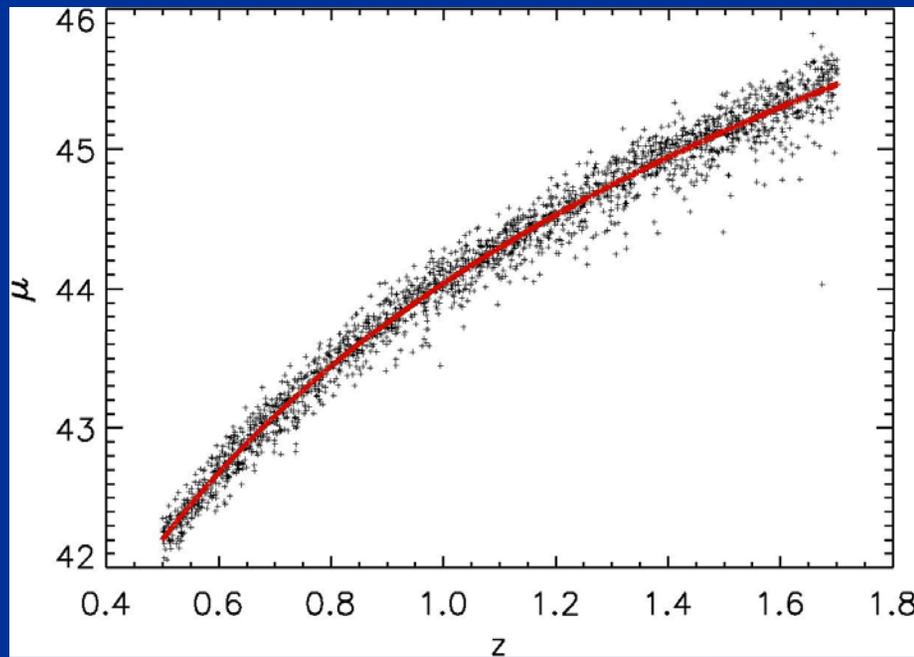
- **Assumption:** flat geometry

$$\Omega_m + \Omega_\Lambda = 1$$



Simulated SNAP-like experiment

- Once the pdf for the two contributions $\delta\mu_{\text{int}}$ and $\delta\mu_{\text{cos}}$ are available, we generated a synthetic sample of 2000 SNIa in the redshift range $z=[0.5,1.7]$.

