

# String Inflation: Why?

*Convoluting models  
walk the Planck*



*C.P. Burgess*



# Outline

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- Context: *Occam vs Wilson*
  - Why build models only a mother could love?
- Some UV – sensitive issues
  - Naturalness;
  - Large field excursions and tensor modes;
  - Robustness
  - Reheating vs insulation; mind broadening; initial conditions; ...
- A tentative scorecard

# Outline

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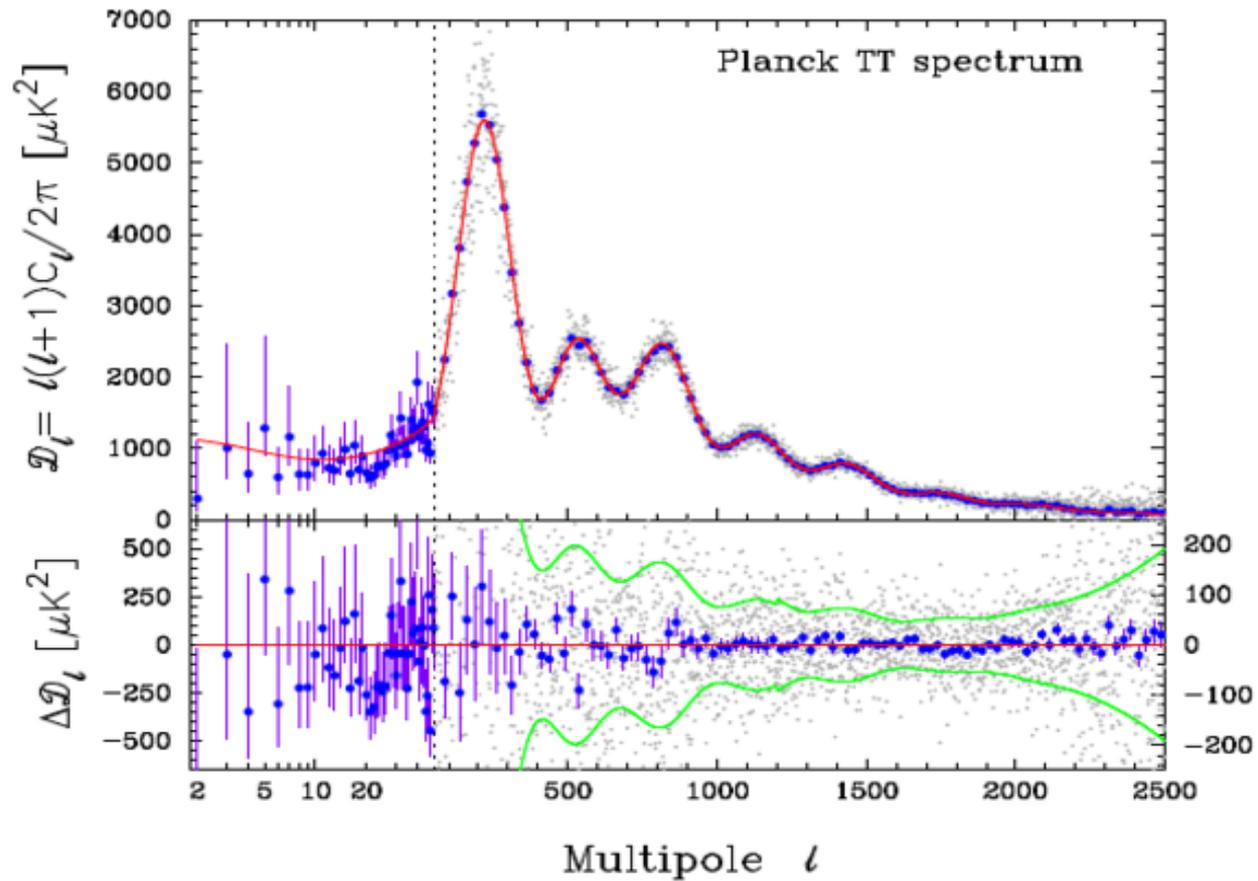
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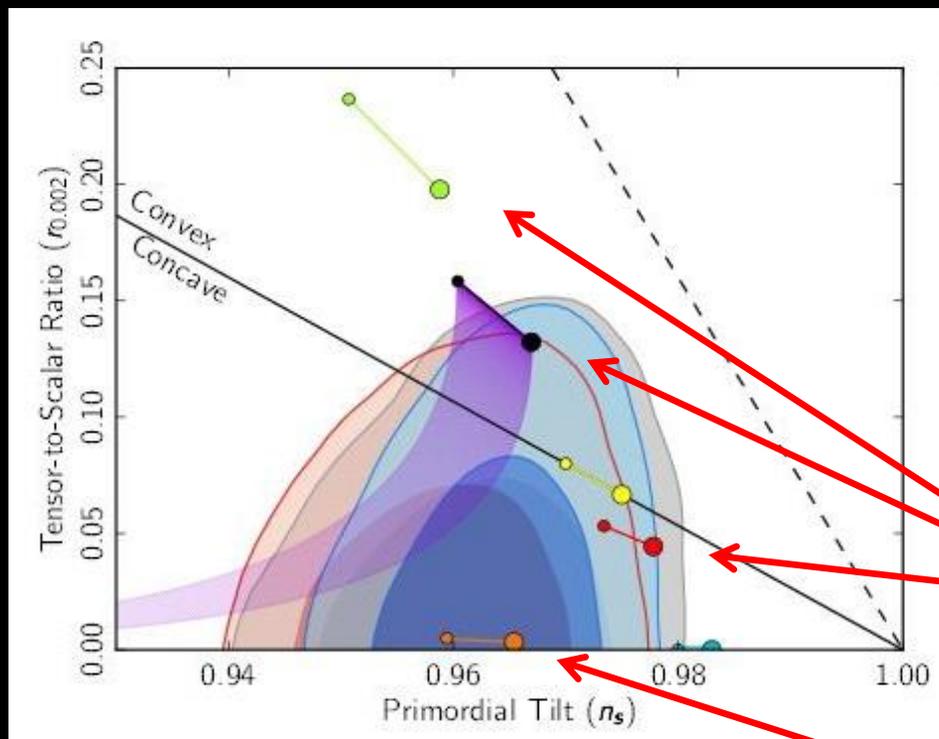


# Context



# Context

Planck 2013



$$P_S(k) = A_S k^{n_s-1}$$

$$r = \frac{A_T}{A_S}$$

$$V(\varphi) \propto \varphi^p$$

Simple models work well

$$L = M^2 R + \gamma R^2$$

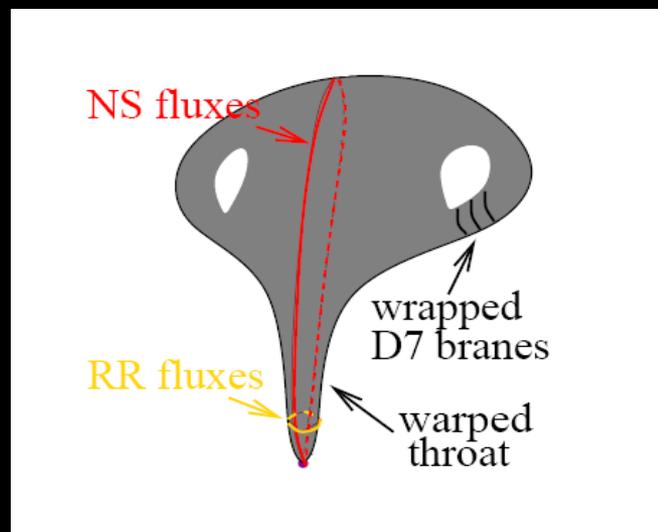
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String cosmology involves complicated constructions (extra dimensions, wrapped branes, fluxes, gaugino condensation, supersymmetry, ...)

