

# Giant planets vs Brown Dwarfs: who's who ?

Gilles Chabrier

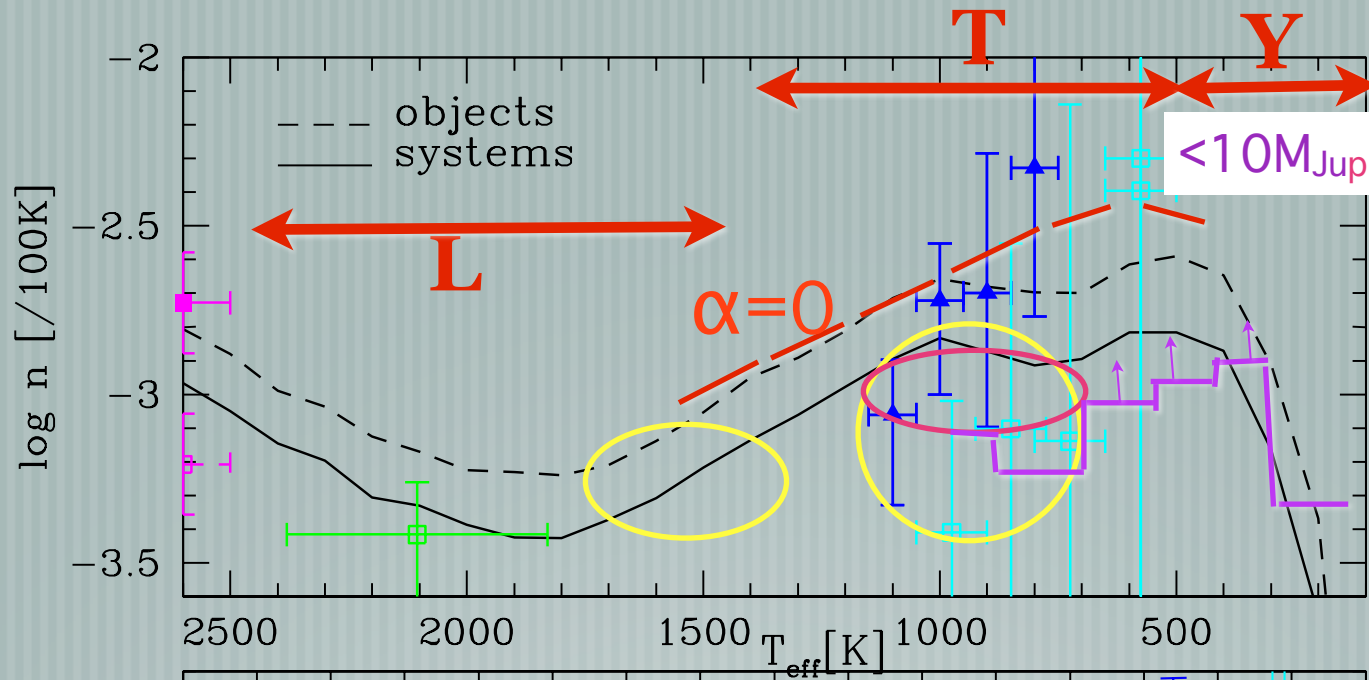


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# I) THE MASS FUNCTION

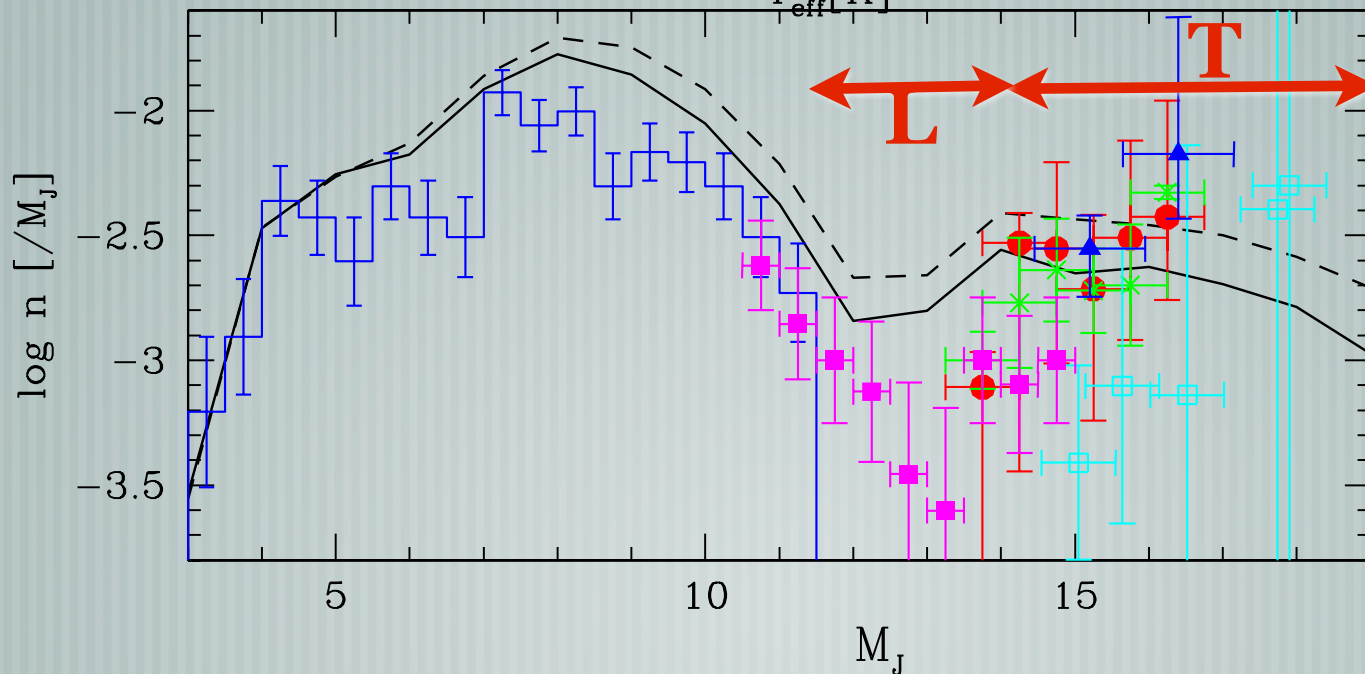
# Brown dwarf mass function

$n(T_{\text{eff}})$



Reid et al. '04  
+ Cruz et al. '07  
Burgasser '04  
Burningham et al. '10  
Gizis et al. '00  
Metchev et al. '08  
Reylé et al. '10  
Kirkpatrick et al. '12