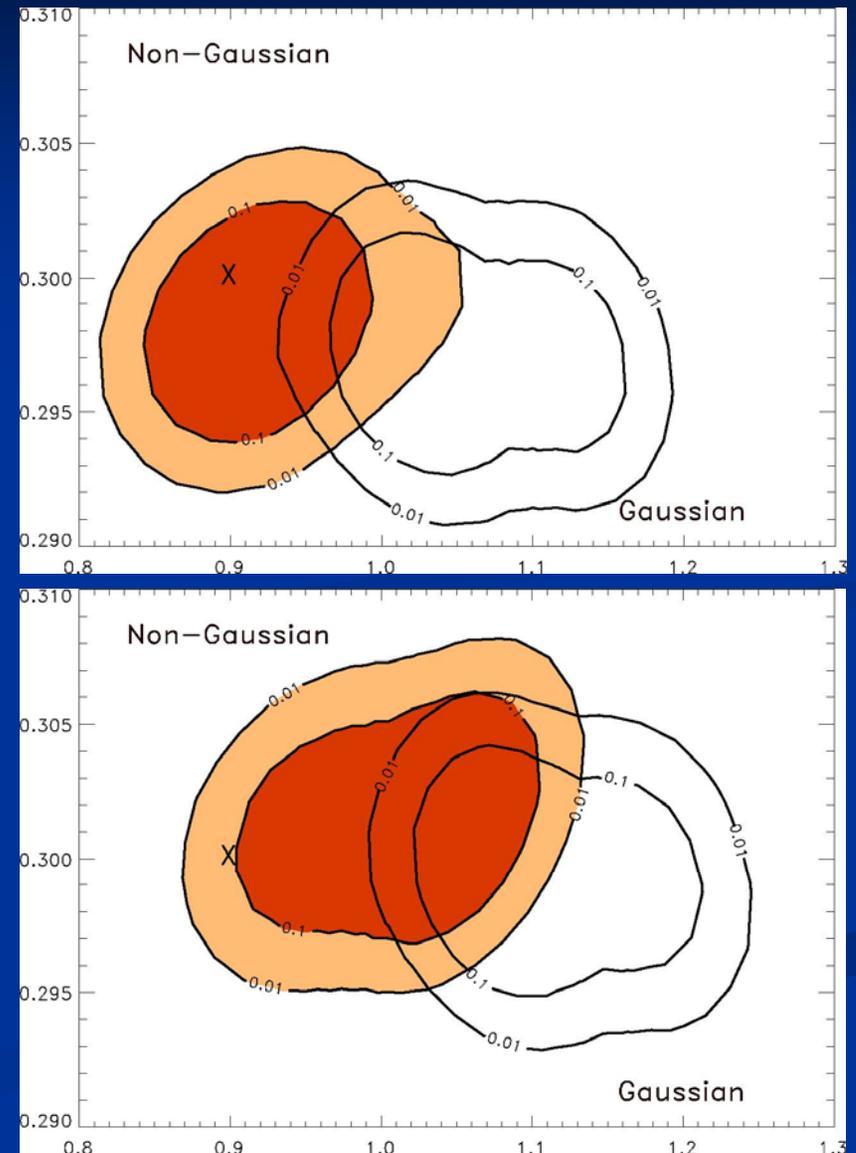


Likelihood Analysis: 2 pdfs

- The data generated are then analyzed using for the distribution of $\delta\mu_{\text{cos}}$
 - the fiducial pdf derived from Wang, Holz and Munshi [2002] and used to generate the data.
 - the gaussian pdf. This is done to determine whether a cosmological parameter extraction performed neglecting the non-gaussianity of the pdf would lead to biased estimates.
- We marginalize over the value of σ_{int} . No knowledge of its amplitude is thus assumed, but only its redshift independence.

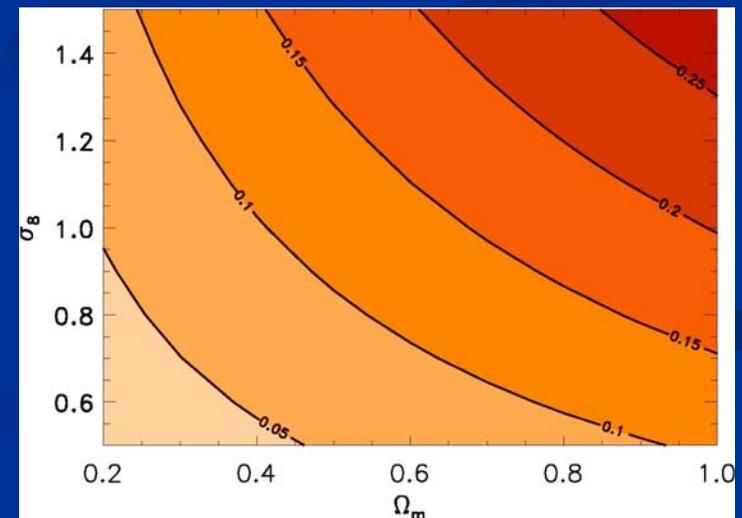
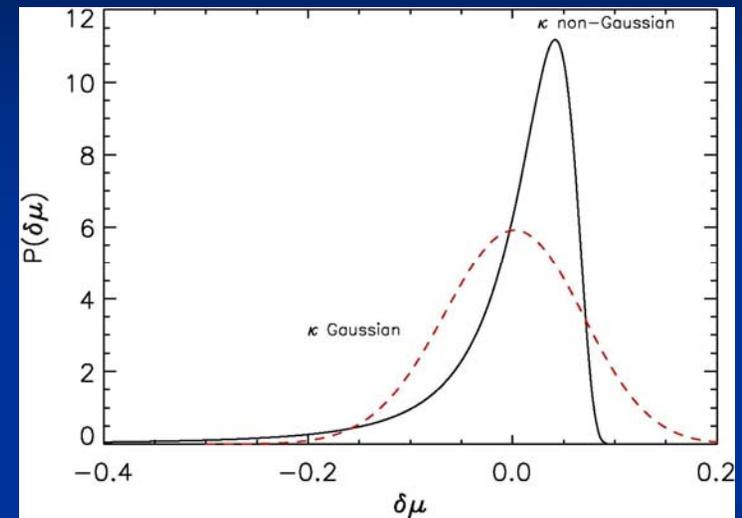
Likelihood Analysis: Results

- Ω_m is tightly constrained and unbiased (mostly by μ_0).
Holz and Linder [2004]
- The clustering parameter σ_8 is constrained to about 5% *but* is sensitive to the actual pdf assumed because it solely depends on σ_{cos} .
- 100 runs performed
 - The use of the correct non-gaussian pdf leads on average to the correct cosmology
 - The use of the gaussian pdf leads to an average bias of $\Delta\sigma_8=0.12$



The Origin of the Bias

- In brief, the bias is due to the wrong choice of the probability distribution for $\delta\mu_{\text{cos}}$.
- More specifically
 - The non-gaussian pdf generates data with a larger variance, thanks to its fat tail
 - This data can be explained by a gaussian with a larger cosmic dispersion σ_{cos}
 - But a larger variance can only be explained by a larger value of σ_8 .