eg:

$$K = -2 \ln V = -2 \ln(\tau_1^{3/2} - \tau_2^{3/2})$$

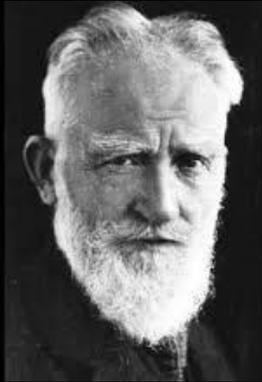
$$W = W_0 + A_1 e^{a_1 T_1} + A_2 e^{a_2 T_2}$$





# Context

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*Divided by a common language*

# Context



*Divided by a common language*



# Context

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- **Occam:** *What is the simplest possible model that the data requires?*
- **Wilson:**

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- **Occam:** *What is the simplest possible model that the data requires?*
- **Wilson:** *Low energy limit is often messy. What is generic and stable?*

*eg SUSY vs simple dark matter model*

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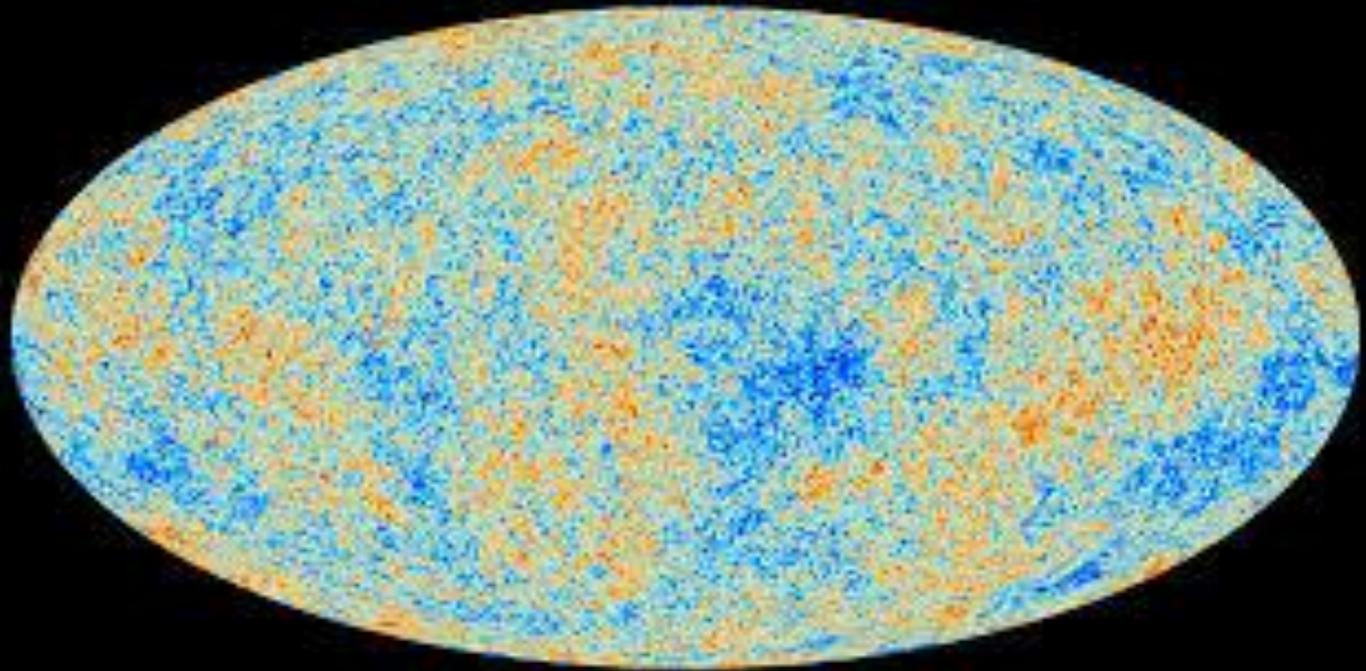
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- **Occam:** *What is the simplest possible model that the data requires?*
- **Wilson:** *Low energy limit is often messy. What is generic and stable?*

*Why embed into UV theory? Conceptual problems of simple models often require it...*

# Context

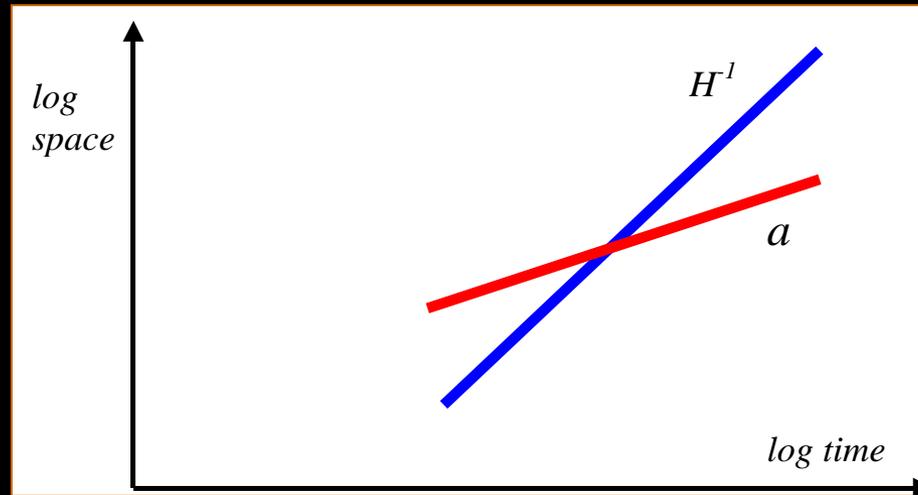
*Planck 2013*



*Primordial fluctuations*

Paris Dec 2014

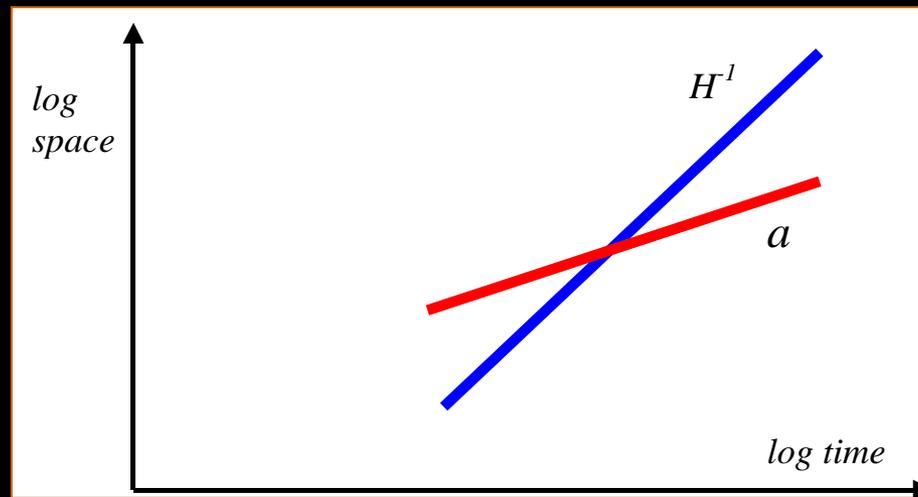
# Context



The problem: *how to generate correlations over scales that were super-Hubble in size?*

