

New results on sub-Chandrasekhar mass explosions of White Dwarfs

Markus Kromer

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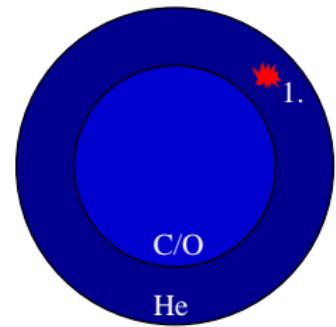
Max-Planck-Institut für Astrophysik, Garching, Germany

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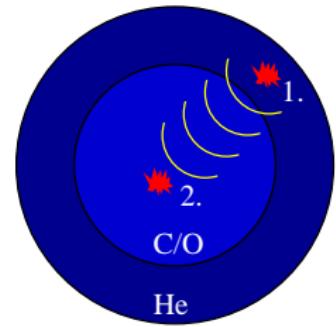
The double detonation scenario

- CO WD accretes He from a He-rich companion star
- He flash triggers a detonation in the shell



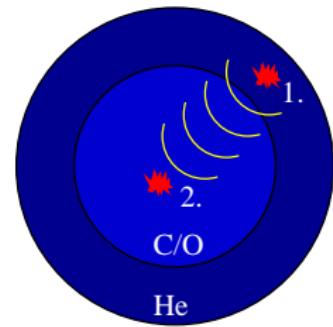
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- Shock-compression ignites a secondary detonation in the core



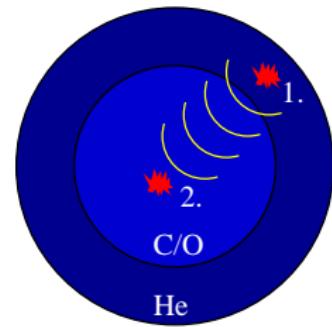
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- + WD mass direct parameter to explain range of observed ^{56}Ni yields
 - + Progenitors should be frequent
 - + Possible link to stellar population



The double detonation scenario

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- + WD mass direct parameter to explain range of observed ^{56}Ni yields
- + Progenitors should be frequent
- + Possible link to stellar population
- Robustness of core ignition
- Problems in fitting observational data (Höflich & Khokhlov 1996, Nugent et al. 1997)

New hydro simulations

- Less massive helium shells than previously thought may detonate (Bildsten et al. 2007)
- Can these shell detonations trigger core detonations?

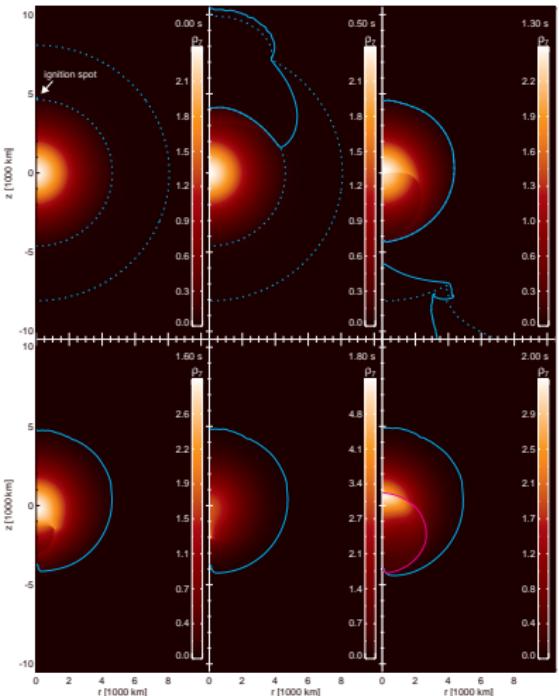
New hydro simulations

- Less massive helium shells than previously thought may detonate (Bildsten et al. 2007)
- Can these shell detonations trigger core detonations?
- Fink et al. (2010) investigate six models

Model	M_{tot}/M_{\odot}	$M_{\text{core}}/M_{\odot}$	$M_{\text{shell}}/M_{\odot}$
1	0.936	0.810	0.126
2	1.004	0.920	0.084
3	1.080	1.025	0.055
4	1.164	1.125	0.039
5	1.293	1.280	0.013
6	1.389	1.385	0.004

New hydro simulations

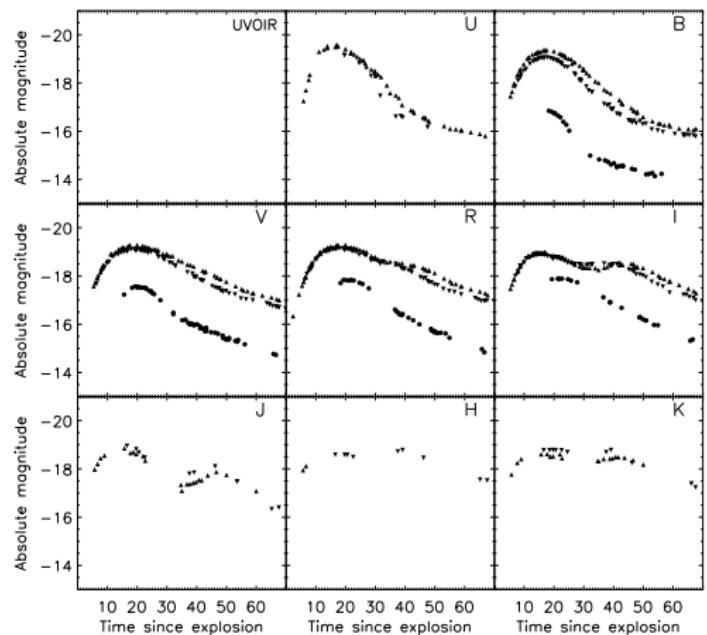
- Less massive helium shells than previously thought may detonate (Bildsten et al. 2007)
- Can these shell detonations trigger core detonations?
- Fink et al. (2010) investigate six models
 - Surface reduction increases shock strength
 - All models successfully ignite a detonation in the C/O core



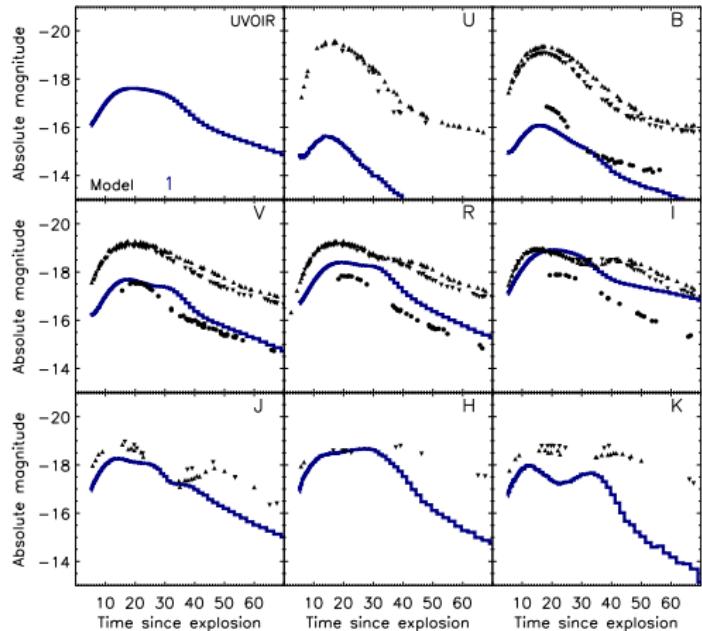
Now calculate synthetic observables for these models applying the ARTIS code (Kromer & Sim 2009, Sim 2007)

- Multi-wavelength: γ to NIR
 - Time-dependent
 - Fully 3D
 - Detailed solution of ionisation and thermal balance equation
 - Detailed treatment of radiation/matter interactions
- ⇒ Parameter-free prediction of synthetic observables

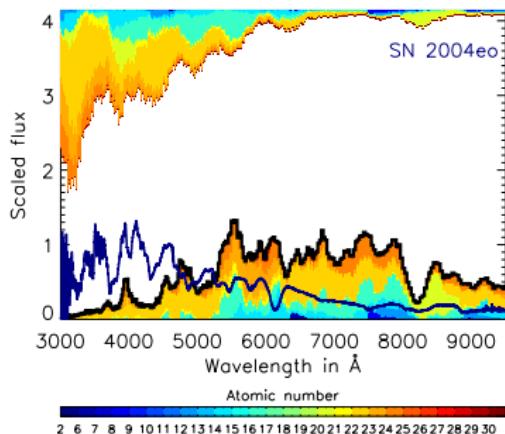
Synthetic observables (Kromer et al., ApJ in press)