$n(M_J)$



Reylé et al. '10 CFHBS  
Allen et al. '05  
Reid et al. '04  
+ Cruz et al. '07  
Burgasser '04  
Burningham et al. '10

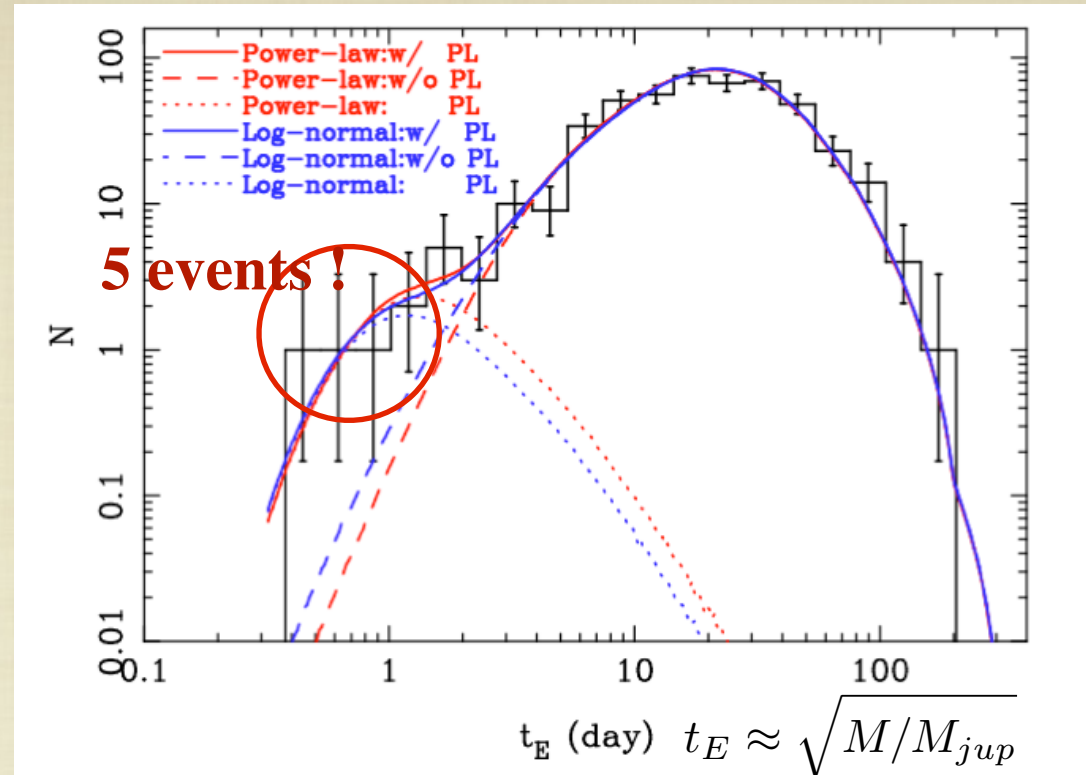
<i>Parameter</i>		<i>Disk</i>	<i>Spheroid</i>	<i>Dark halo</i>
$n_{BD}$		$2.6 \times 10^{-2}$	$3.5 \times 10^{-5}$	
$\rho_{BD}$		$1.0 \times 10^{-3}$	$\lesssim 2.3 \times 10^{-6}$	
$n_*$		$(9.3 \pm 2) \times 10^{-2}$	$\leq (2.4 \pm 0.1) \times 10^{-4}$	
$\rho_*$		$(3.4 \pm 0.3) \times 10^{-2}$	$\leq (6.6 \pm 0.7) \times 10^{-5}$	$\ll 10^{-5}$
$n_{rem}$		$(0.7 \pm 0.1) \times 10^{-2}$	$\leq (2.7 \pm 1.2) \times 10^{-5}$	
$\rho_{rem}$		$(0.6 \pm 0.1) \times 10^{-2}$	$\leq (1.8 \pm 0.8) \times 10^{-5}$	$< 10^{-4}$
$n_{tot}$		$0.13 \pm 0.03$	$\leq 3.0 \times 10^{-4}$	
$\rho_{tot}$		$(4.1 \pm 0.3) \times 10^{-2}$	$\leq (9.4 \pm 1.0) \times 10^{-5}$	$< 10^{-4}$
BD:	$\mathcal{N}; \mathcal{M}$	0.20; 0.02	0.10; 0.03	
LMS( $\leq 1 M_\odot$ ):	$\mathcal{N}; \mathcal{M}$	0.71; 0.68	0.80; 0.77	
IMS(1-9 $M_\odot$ ):	$\mathcal{N}; \mathcal{M}$	0.03; 0.15	0.; 0.	
WD+NS:	$\mathcal{N}; \mathcal{M}$	0.06; 0.15	0.10; 0.20	

<sup>a</sup>The number densities  $n$  are in [ $\text{pc}^{-3}$ ], the mass densities  $\rho$  are in [ $M_\odot \text{pc}^{-3}$ ].