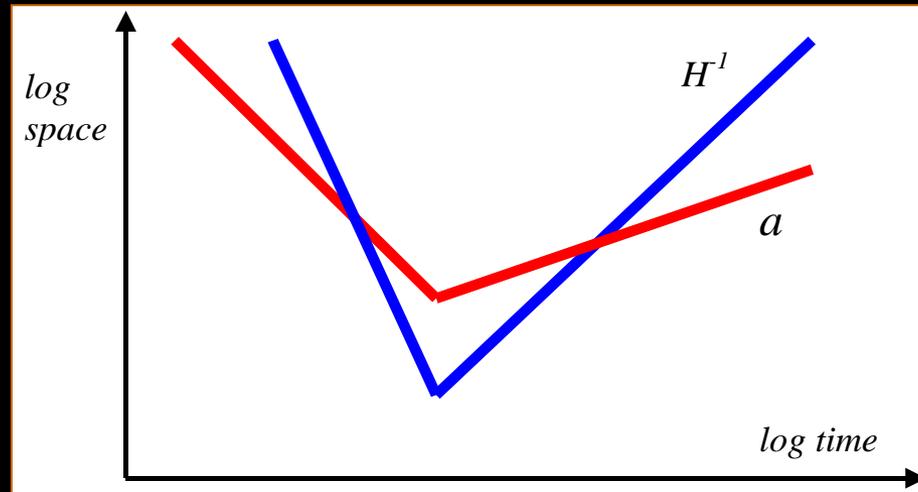


# Context



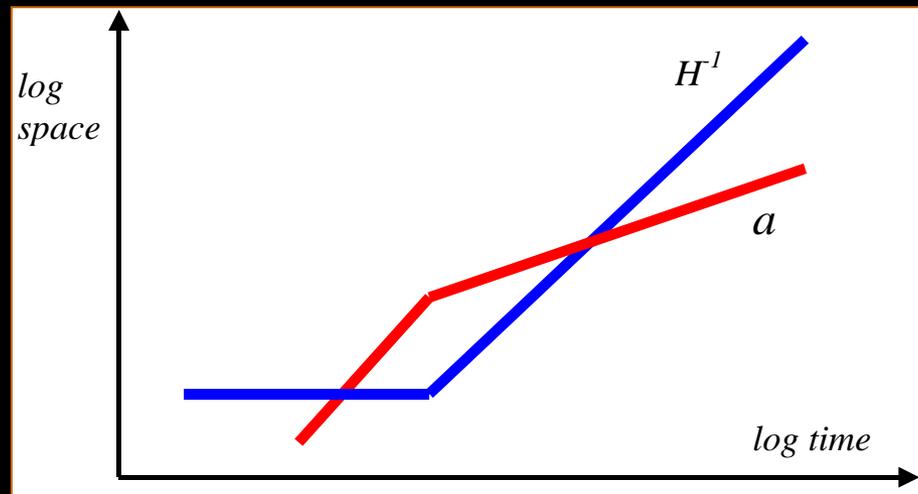
Usual way out: *is the extrapolation wrong?*

# Context



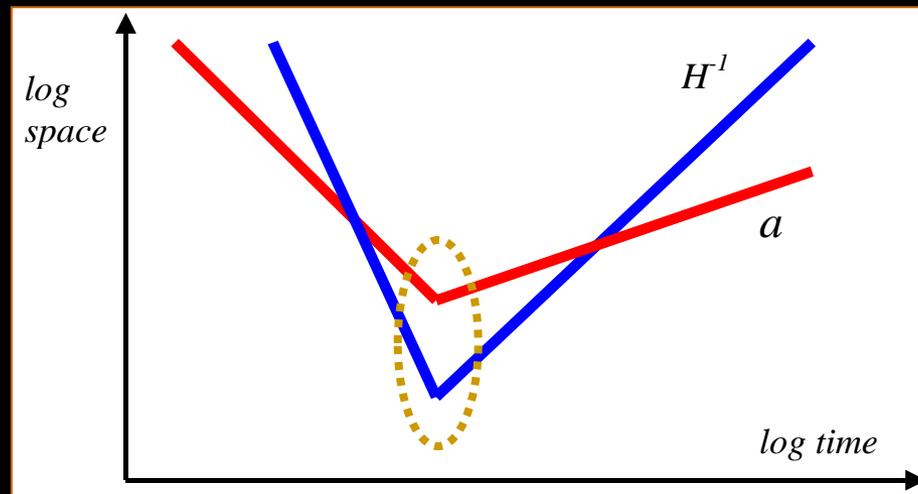
Possible solution: *the expansion could have accelerated...eg a bounce*

# Context



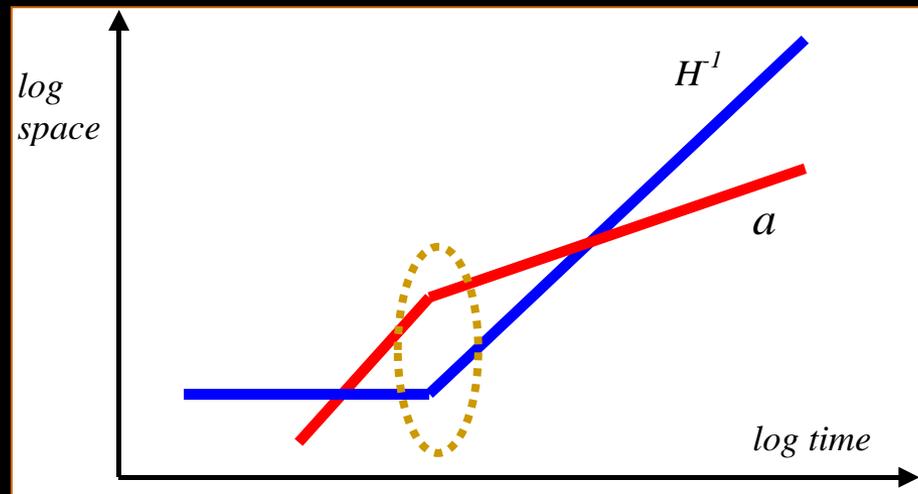
Possible solution: *the expansion could have accelerated...eg inflation*

# Context



*But: can and why should it bounce?*

# Context



But: *can and why should inflation start or end?*

# Context

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Some problems best addressed with a UV completion:

*If a bounce, why is the bounce semi-classical?*

*Why does inflation start or end?*

*Why are scalars light (why scale invariant)?*

*What does one do with a landscape/multiverse?*

*Why doesn't trans-Planckian physics intrude?*

*Why doesn't energy dissipate?*

*Why aren't initial conditions important?*

*and so on.....*

# Context

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Inflation: *impressive phenomenology seeking a fundamental theory (on the edge of the quantum gravity regime) within which to sit;*

# Context

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*Inflation: impressive phenomenology seeking a fundamental theory (on the edge of the quantum gravity regime) within which to sit;*

*String Theory: an impressive fundamental theory seeking experimental consequences to test*



# Context

Rocky Kolb

*Happy Valentine's Day!*

Inflation and strings were made for each other. – *Cliff Burgess*

From *Physical Review D* Personal ads:

Mature paradigm with firm observational support seeks a fundamental theory in which to be embedded. No loop quantum gravity theories, please. Contact [alan@mit.edu](mailto:alan@mit.edu).