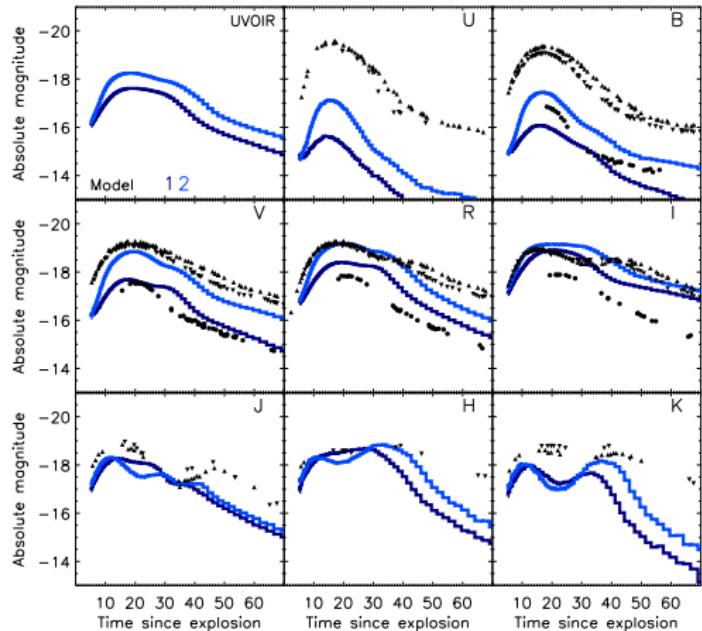
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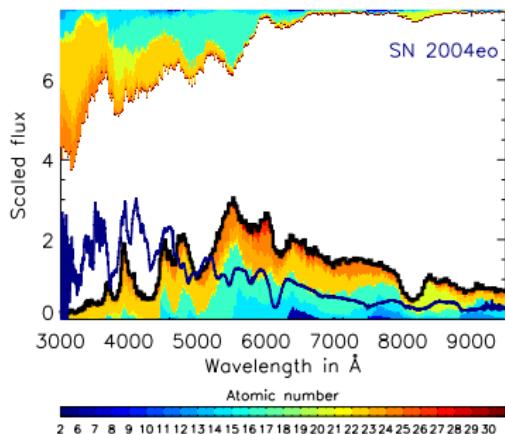
Model 1	Core	Shell
M	0.810	0.126
$M(^{56}\text{Ni})$	1.7×10^{-1}	8.4×10^{-4}
$M(\text{Ti})$	3.9×10^{-4}	1.1×10^{-2}
$M(\text{Si})$	2.7×10^{-1}	4.8×10^{-4}



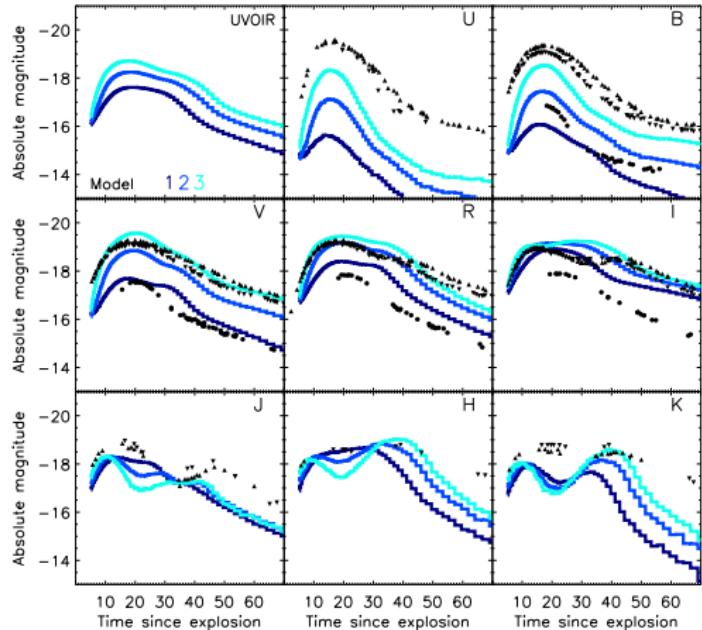
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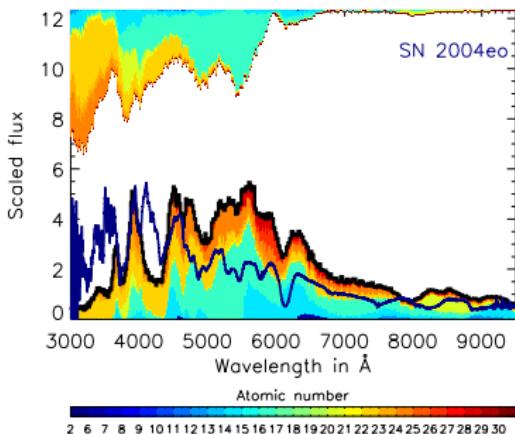
Model 2	Core	Shell
M	0.920	0.084
$M(^{56}\text{Ni})$	3.4×10^{-1}	1.1×10^{-3}
$M(\text{Ti})$	4.6×10^{-4}	7.8×10^{-3}
$M(\text{Si})$	2.5×10^{-1}	2.5×10^{-4}



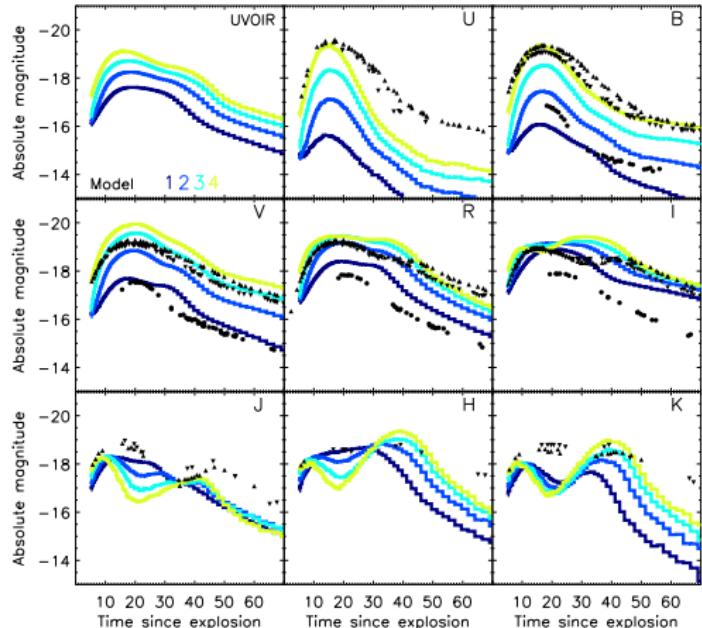
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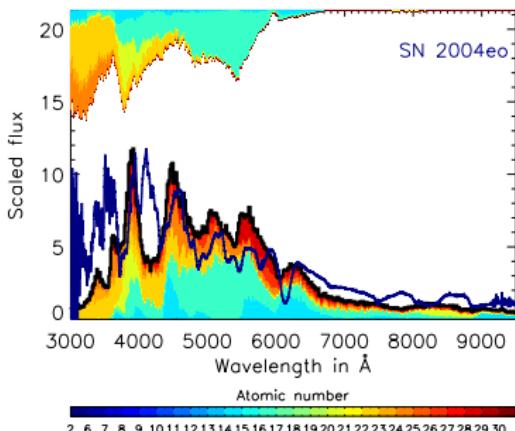
Model 3	Core	Shell
M	1.025	0.055
$M(^{56}\text{Ni})$	5.5×10^{-1}	1.7×10^{-3}
$M(\text{Ti})$	4.5×10^{-4}	4.4×10^{-3}
$M(\text{Si})$	2.1×10^{-1}	1.4×10^{-4}



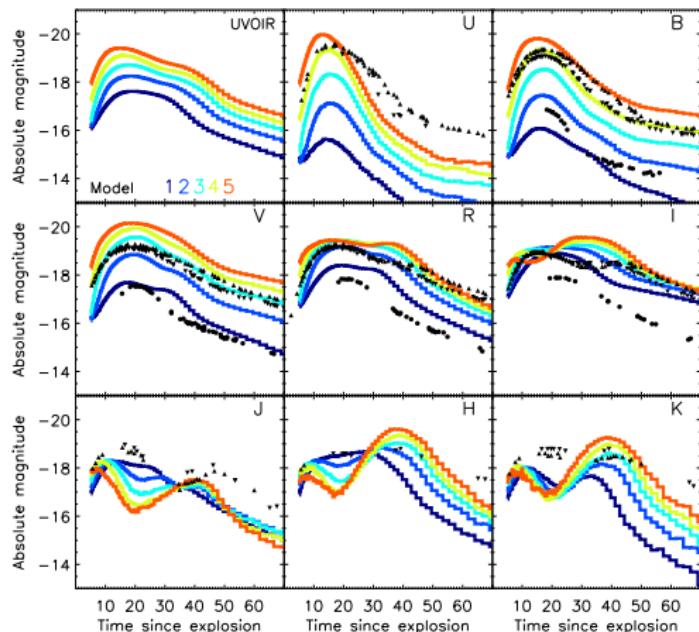
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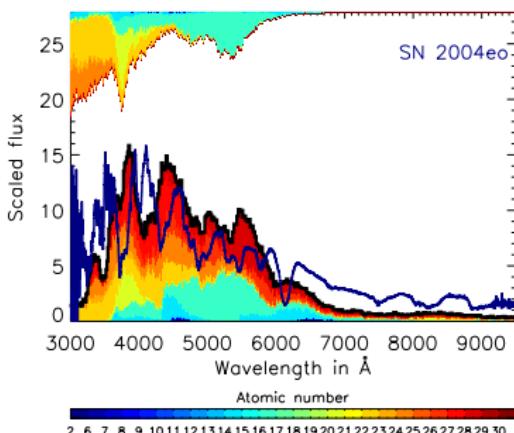
Model 4	Core	Shell
M	1.125	0.039
$M(^{56}\text{Ni})$	7.8×10^{-1}	4.4×10^{-3}
$M(\text{Ti})$	3.8×10^{-4}	2.2×10^{-3}
$M(\text{Si})$	1.4×10^{-1}	4.7×10^{-4}