$$N_{BD}/N_* = 1/4 - 1/3$$

$$\text{WISE (now !)} \sim 1/5$$

# Excess of very-low-mass objects ? Sumi et al. '11



- 5 « anomalous » events: no statistical significance

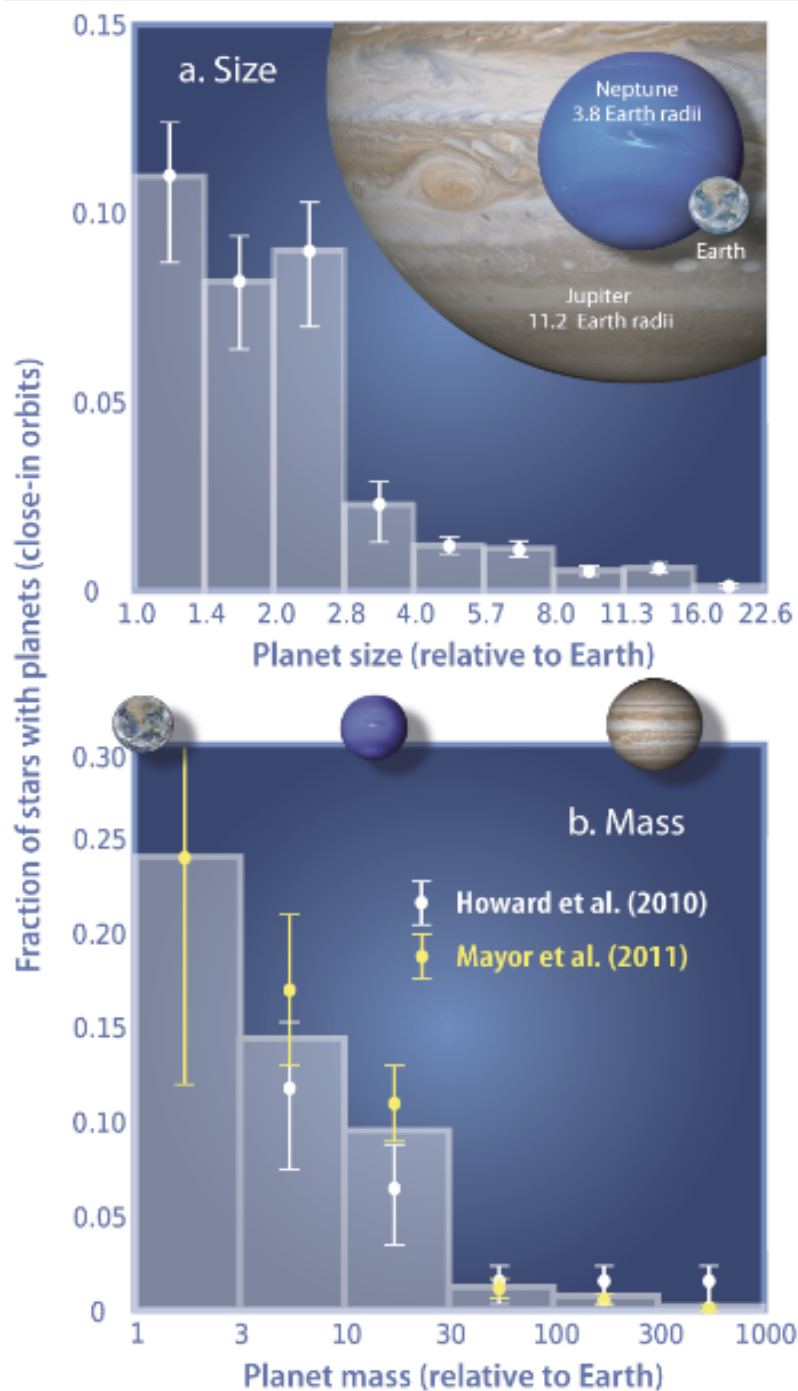
- $n_{jup} \sim 2 \times n_{\star}$   $dN/dM \propto M^{-1.3}$  ( $M < 10 M_{jup}$ )

excluded by observations in SFRs (Scholz et al '12, Caballero et al. '07,...). See also Quanz et al. '12

- when observing a short duration isolated event with low S/N and non-complete coverage, the symmetry of the event, a characteristic property of lensing events, is very difficult to assess properly, in contrast to the case when the planetary event occurs as a caustic during a stellar lensing event.
- **Free-floating planet candidate events are thus much more uncertain than star-planet candidate events and should be considered with extreme caution.**

# Planet mass function

Mayor et al. (2011),  
Howard et al. (2013)



Planets **INCREASE**  
with decreasing size/mass

Brown dwarfs **DECREASE**  
with decreasing mass

in the same mass range

- Expected from core-accretion
- GP Mass dist'n  $\neq$  BD Mass dist'n

## 2) THE FORMATION

# GRAVOTURBULENT FRAGMENTATION

Padoan & Nordlund; Hennebelle & Chabrier; Hopkins  
(see Chabrier, Johansen, Janson & Rafikov 2014 for more details)

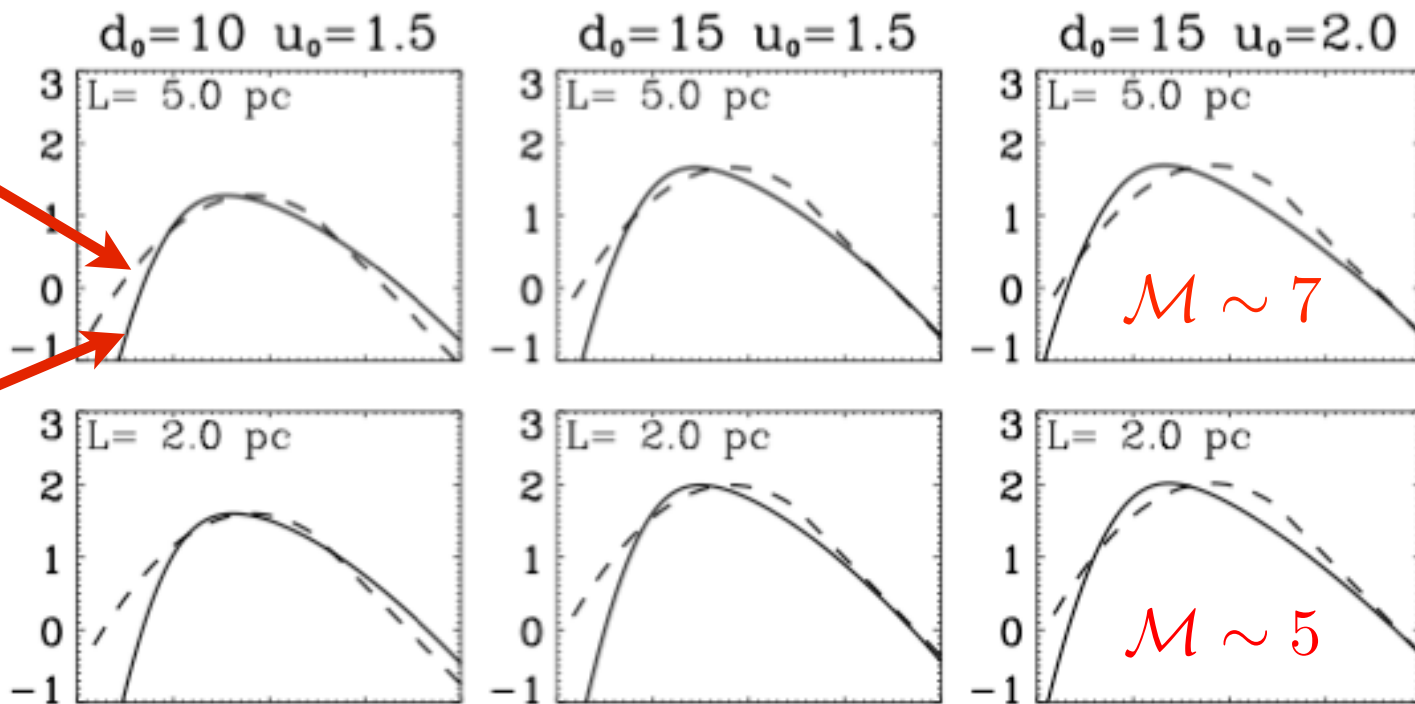
- Large-scale **turbulence** (density PDF lognormal) sets up a lognormal distribution of overdense regions at all scales in the cloud
- Virial: regions  $M(R)$  with  $|E_{\text{grav}}(R)| > (E_{\text{th}} + E_{\text{rms}}(R) + E_{\text{mag}})$  **grav. collapse**  
-> leads to the prestellar **core** mass function (**CMF**)
- magneto-centrifugally outflows expel  $\sim 30\text{-}50\%$  of the mass  
-> leads to the **star** mass function (**IMF**)
- **turbulence sets up the density fluctuations**  
at the very early stages of star formation
- the IMF is imprinted in the very cloud conditions ( $\langle n \rangle$ ,  $T$ , Mach)