Elegant theory of everything desires to explore the landscape with a phenomenon in the hope that it will lead to a prediction. Let's get physical! Contact [ed@ias.edu](mailto:ed@ias.edu).

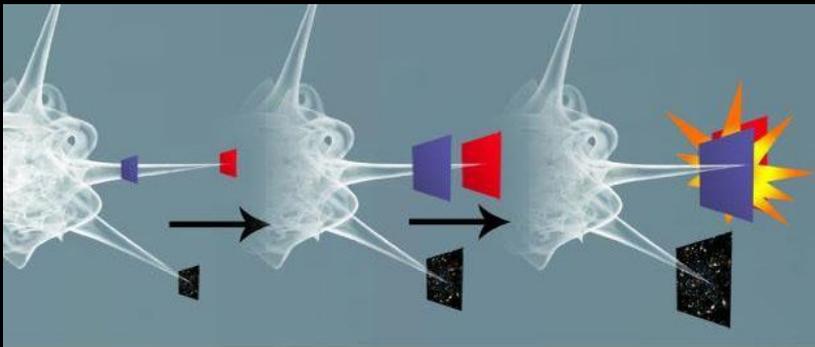


# Context

Superficially easy to find inflation in string theory:

It has many 4D scalars. (eg: brane positions; moduli)

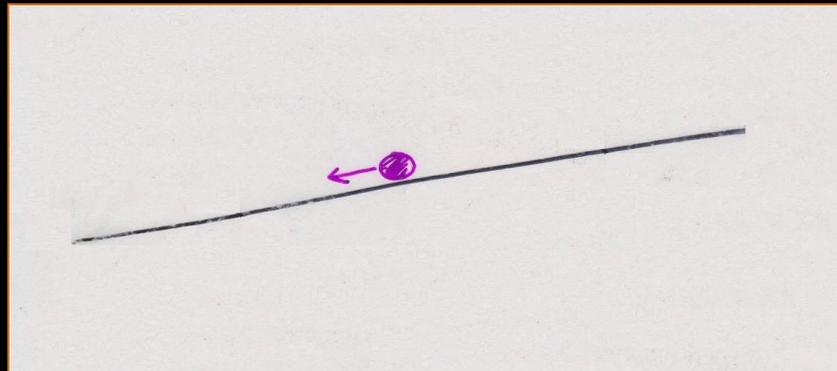
$$g_{mn}(x) = g_{mn}(\alpha^i; x)$$



# Context

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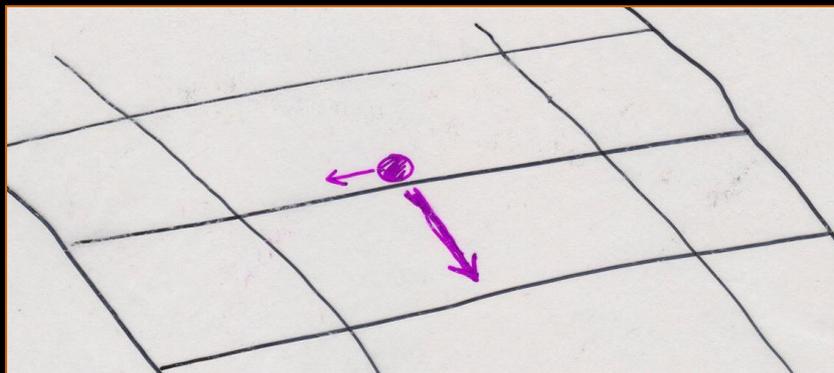


# Context

Superficially easy to find inflation in string theory:

It has many 4D scalars; supersymmetry can give these very flat potentials;

**BUT:** a convincing case for inflation requires reliably knowing the potential that stabilizes *all* moduli



# Context

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Giddings, Kachru & Polchinski  
Kachru, Kallosh, Linde & Trivedi

It has become possible to compute the low-energy potential for *all* moduli for some string constructions since around 2000.

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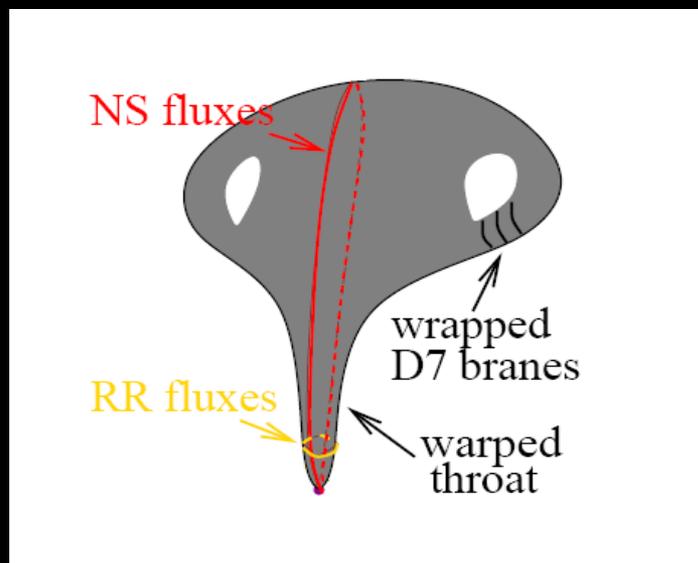
Giddings, Kachru & Polchinski  
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It has become possible to compute the low-energy potential for *all* moduli for some string constructions since around 2000. First done for Type IIB vacua

Key ingredients:

*D3 and D7 branes & orientifolds*

*Fluxes sourced by these branes, wrapping extra dimensions*



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# UV Issues

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- Naturalness (why are scalars light?)
- Trans-Planckian fields and tensor modes?
- Robustness
  - Multiple scalars and non-Gaussianity
  - Decoupling vs trans-Planckianism?
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# UV Issues: naturalness

---

It is unusual to find scalars that are systematically lighter than the other scales in the problem:

*ways out: supersymmetry, goldstone symmetry*

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*ways out: supersymmetry, goldstone symmetry*

Slow-roll condition  $M^2 V'' \ll V \sim \mu^4$  means inflation requires scalar is *extremely* light:  $m \ll H \sim \mu^2/M$

Unusually: sensitive even to ‘Planck slop’:

if  $L \supset V_0$  then what forbids  $L \supset V_0 \varphi^2/M^2 \sim H^2 \varphi^2$  ?

# UV Issues: naturalness

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So far as we can tell inflation seems no easier to get in string theory than for generic field theories

But we also tend only to seek inflation in a regime with strings well-described by a field theory anyway.