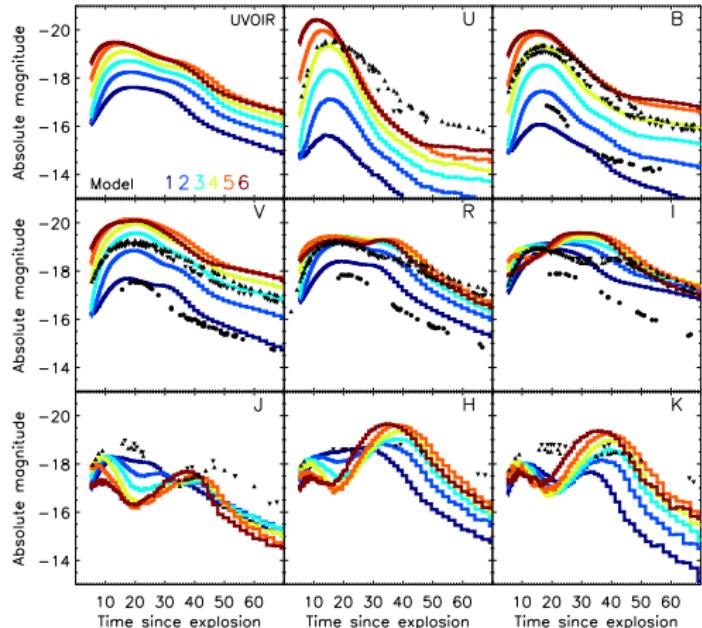
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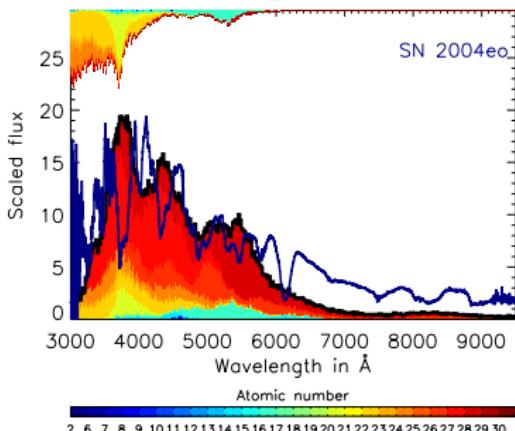
Model 5	Core	Shell
M	1.280	0.013
$M(^{56}\text{Ni})$	1.05	1.5×10^{-3}
$M(\text{Ti})$	2.1×10^{-4}	6.8×10^{-4}
$M(\text{Si})$	6.1×10^{-2}	1.6×10^{-4}



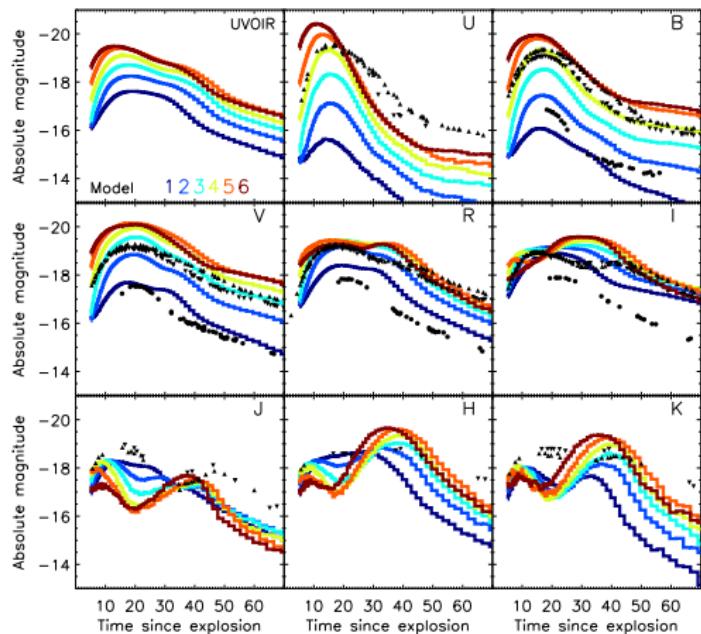
Synthetic observables (Kromer et al., ApJ in press)



Model 6	Core	Shell
M	1.385	0.004
$M(^{56}\text{Ni})$	1.10	5.7×10^{-4}
$M(\text{Ti})$	7.1×10^{-5}	1.5×10^{-4}
$M(\text{Si})$	1.5×10^{-2}	1.3×10^{-4}



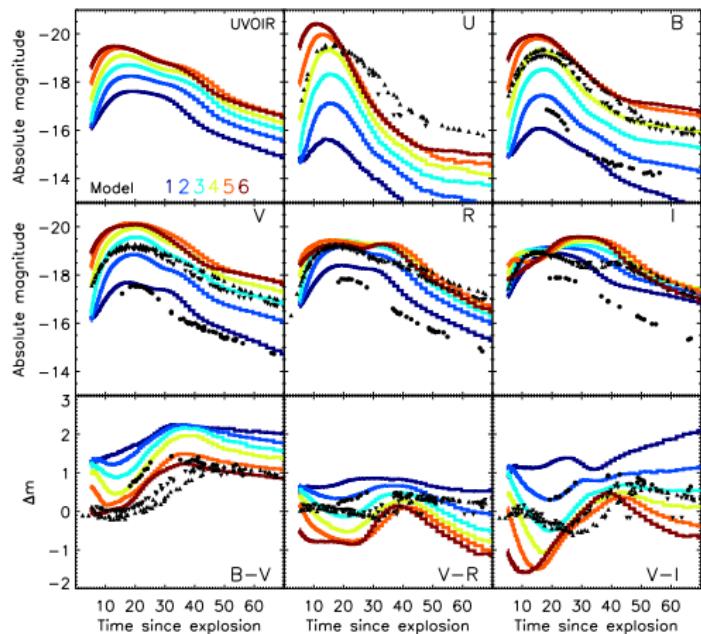
Synthetic observables (Kromer et al., ApJ in press)



Promising models

- + Populate a large range in brightness
- + Despite low mass, time-evolution OK

Synthetic observables (Kromer et al., ApJ in press)



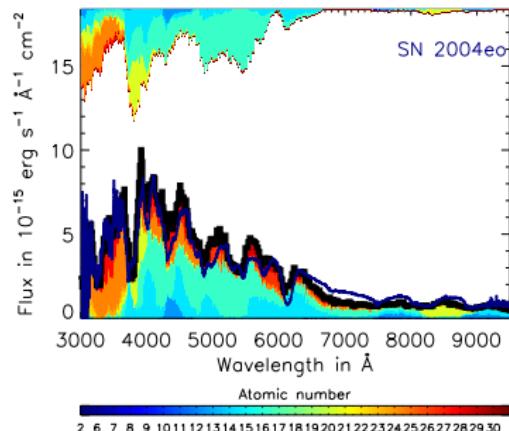
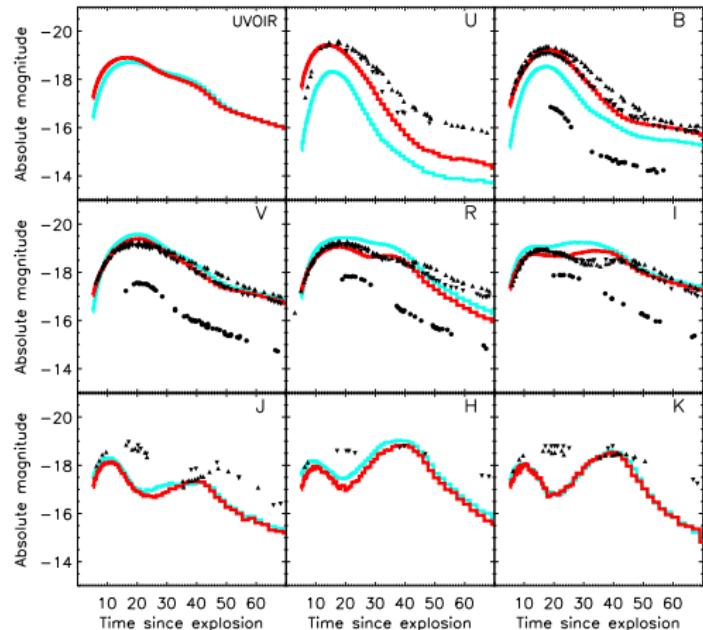
Promising models