$$\bar{n} = (d_0 \times 10^3 \text{ cm}^{-3}) \left( \frac{L}{1 \text{ pc}} \right)^{-0.7}, \quad V_{\text{rms}} = (u_0 \times 0.8 \text{ km s}^{-1}) \left( \frac{L}{1 \text{ pc}} \right)^\eta.$$

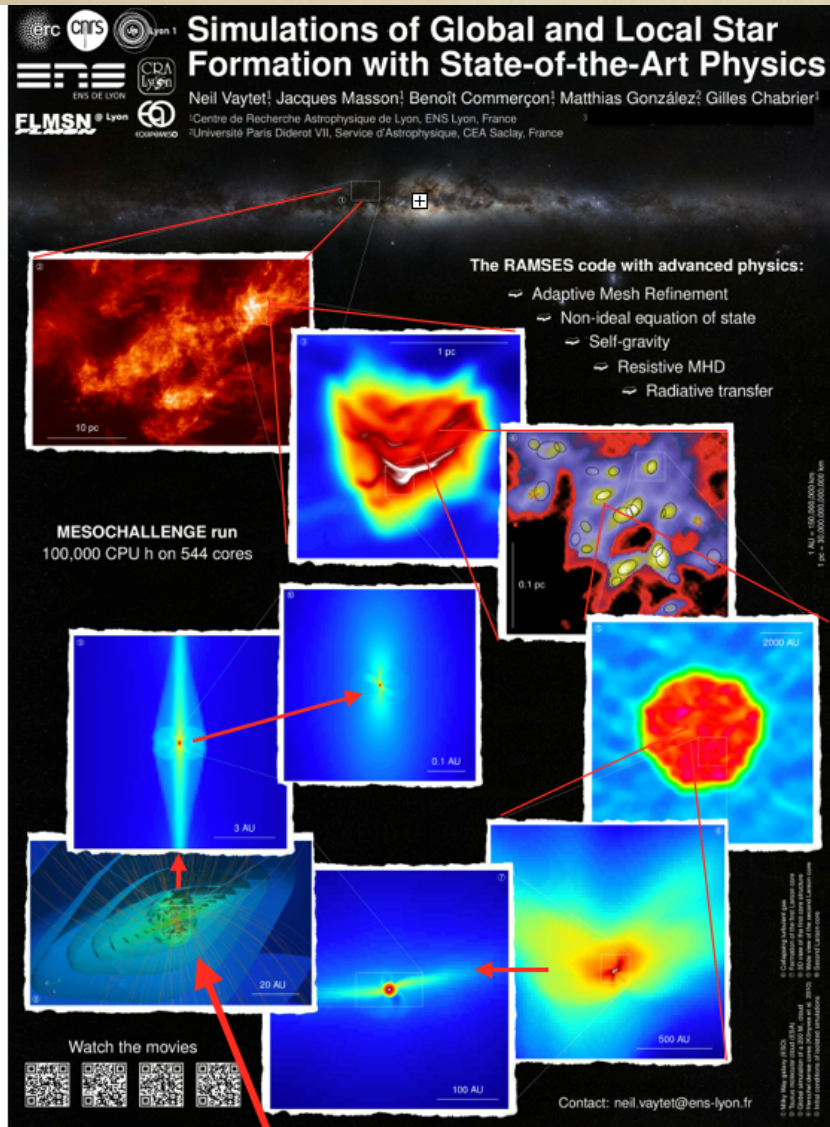
Chabrier  
system IMF

HC analytical  
IMF



Hennebelle & Chabrier '09, '13

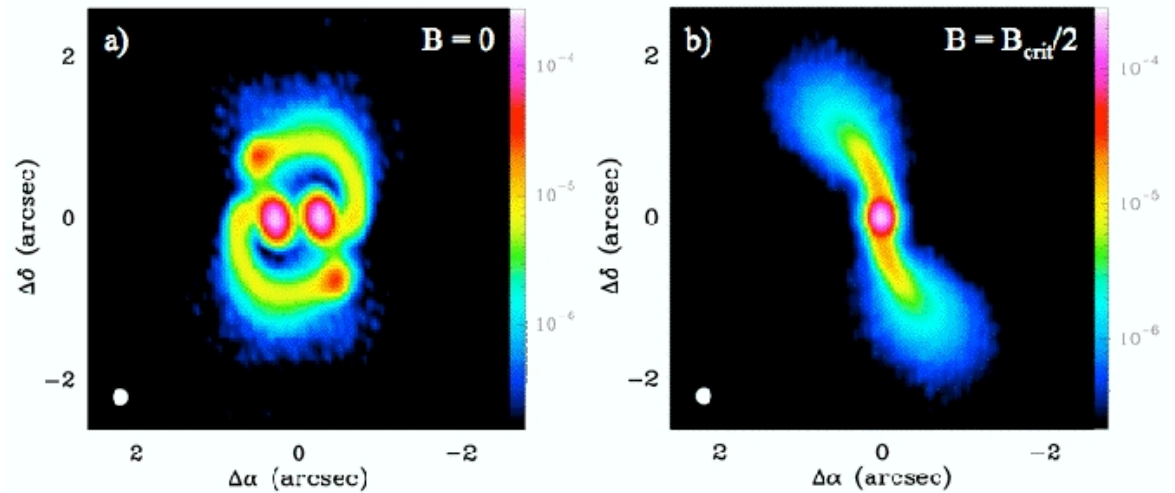
# Star/Brown dwarf formation



Can we test these ideas ?