# UV Issues: naturalness

*Brandenberger & Vafa*

So far as we can tell inflation seems no easier to get in string theory than for generic field theories

But we also tend to  
with strings well-de

Some exceptions, like string-gas cosmology, but the good news (*you leave the field theory limit*) is also the bad news (*hard to calculate reliably because have left the field theory limit*)

# UV Issues: naturalness

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So far as we can tell inflation seems no easier to get in string theory than for generic field theories

Scalars can be light, usually by exploiting known mechanisms (like super- or Goldstone-symmetries)

de Sitter space tends to ruin supersymmetry, so often find inflaton as pseudo-Goldstone bosons

# UV Issues: naturalness

So far  
string

*SUSY often combines two types of pGB:*

$$X = e^\varphi + ia$$

*where:*

Scale  
these

- *a is an axion: in  $G/H$  with  $G$  compact*

$$V(a) = V_0 + V_1 \cos(a/f) + \dots$$

*with  $f \sim M_s \ll M_p$*

de S  
find

# UV Issues: naturalness

Freese, Frieman & Olinto

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*Inflation requires  $f \gg M_p$*

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-  *$\phi$  is a modulus: in  $G/H$  with  $G$  noncompact*

$$V(\varphi) = V_0 + V_1 e^{-\varphi/f} + \dots$$

*with  $f \sim M_p$*



# UV Issues: naturalness

BMQRZ th/0111025

Arises, eg, as radius of extra-dimensional cycle:  $r = \ell e^{\phi/f}$  with energy given as series in  $1/r$  given a kinetic term  $L = M^2 [R + (\partial r/r)^2]$

Calculable regime:  $r \gg \ell$  and so  $\phi \gg f$

of pGB:

compact

) + ...

de S  
find

-  $\phi$  is a modulus: in  $G/H$  with  $G$  noncompact

$$V(\phi) = V_0 + V_1 e^{-\phi/f} + \dots$$

with  $f \sim M_p$

# UV Issues: naturalness

Goncharov & Linde

*Inflation requires  $\varphi \gg f$  which is the same as the calculable regime; with no condition on size of  $f$ .*

*BUT inflation also requires  $V_0 \neq 0$*

*of pGB:*

*compact*

*) + ...*

de S  
find

*-  $\phi$  is a modulus: in  $G/H$  with  $G$  noncompact*

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# UV Issues: naturalness

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*Should take pGB potentials very seriously!!*

*In particular exponential potentials:*

$$V(\varphi) = V_0(1 - e^{-k\varphi} + \dots)$$

*so*

$$\epsilon = e^{-2k\varphi} \quad \text{and} \quad \eta = e^{-k\varphi}$$

*so slow roll is same as large field*

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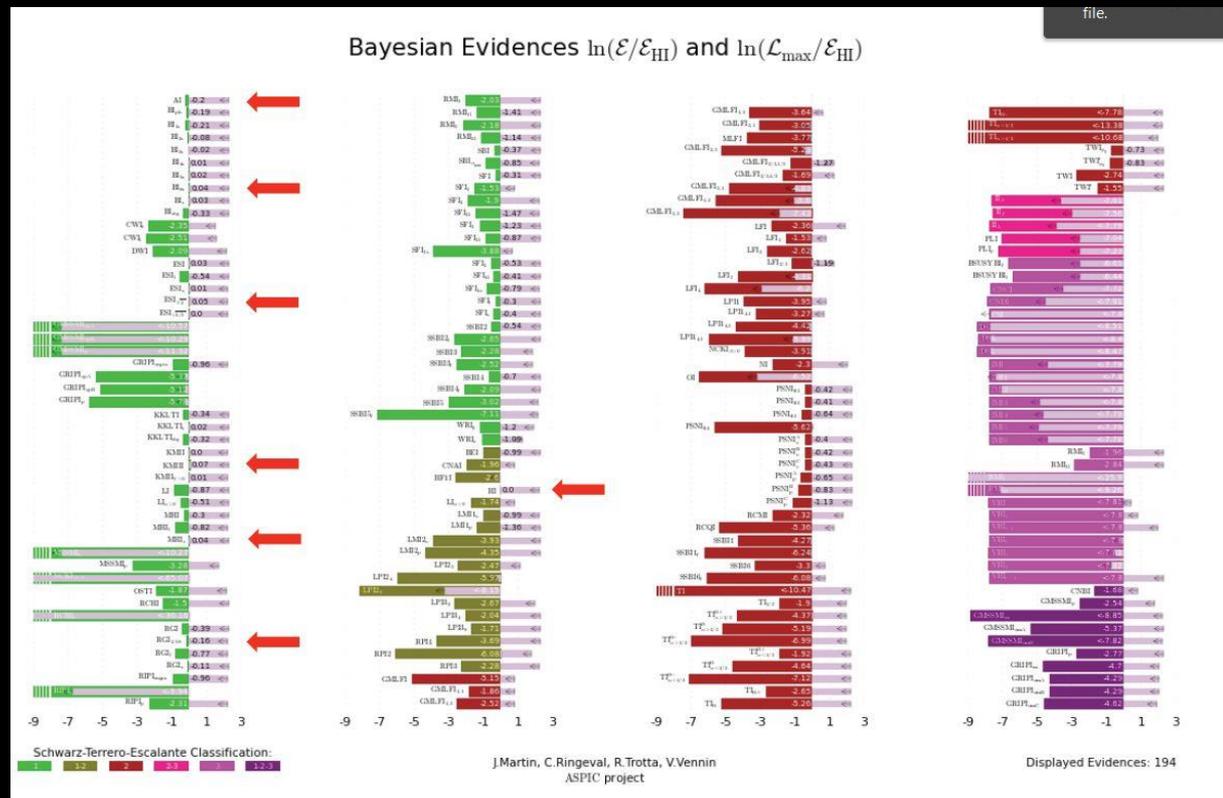
*so*

$$\epsilon = e^{-2k\varphi} \quad \text{and} \quad \eta = e^{-k\varphi}$$

*since  $\epsilon \sim \eta^2$  get prediction  $r \sim (n_s - 1)^2$*

# UV Issues: naturalness

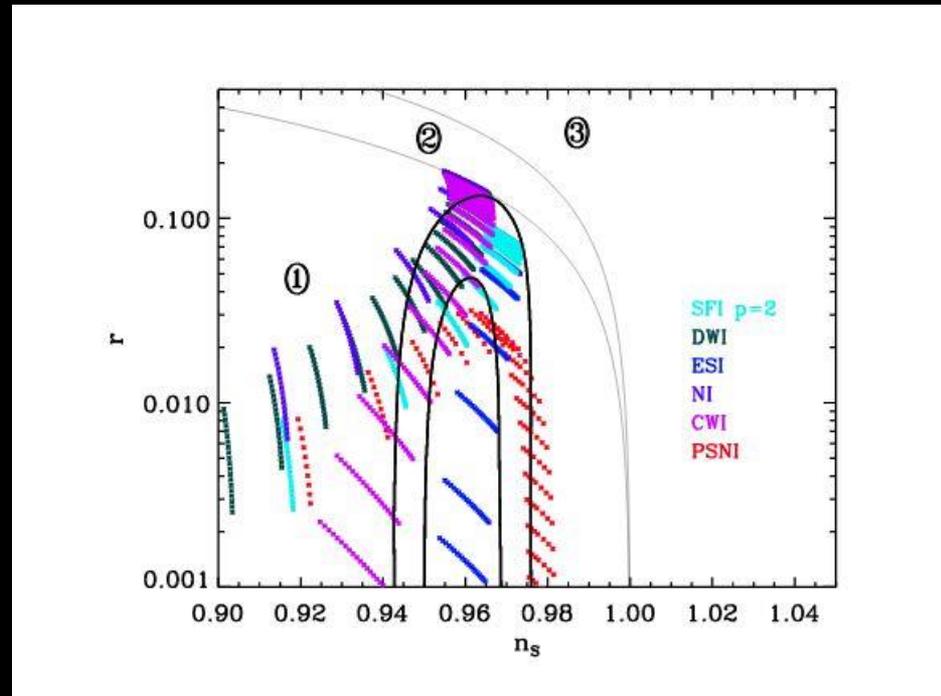
Martin, Ringeval & Vennin 2013



*Exponentials do well in comparison of models*

# UV Issues: naturalness

*Martin, Ringeval & Vennin 2013*



*Why some do better than others*

# UV Issues: naturalness

*Starobinsky  
Bezrukov & Shaposhnikov*

Do string vacua also give  $R^2$  or Higgs-type inflation?

$$L = M^2 R + \gamma R^2 + \frac{R^3}{M^2} + \dots$$

with  $M = M_p$  and  $\gamma \gg 1$ ?