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Problems

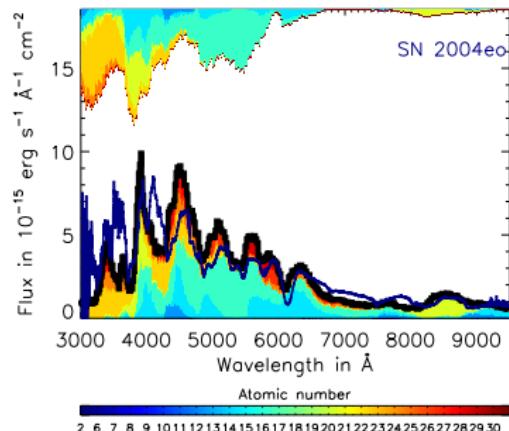
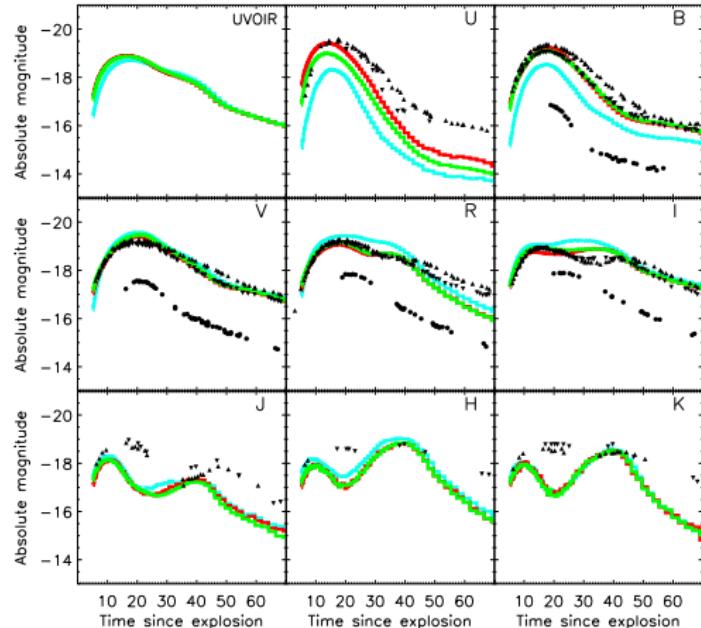
- Peculiar light curves and spectra
- Colours too red

Influence of the helium shell



Model 3
Shell-less toy version of Model 3

Prospects: modifying the initial shell composition



Model 3

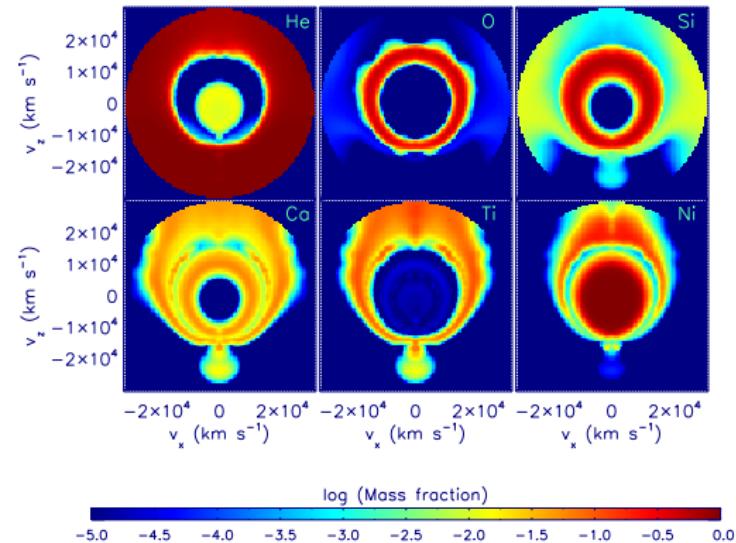
Shell-less toy version of Model 3

Model 3, initially 34 % ^{12}C in shell

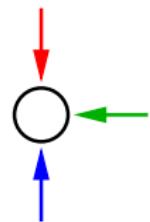
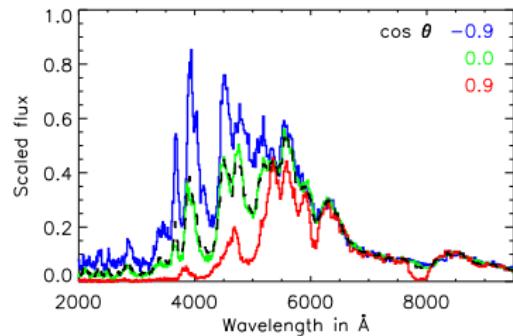
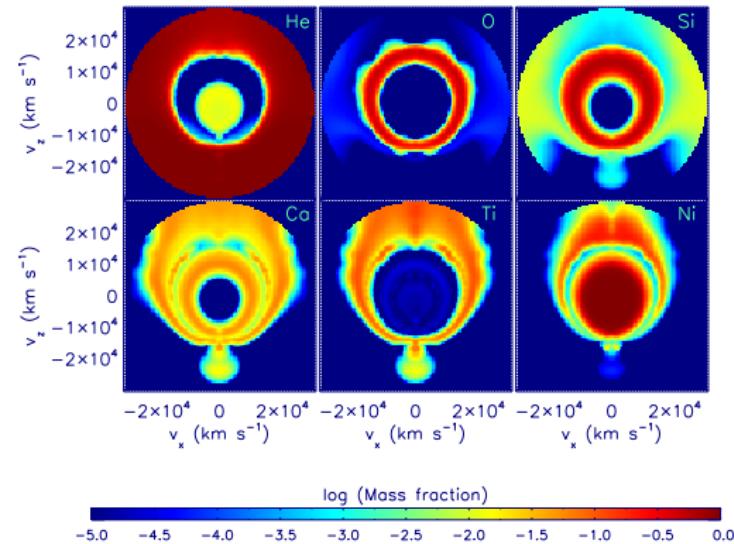
Summary

- Secondary core detonation ignites even for minimum helium shell masses (Fink et al. 2010)
- Models predict correct range of brightness and rise times (Kromer et al. arXiv:1006.4489)
- However, spectral features and colours do not agree
- Origin of this discrepancy is the helium shell material
- Tiny amounts of iron-group elements in outer layers cause severe problems
- Nucleosynthesis?
- Radiative transfer?

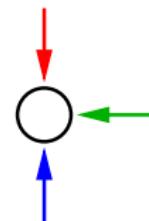
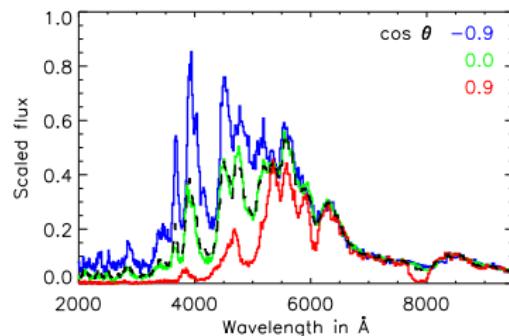
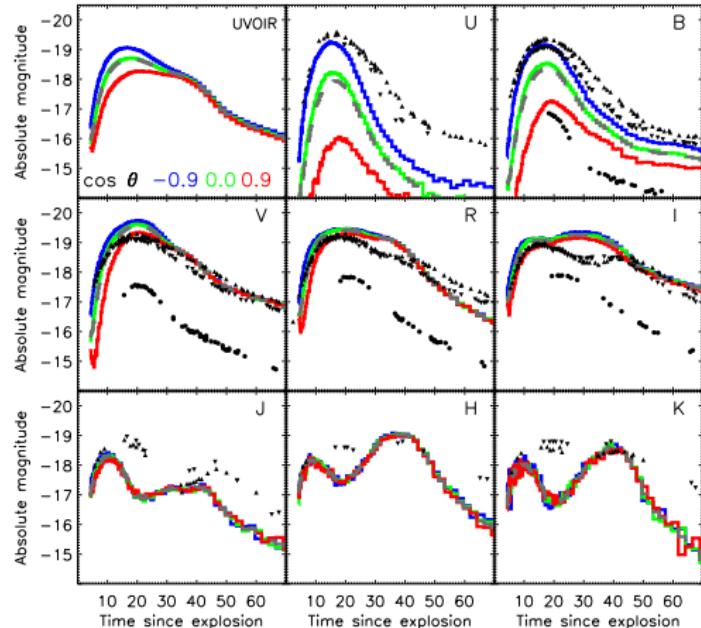
Model 3 as an example for line-of-sight effects



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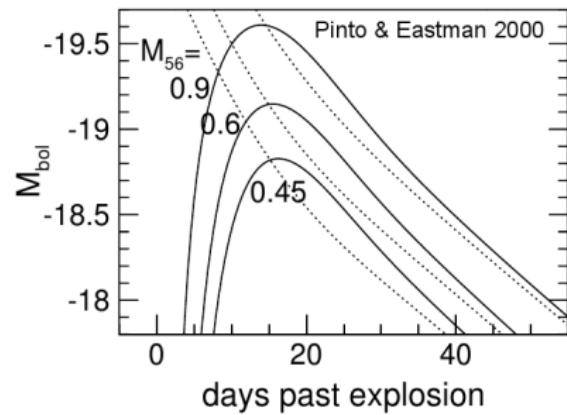