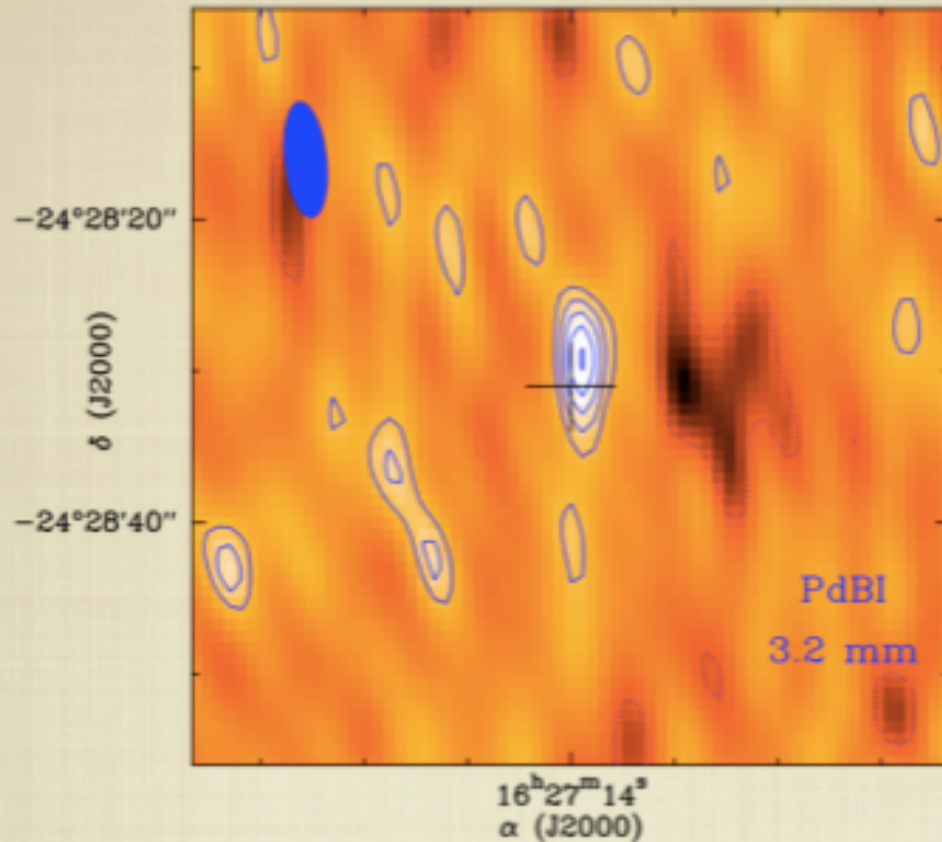
ALMA will offer the spatial resolution required

Synthetic observations done with the ALMA simulator  
 Included in the Gildas software



Commerçon et al. 2012

# Observed **isolated Pre-Brown dwarf core**



Oph B-11  
 $M \sim 30 M_{\text{Jup}}$   
 $R < 460 \text{ AU}$   
 $n \sim 10^7 - 10^8 \text{ cm}^{-3}$

André et al., 2012, Science 337, 69

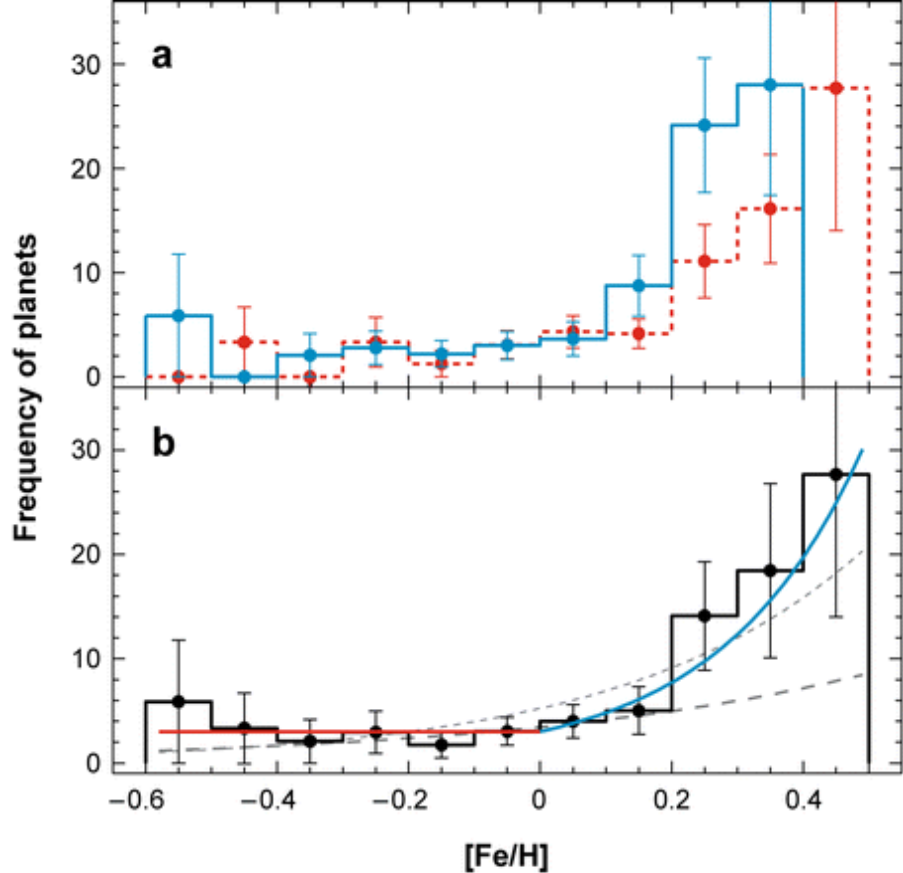
emerging observations of **isolated proto-brown dwarfs**