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*Actually do get curvature expansions, but usually with  $M_s$  everywhere, so do not get unusually large  $R^2$  term.*



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*Actually do get curvature expansions, but usually with  $M_s$  everywhere, so do not get unusually large  $R^2$  term.*

*But Higgs and  $R^2$  inflation can also be regarded as exponential potentials, which seem easier to get.*

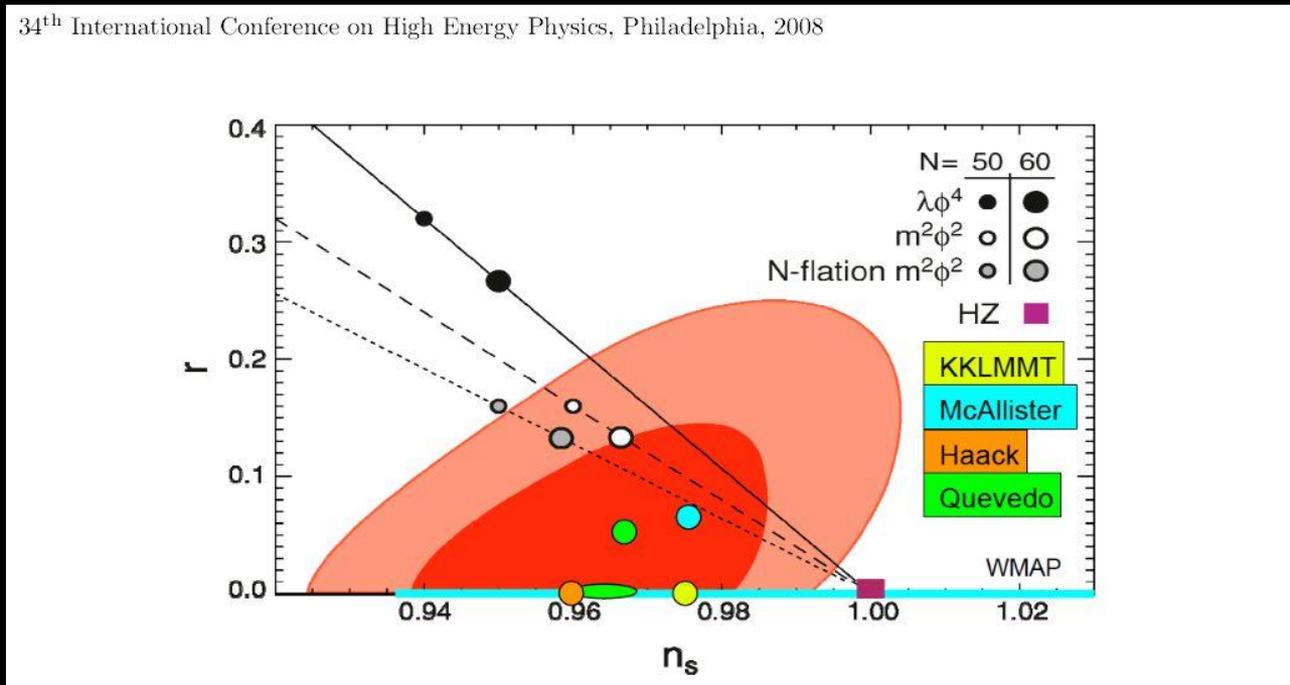
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# UV Issues: large fields and tensors

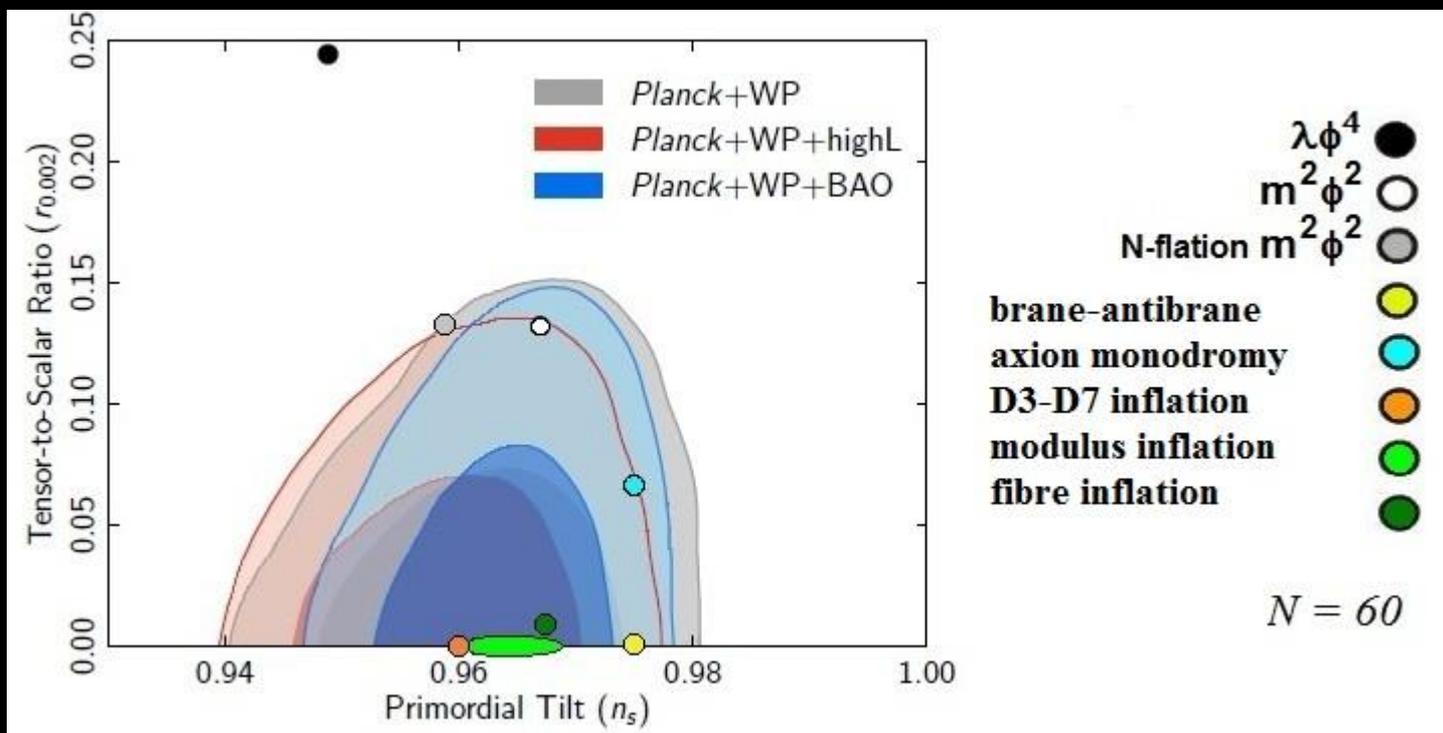
Early models tend to share a preference for small  $r$