Model 3 as an example for line-of-sight effects

Sim et al. (2010) investigated pure detonations of naked sub-Chandrasekhar mass WDs

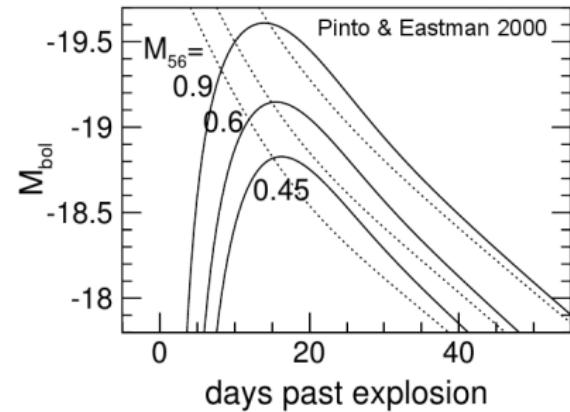
- Set of 1D models ($0.88 - 1.15 M_{\odot}$)
- Hydrodynamics, nucleosynthesis and radiative transfer
- Compared to observed SNe Ia the models explain
 - Range of brightness
 - Rise times
 - Colours
 - Silicon-line ratio
 - Velocity evolution of silicon lines
 - Width-luminosity relation

Outline of the problem



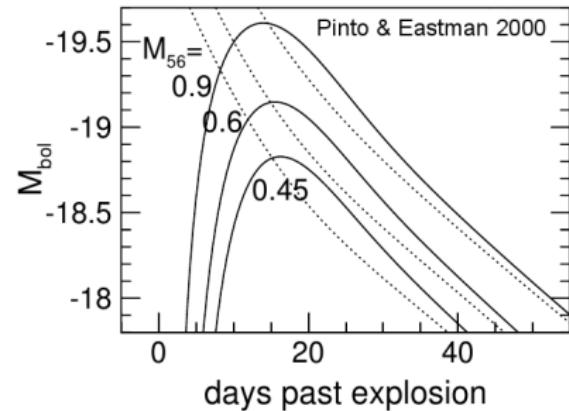
Outline of the problem

- Multi-wavelength



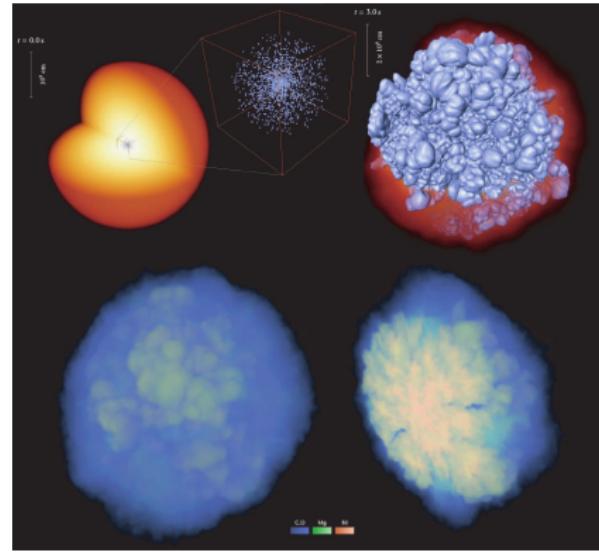
Outline of the problem

- Multi-wavelength
- Time-dependent



Outline of the problem

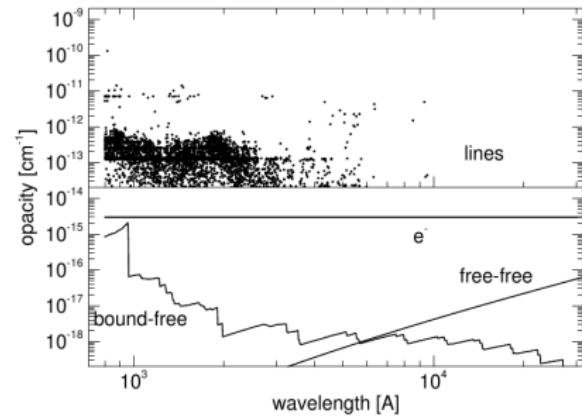
- Multi-wavelength
- Time-dependent
- Multi-dimensional



Röpke et al. 2007

Outline of the problem

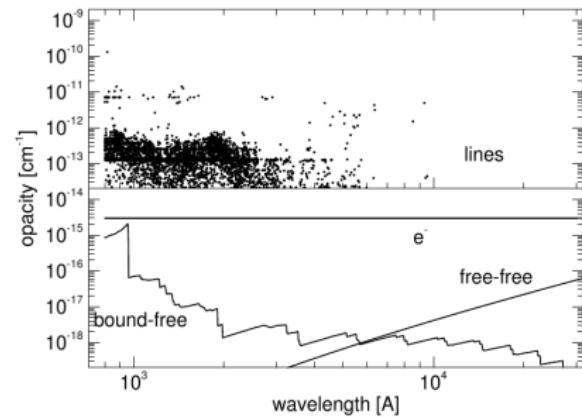
- Multi-wavelength
- Time-dependent
- Multi-dimensional
- Opacity dominated by lines



Pinto & Eastman 2000

Outline of the problem

- Multi-wavelength
- Time-dependent
- Multi-dimensional
- Opacity dominated by lines
- Non-LTE effects important



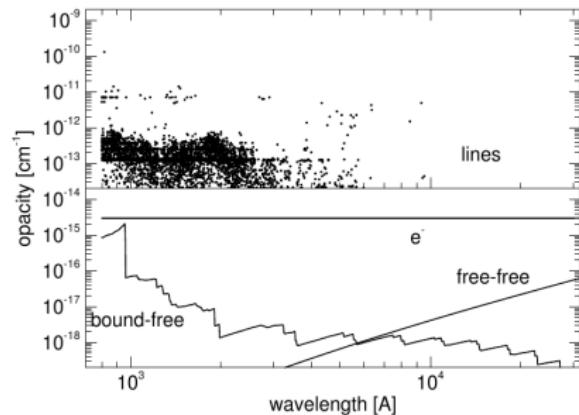
Pinto & Eastman 2000

Outline of the problem

- Multi-wavelength
- Time-dependent
- Multi-dimensional
- Opacity dominated by lines
- Non-LTE effects important

But some simplifications

- Homologous expansion
- Sobolev approximation
- Statistical and thermal equilibrium

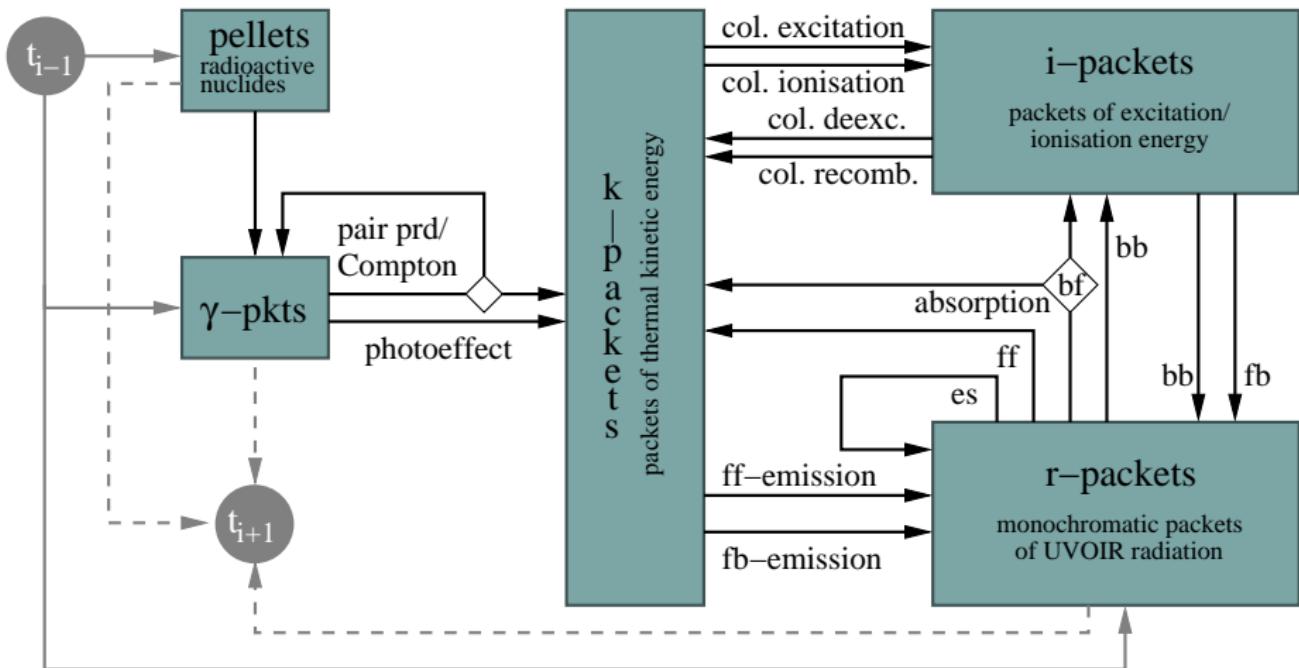


Pinto & Eastman 2000

Monte Carlo method

- Based on quantized energy flow: energy packets
- Follow the packets propagation through the ejecta
- Microphysical description of radiation/matter interactions
 - ⇒ Purely local
 - ⇒ Suitable for complex geometries & time-dependence
- Extract spectra and light curves by binning of escaping packets
- Use **indivisible energy packets** (Abbott & Lucy 1985; Mazzali & Lucy 1993; Lucy 1999, 2005)
 - ⇒ Implicit energy conservation
 - ⇒ Statistical and thermal equilibrium enforceable (Lucy 2002, 2003)