**VeLLO L1148-IRS** (Kauffmann et al. '11); **IRAS 16253-2429** (Wiseman et al.)

**J042118 + J041757** 1-5  $M_{\text{Jup}}$  (Palau et al. '12); **L328-IRS** <50  $M_{\text{Jup}}$  (C-W. Lee et al. '13)

**Growing theoretical and observational evidence that BDs dominantly form like stars**  
(Luhman, ARA&A 2012, Chabrier et al., PPVI review 2014)

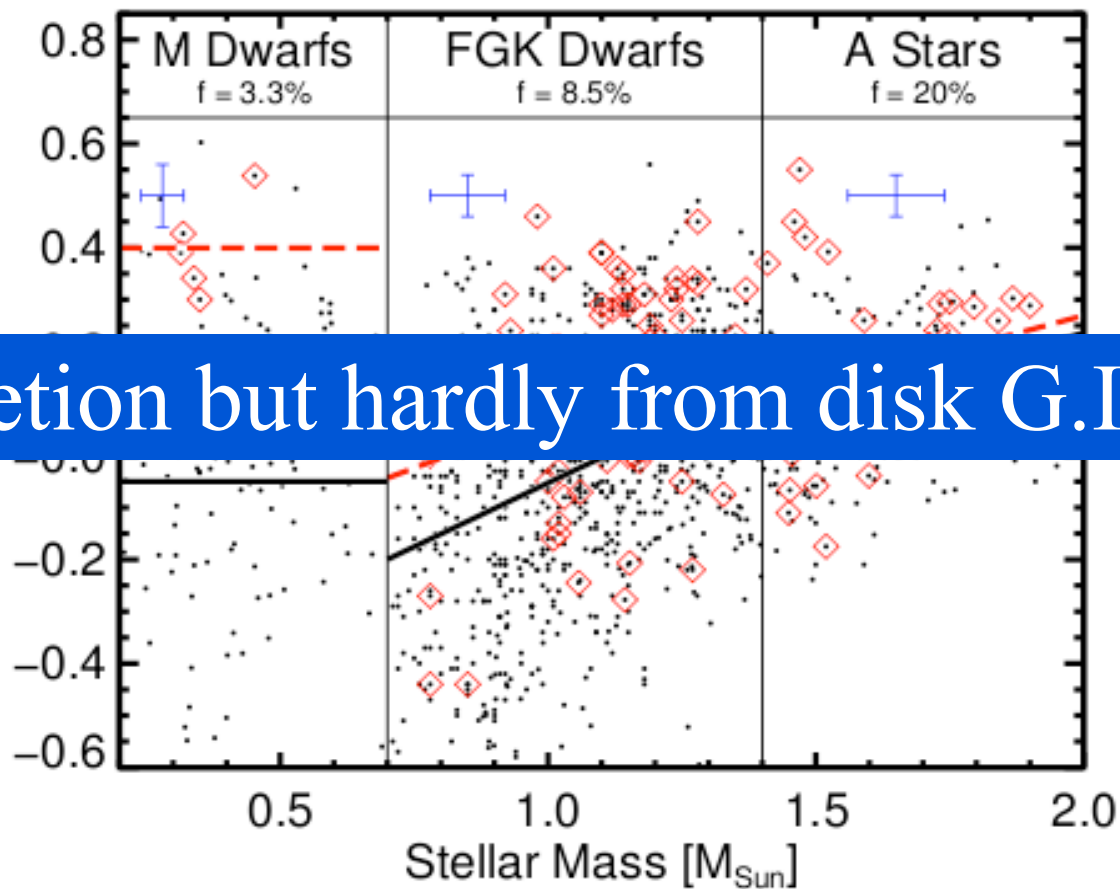
# Metallicity



Giant planet frequency is strongly dependent on stellar host metallicity

AR Udry S, Santos NC. 2007.  
Annu. Rev. Astron. Astrophys. 45:397–439

Johnson et al. (2010)