Prospects for the Future and Conclusions

- More work ahead:
 - Include the effect of baryons on small scales
 - Determine whether the scatter in the intrinsic luminosity of SNIa has a redshift dependence
 - Further investigation of the pdf of $\delta\mu_{\text{cos}}$: extend to the case of modified gravity theories
- Conclusions
 - The scatter of SNIa distance modulus is not just a nuisance
 - it contains information about cosmological parameters
 - which can be successfully mined (provided the correct pdf are used)

“There is (at least some) gold in them hills”... and it is for free

A Few Details on the pdf Calculation

Given the reduced convergence η

$$\eta = 1 + \frac{\kappa}{|\kappa_{\min}|}$$

Wang et al. [2002] use the results of n-body simulations to calibrate the following fit for $P_{\eta}(\eta)$

$$P_{\eta}(\eta) = C \exp \left[- \left(\frac{\eta - \eta_p}{w \eta^q} \right)^2 \right]$$

Where, letting

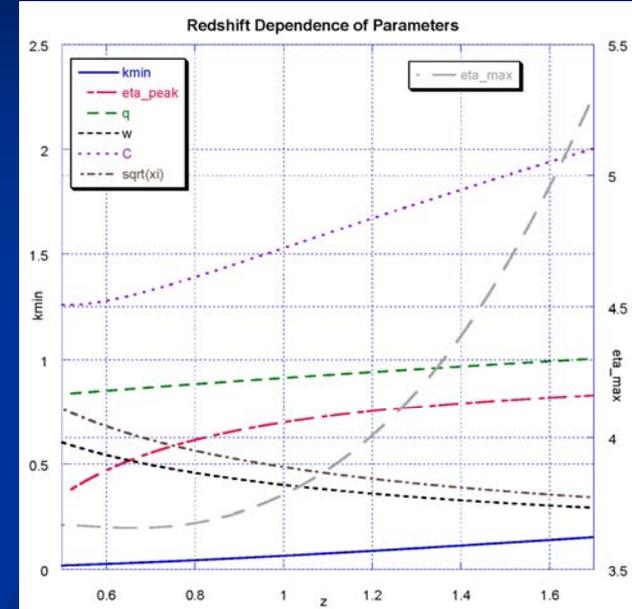
$$\xi_{\eta} = \frac{\langle \kappa^2 \rangle}{|\kappa_{\min}|^2}$$

The pdf parameters are then fitted to

$$\eta_p(\xi_{\eta}) = 1.002 - 1.145 \left(\frac{\sqrt{\xi_{\eta}}}{5} \right) - 20.427 \left(\frac{\sqrt{\xi_{\eta}}}{5} \right)^2$$

$$w(\xi_{\eta}) = 0.028 + 3.952 \left(\frac{\sqrt{\xi_{\eta}}}{5} \right) - 1.262 \left(\frac{\sqrt{\xi_{\eta}}}{5} \right)^2$$

$$q(\xi_{\eta}) = 0.702 + 0.509 \left(\frac{1}{5\sqrt{\xi_{\eta}}} \right) + 0.008 \left(\frac{1}{5\sqrt{\xi_{\eta}}} \right)^2$$



But it is possible to show that in the weak lensing limit

$$\eta = 1 + \frac{10^{\delta\mu/5}}{|\kappa_{\min}|}$$

Which allows the derivation of the pdf for $\delta\mu_{\text{cos}}$

Weak Lensing as a source of noise: Baryon Acoustic Oscillations

- Baryon Acoustic Oscillation are cosmological standard rulers. They appear as a peak in the correlation function of galaxies and quasars.
- Weak lensing alters the observed position of a source
⇒ **lensing-induced correlation**
This is an unavoidable source of error.
- Weak lensing alters the observed magnitude of a source
⇒ **magnification bias**
This suggests the class of objects to be used for the measurement