The framework of ARTIS (Kromer & Sim 2009)



Calculation of transition probabilities requires

- ① Specification of atomic data
- ② Population numbers (excitation/ionization state of the plasma)
- ③ Local radiation field J_ν

Calculation of transition probabilities requires

- ① Specification of atomic data
 - CD23: 4×10^5 bound-bound transitions
 - BIG: 8×10^6 bound-bound transitions
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Calculation of transition probabilities requires

① Specification of atomic data

- CD23: 4×10^5 bound-bound transitions
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② Population numbers (excitation/ionization state of the plasma)

- Complete set of NLTE rate equations too expensive
- Instead approximate NLTE treatment (**detailed**)
 - Consistent solution of photoionization and thermal balance
 - Boltzmann excitation formula
- For comparison: LTE treatment (**simple**)
 - Saha ionization formula
 - Boltzmann excitation formula

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③ Local radiation field J_ν

- Extractable from MC simulation, but computationally prohibitive
- Nebular approximation for **detailed** treatment: $J_\nu = WB_\nu(T_R)$
- Black body approximation for **simple** treatment: $J_\nu = B_\nu(T_J)$

Excitation/ionisation treatment

- detailed solution of the ionisation balance
 - assume photoionisation equilibrium

$$\frac{N_{j,k}}{N_{j+1,k} n_e} = \frac{\alpha_{j,k}^{\text{sp}}}{\Gamma_{j,k}}$$

- derive $\Gamma_{j,k}$ from Monte Carlo simulation

$$\Gamma_{j,k} \equiv \frac{g_{0,j,k}}{U_{j,k} n_{0,j,k}} \cdot \sum_{i=0}^{N_{j,k}} n_{i,j,k} \gamma_{i,j,k}$$

- simultaneous solution of the thermal balance equation $\Rightarrow T_e$
 - heating rates from Monte Carlo simulation
 - cooling rates evaluated at T_e
- use Boltzmann formula evaluated at $T_J = \frac{\pi}{\sigma^4} \langle J \rangle$ for excitation

Radiation field

- exact radiation field extractable by Monte Carlo estimators

$$J_\nu d\nu = \frac{1}{4\pi\Delta t V} \sum_{d\nu} \epsilon_\nu^{\text{cmf}} d\mathbf{s}$$

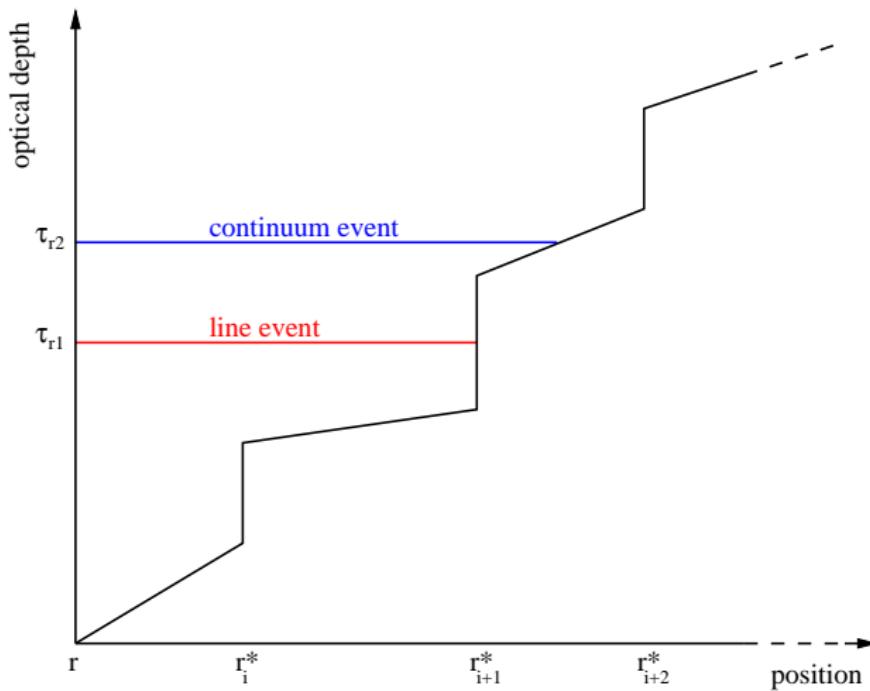
- but:* computationally prohibitive
⇒ parameterise local radiation field in **nebular approximation**

$$J_\nu = W \cdot B_\nu (T_R)$$

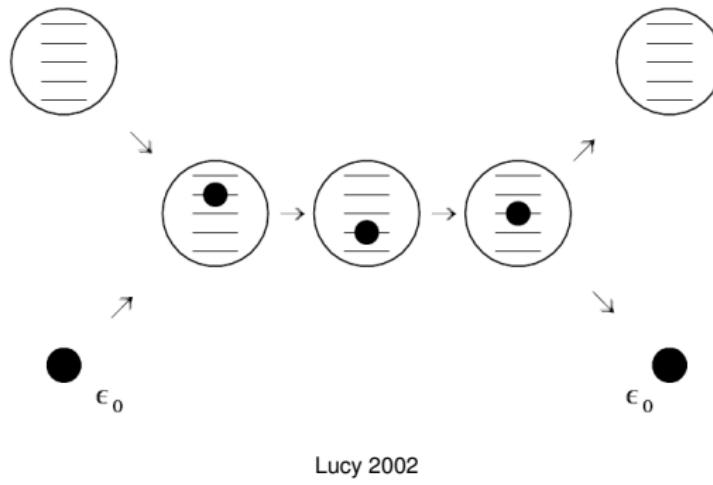
dilution factor W and radiation temperature T_R defined as

$$W = \frac{\pi}{\sigma T_R^4} \langle J \rangle \quad T_R = \frac{h \langle \nu \rangle}{3.832 k_B}$$