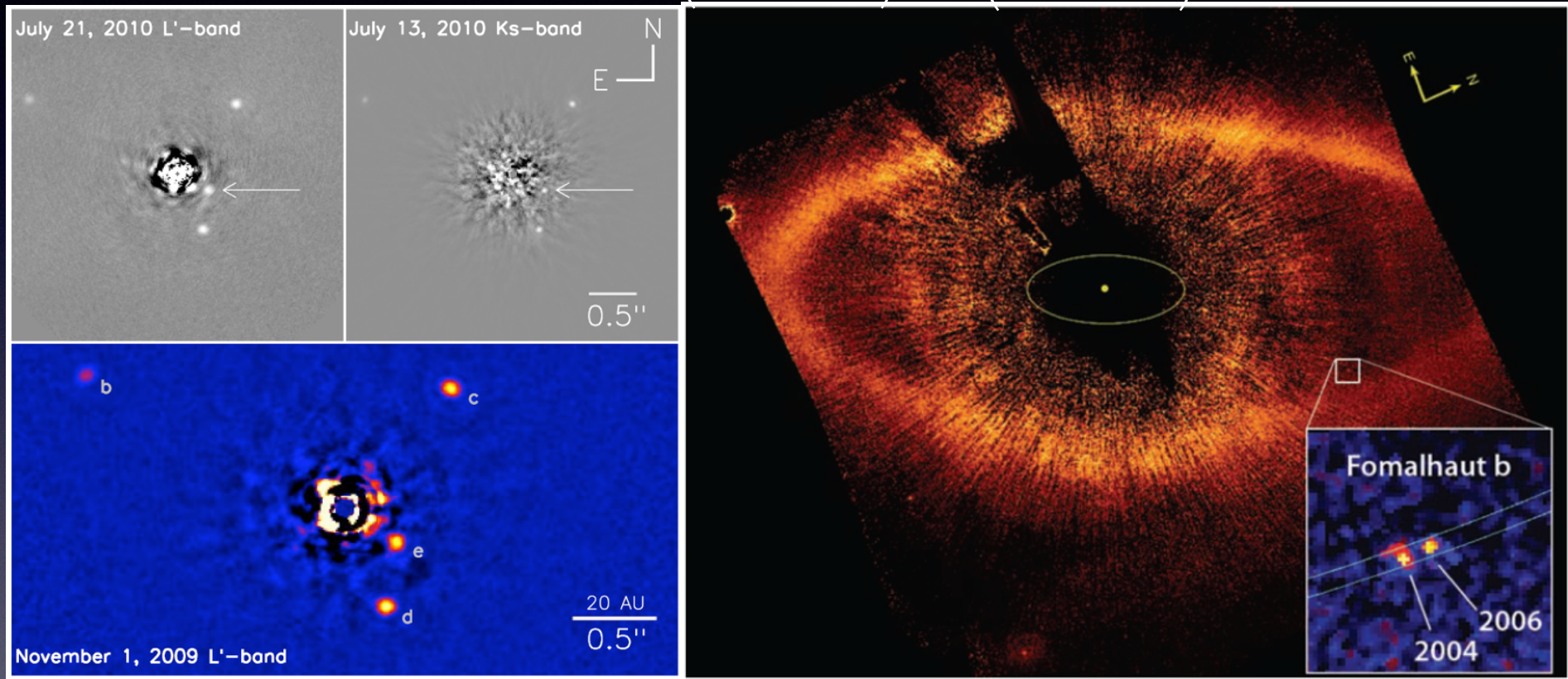


Expected from core-accretion but hardly from disk G.I

Correlation is weaker for lower-mass planets but stronger for lower-mass stars

# PROBLEM FOR C.A. : PLANETS ON WIDE ORBITS

$$a < 44 \text{ AU} \kappa_{0.1}^{-0.09} \left( \frac{\tau_{neb}}{3 \text{ Myrs}} \right)^{0.75} \left( \frac{\Sigma}{\Sigma_{MMSN}} \right)^{2/3} \quad (\text{Rafikov 2010})$$



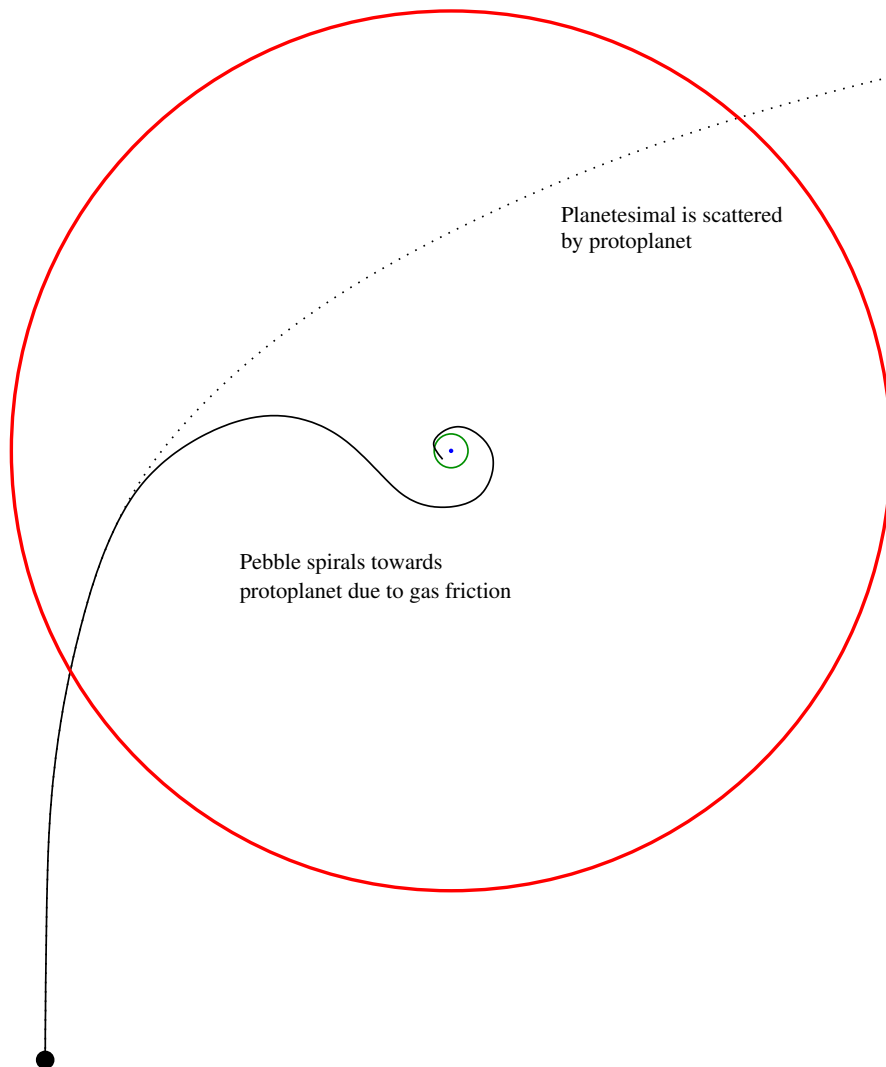
HR 8799 (4 planets at 14.5, 24, 38, 68 AU)

Fomalhaut (1 controversial planet at 113 AU)

*No way to form the cores of these planets within the lifetime of the protoplanetary disc by standard core accretion*

# Pebble accretion

- Most planetesimals are simply scattered by the protoplanet
- Pebbles spiral in towards the protoplanet due to gas friction
- Pebbles are accreted from the entire Hill sphere
- Growth rate by planetesimal accretion :