*J. Polchinski ICHEP 08 summary talk*

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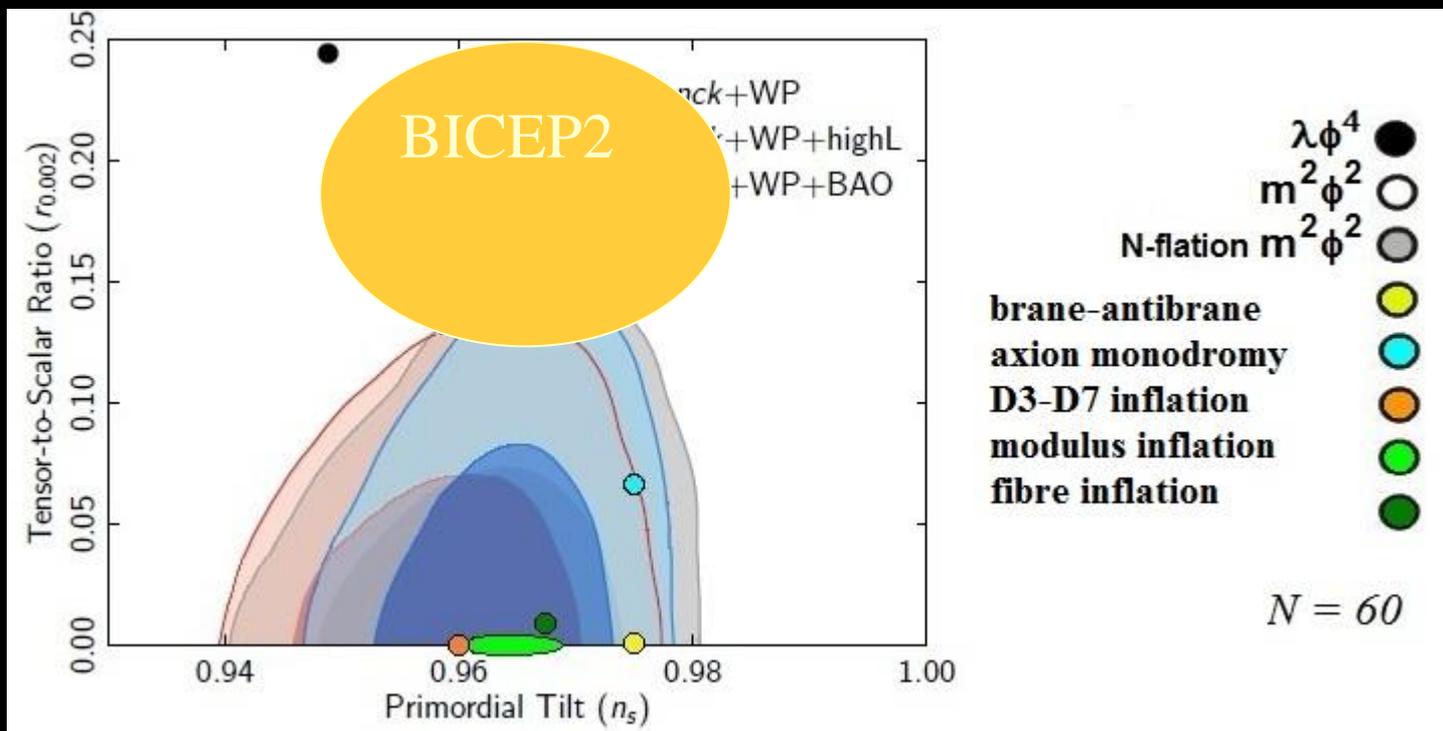
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# UV Issues: large fields and tensors

Early models tend to share a preference for small  $r$



*J. Polchinski ICHEP 08 summary talk*

# UV Issues: large fields and tensors

Lyth

Early models tend



*Usually large  $r$  corresponds to large excursions in field space*

$$\Delta\phi > M_p (r/4\pi)^{1/2}$$

*In examples these turn out to require things like branes rolling further than extra dimensions are large.*

*J. Polchinski ICHEP 08 summary talk*

# UV Issues: large fields and tensors

Early models tend



*Clearly large fields are not in themselves impossible to get, as the exponential potentials show:*

$$r = \ell e^{\Phi/f}$$

*These start life as flat directions so large fields cost little energy (and so make sense in an EFT).*