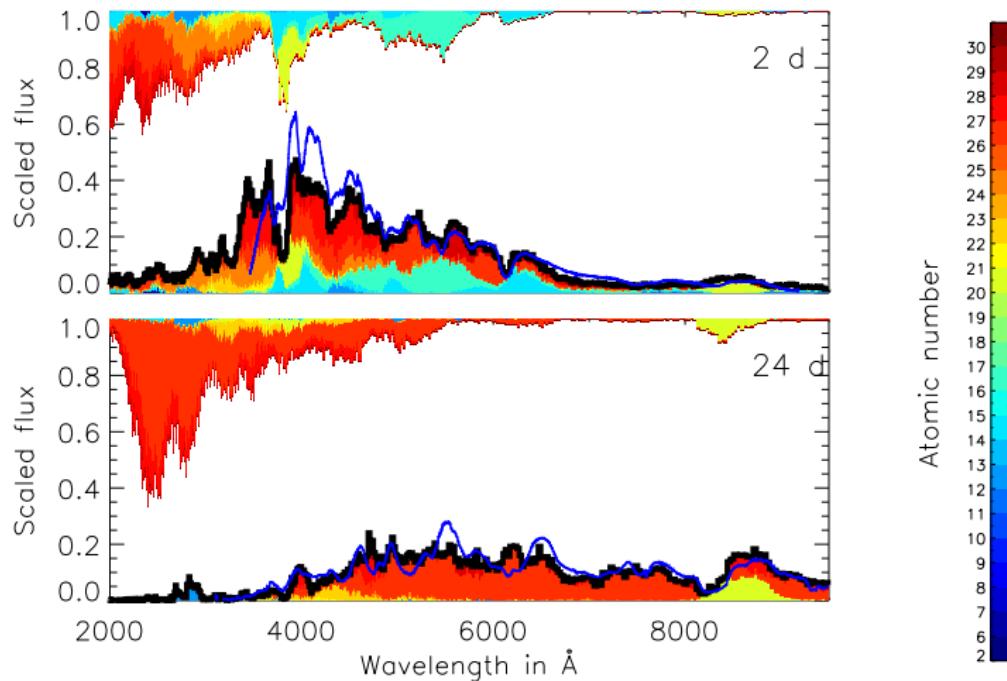
Selecting the next event



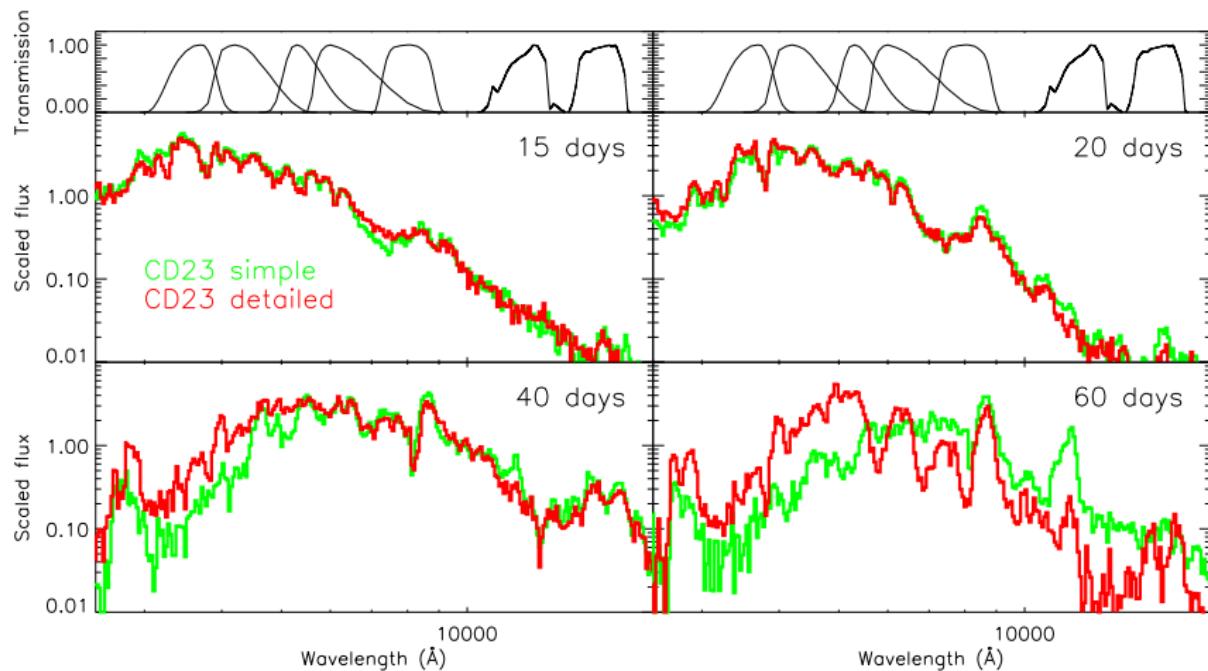
Macro atom formalism



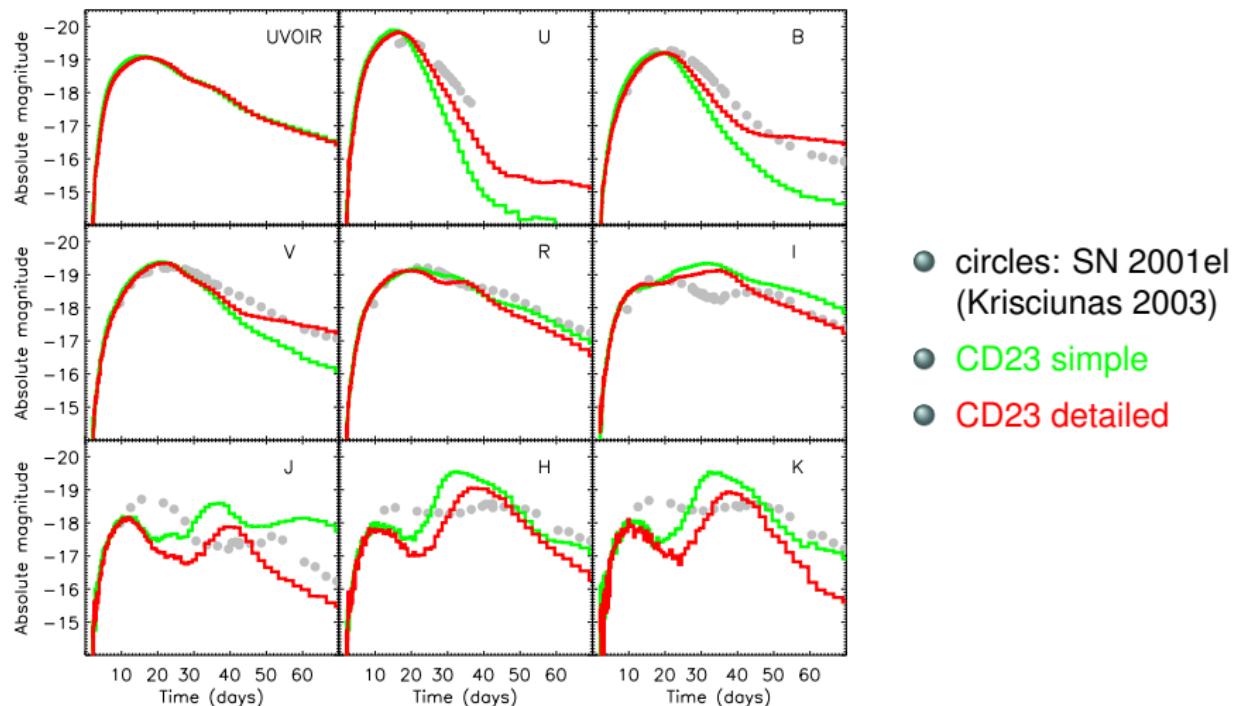
Spectral evolution



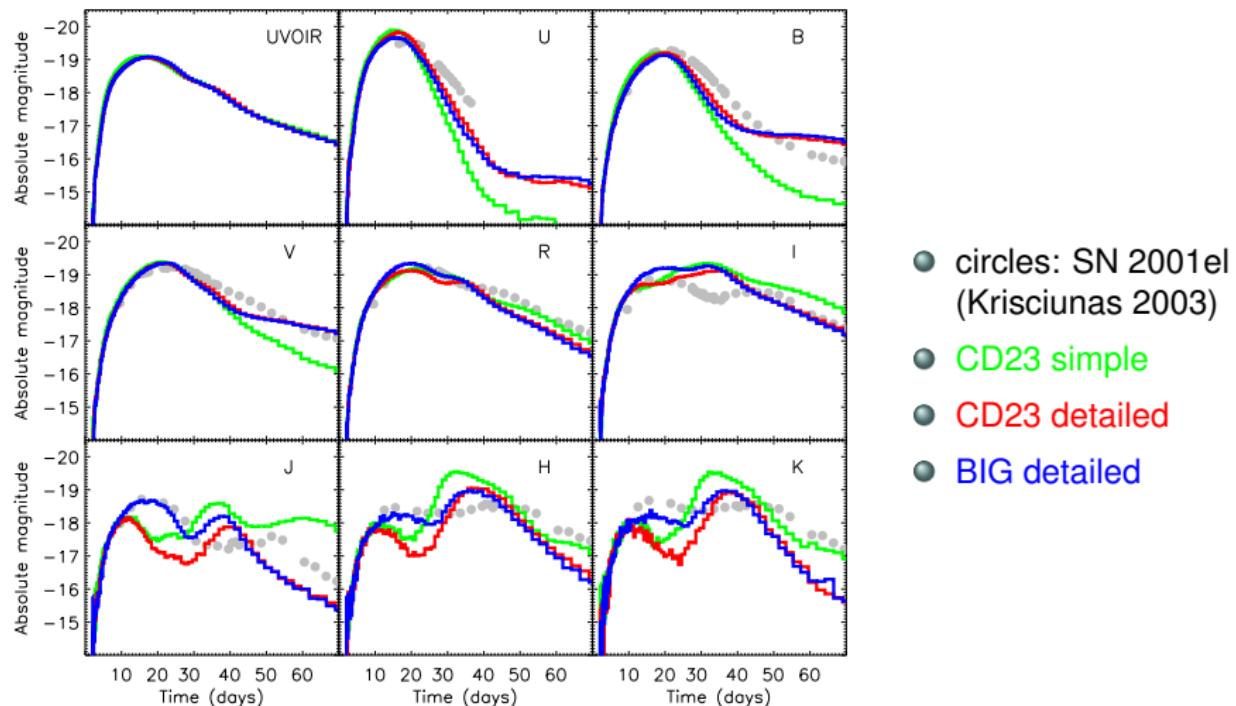
Influence of ionisation treatment



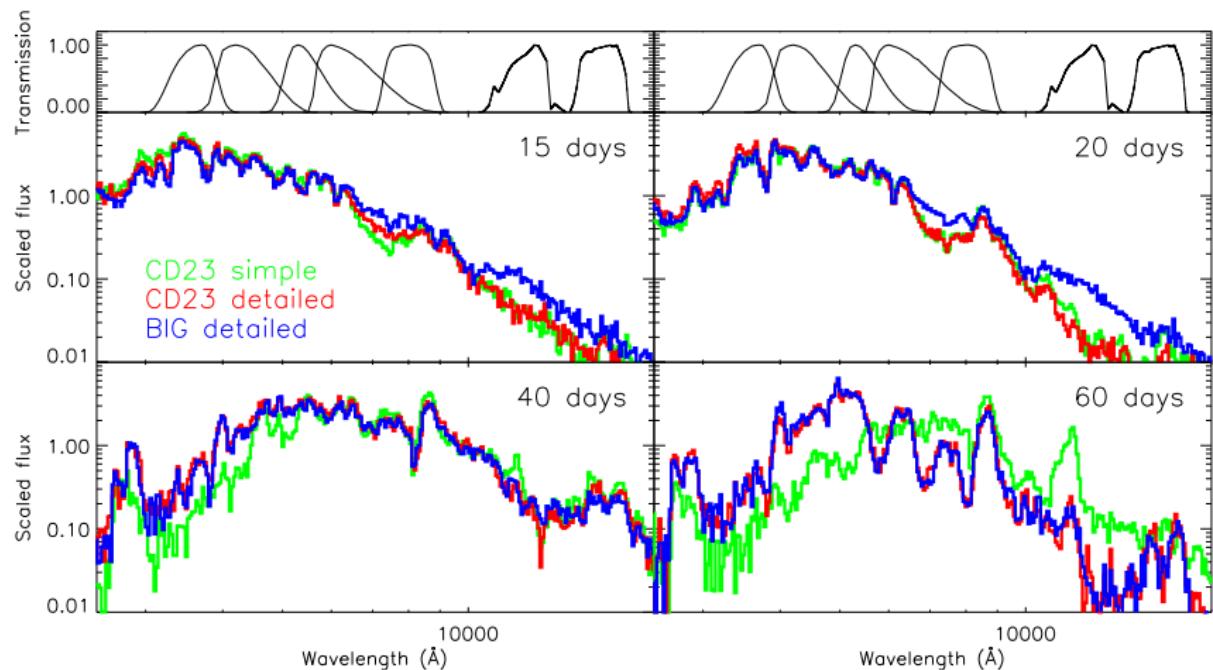
Influence of ionisation treatment



Influence of atomic data

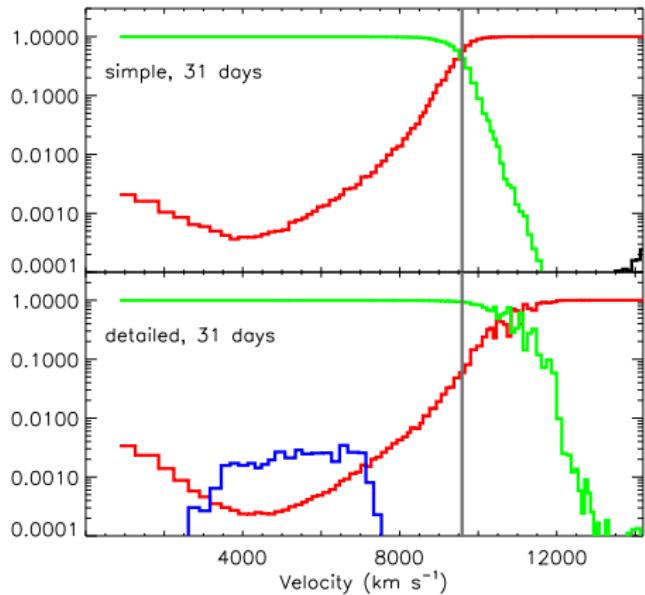


Influence of atomic data



Influence of ionisation treatment

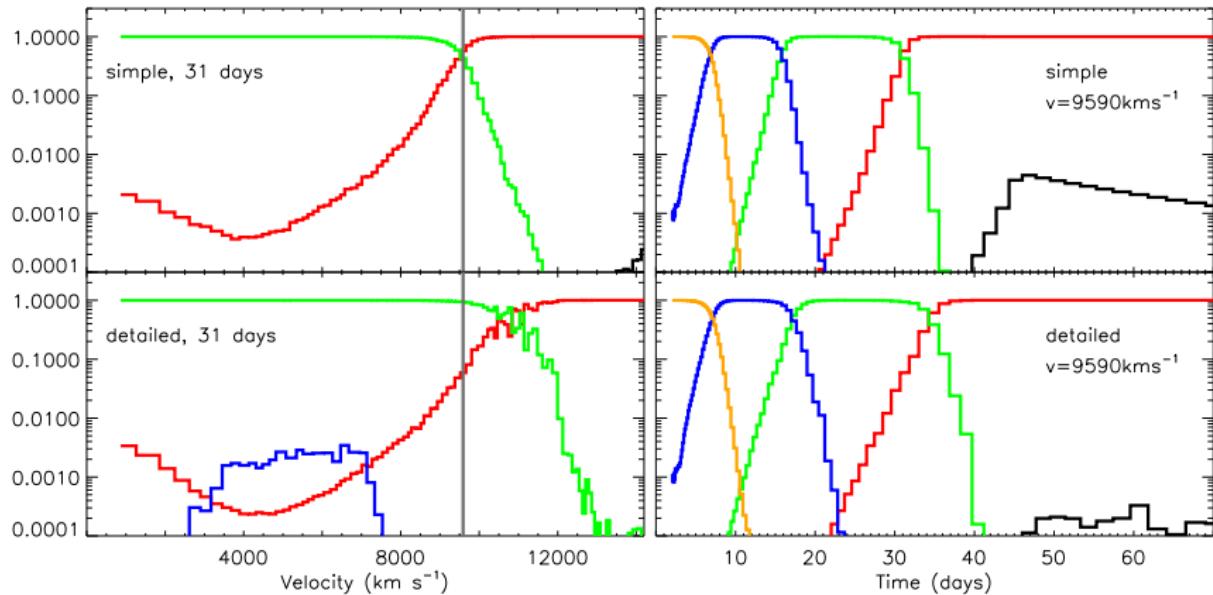
Ionisation fractions of Fe I, II, III, IV, V
versus radial velocity



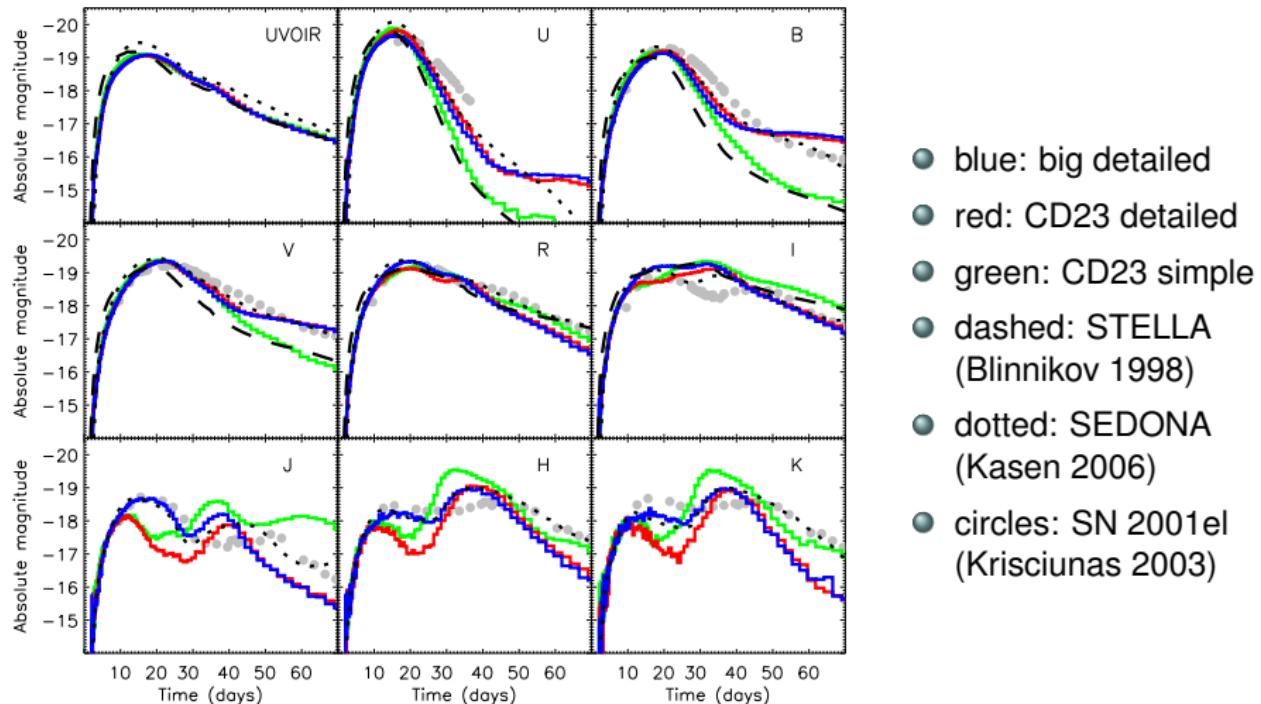
Influence of ionisation treatment

Ionisation fractions of Fe I, II, III, IV, V
versus radial velocity

versus time



Broad-band light curves



Flux redistribution