*But how to inflate?  $V = V_0 + \dots$*

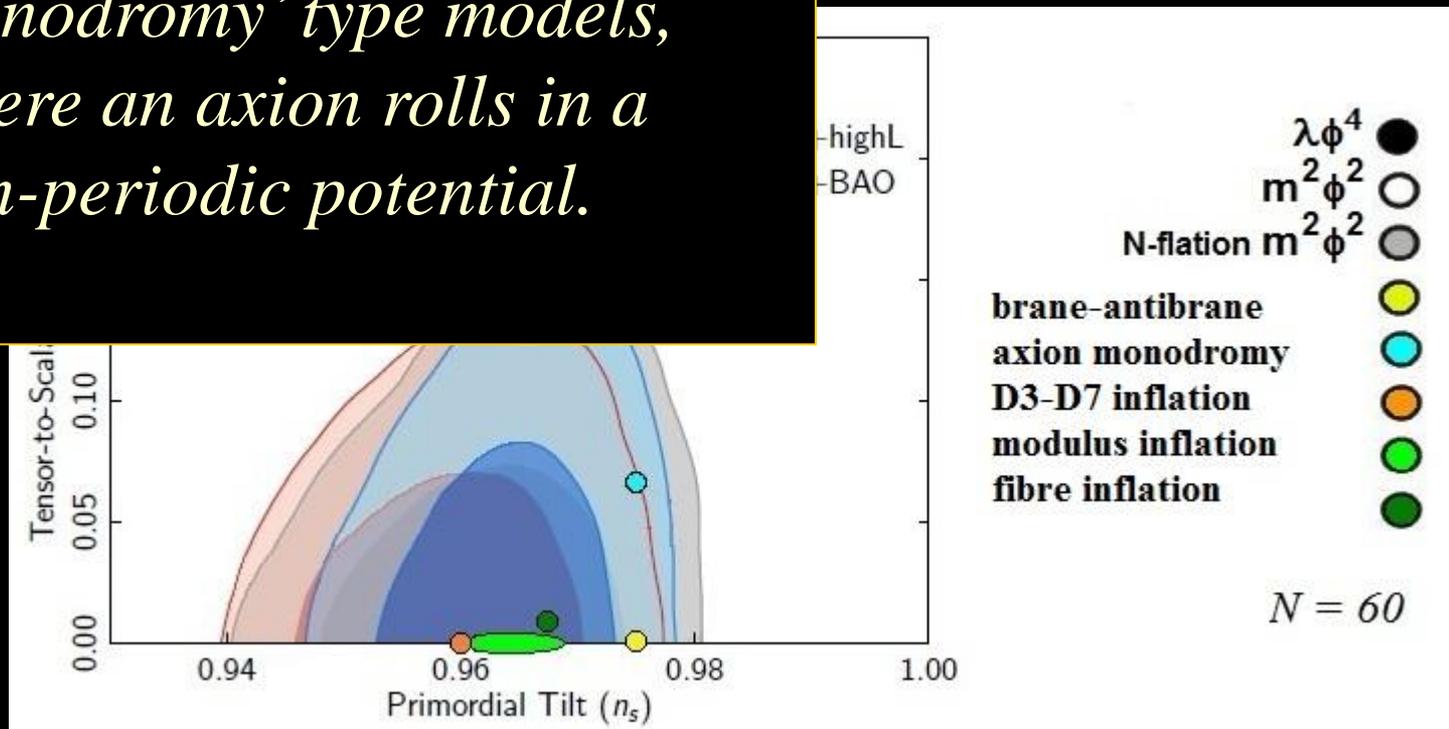
*J. Polchinski ICHEP 08 summary talk*

# UV Issues: large fields and tensors

Silverstein & Westphal  
McAllister, Silverstein & Westphal  
Conlon

*First counter-examples were 'monodromy' type models, where an axion rolls in a non-periodic potential.*

reference for small  $r$



*J. Polchinski ICHEP 08 summary talk*



# UV Issues: large fields and tensors

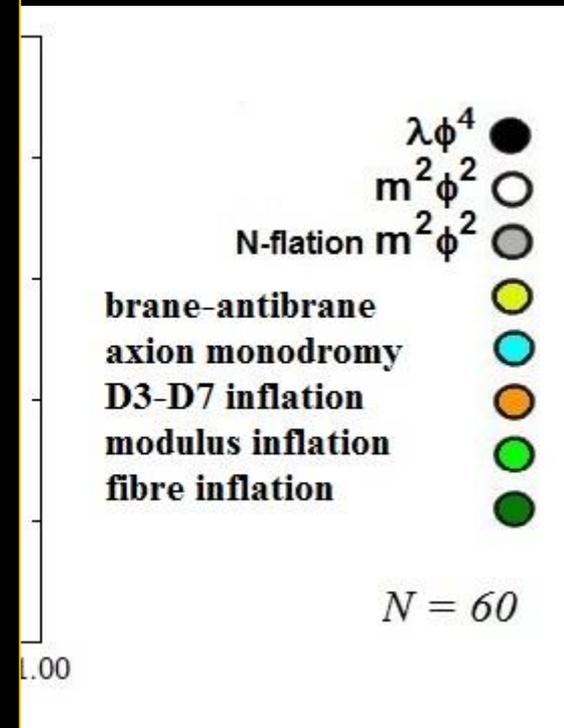
Kim, Niilles & Peloso  
Kappl, Krippendorf & Nilles

*Field-theoretic mechanisms exist for getting larger  $f$  (and so also larger tensors), such as 'alignment' mechanisms.*

$$V(\varphi_1, \varphi_2) = U_1(a\varphi_1 + b\varphi_2) + U_1(c\varphi_1 + d\varphi_2)$$

*By ensuring  $f > M_p$  these can reproduce predictions of simple  $m^2\varphi^2$  models because  $\varphi/f \ll 1$*

ence for small  $r$



*J. Polchinski ICHEP 08 summary talk*

# UV Issues

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# UV Issues: robustness

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Brandenberger & Martin

Do observers need to worry about any of this?

- *If inflation stretches scales, why doesn't it also stretch unknown trans-Planckian physics to observable scales?*

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*String physics in string theory seems to decouple...*

# UV Issues: robustness

McAllister et al

Do observers need to worry about any of this?

- *If inflation stretches scales, why doesn't it also stretch unknown trans-Planckian physics to observable scales?*

String models often have many light scalars if they have one: should we generically expect nongaussianity and isocurvature perturbations?

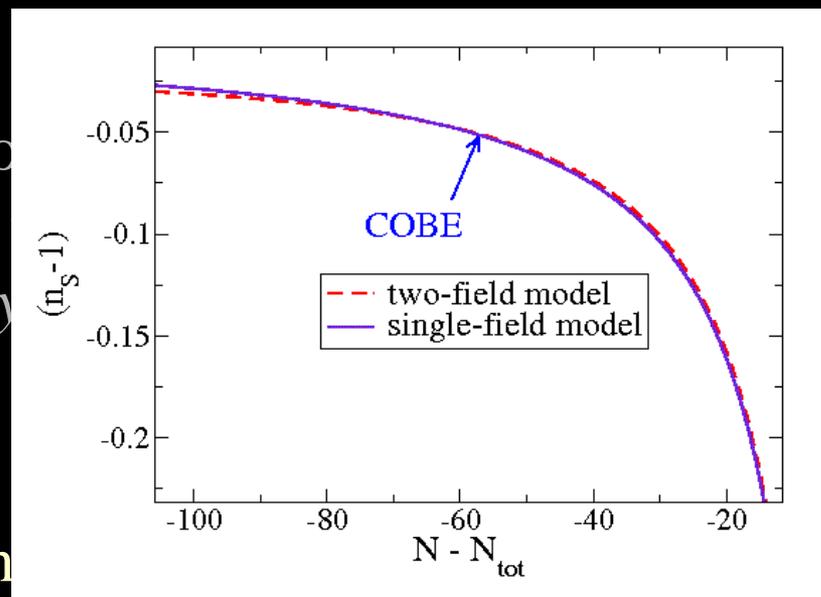
# UV Issues: robustness

'Racetrack 8'

Do observers need to worry about  
- *If inflation stretches scales, unknown trans-Planckian physical scales?*

String models often have many fields  
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*Multi-field models usually are moving along target-space geodesics near horizon exit so are usually effectively single field models...*



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Does string theory have anything useful to say about reheating?

*in principle it should eventually, once we have a single model that both inflates and contains the Standard Model....*



# UV Issues: reheating

Sarangi & Tye  
Copeland, Myers & Polchinski

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*In principle it should eventually, once we have a single model that both inflates and contains the Standard Model....*

*Can bring surprises like post-inflationary cosmic string production, ...*

# UV Issues: reheating

Barnaby, CB & Cline,

Does string theory have anything to say about reheating?



strings already show why studies of a single inflaton oscillating while coupled to SM fields may be inadequate: *a good furnace is not adequate to reheat if you do not also have good insulation. Must know ALL the degrees of freedom at the relevant energies*

# UV Issues

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# UV Issues: mind-broadening

Some novel features to string-motivated models:

CB, Majumdar, Nolte, Rajesh, Quevedo & Zhang

*Inflaton field might not even exist in our present-day world (such as if it describes the distance between a brane-antibrane pair that annihilate at the end of inflation)*

Silverstein & Tong,

*Slow roll need not be inconsistent with relativistic kinematics (such as in DBI inflation)*

$$\mathcal{L}_{\text{DBI}} = -f(\phi)^{-1} \sqrt{1 - 2f(\phi)g^{\mu\nu}\partial_\mu\phi\partial_\nu\phi} + f(\phi)^{-1} - V(\phi)$$

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# Scorecard

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- Mechanisms
- Robustness
- Initial Conditions
- Naturalness
- Reheating & scales
- Signatures
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- Hopefully even more novel physics will emerge...

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# The Cosmological Constant?



*fin*