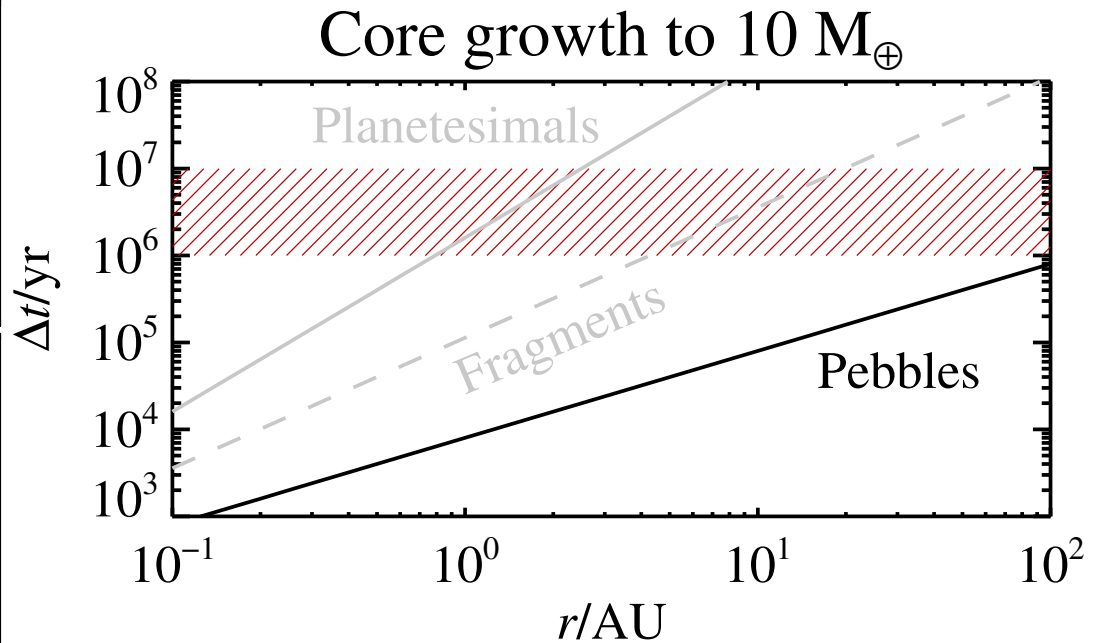
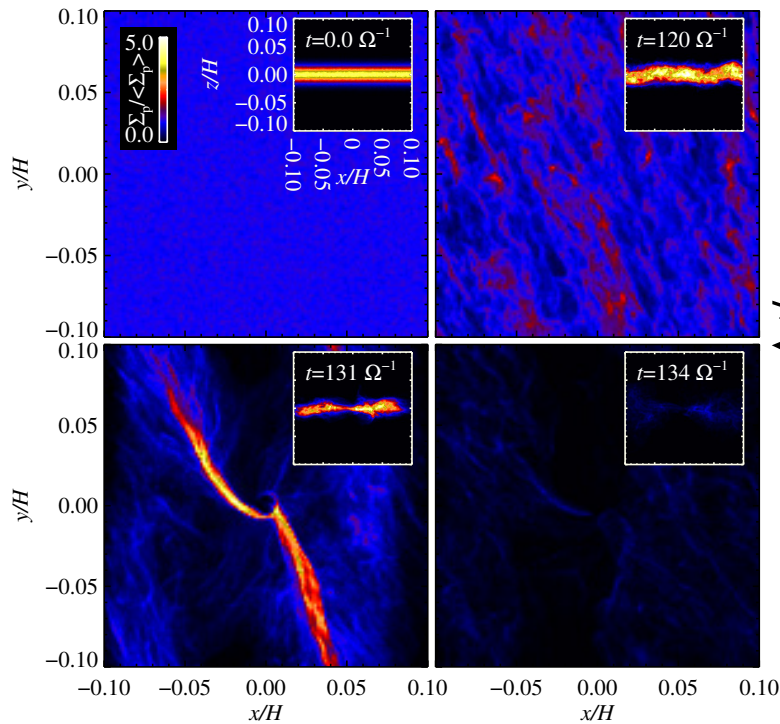
$$\dot{M} = \alpha R_H^2 \mathcal{F}_H$$

$$\frac{R_P}{R_H} \equiv \alpha \approx 0.001 \left( \frac{r}{5 \text{ AU}} \right)^{-1}$$

- Growth rate by pebble accretion :

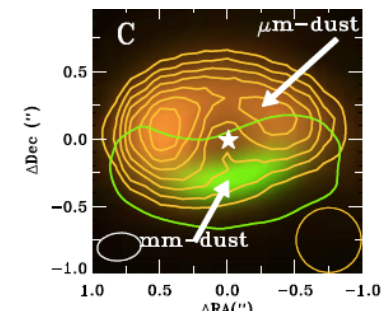
$$\dot{M} = R_H^2 \mathcal{F}_H$$

# Time-scale of pebble accretion



- **Pebble accretion speeds up core formation** by a factor **1,000** at **5 AU** and **10,000** at **50 AU** (Lambrechts & Johansen 2012; Ormel & Klahr 2010; Morbidelli & Nesvorny 2012)
- **Cores form well within the lifetime of the protoplanetary gas disc, even at large orbital distances**
- Requires large planetesimal seeds, consistent with turbulence-aided planetesimal formation (see e.g. A. Johansen et al. 2014, PPVI)

Observations of a «dust trap» in a disk between 45-90 AU  
(Van der Marel et al., Science '13)



- **HD100546 b** (Quanz et al. 2011)

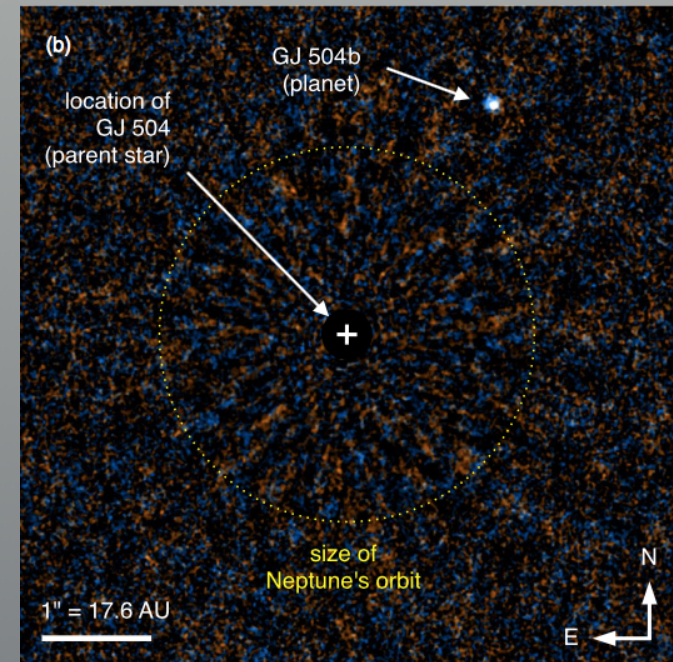
$M_{\star}=2.4 M_{\text{sol}}$   $M_p \sim 5-12 M_{\text{jup}}$  @ 50 AU

Disk not massive enough ( $\sim 0.4\% M_{\star}$ ) for GI.  
 Could be explained by cm-size pebble acc'n

- **GJ504b** (Kuzuhar et al. 2013)

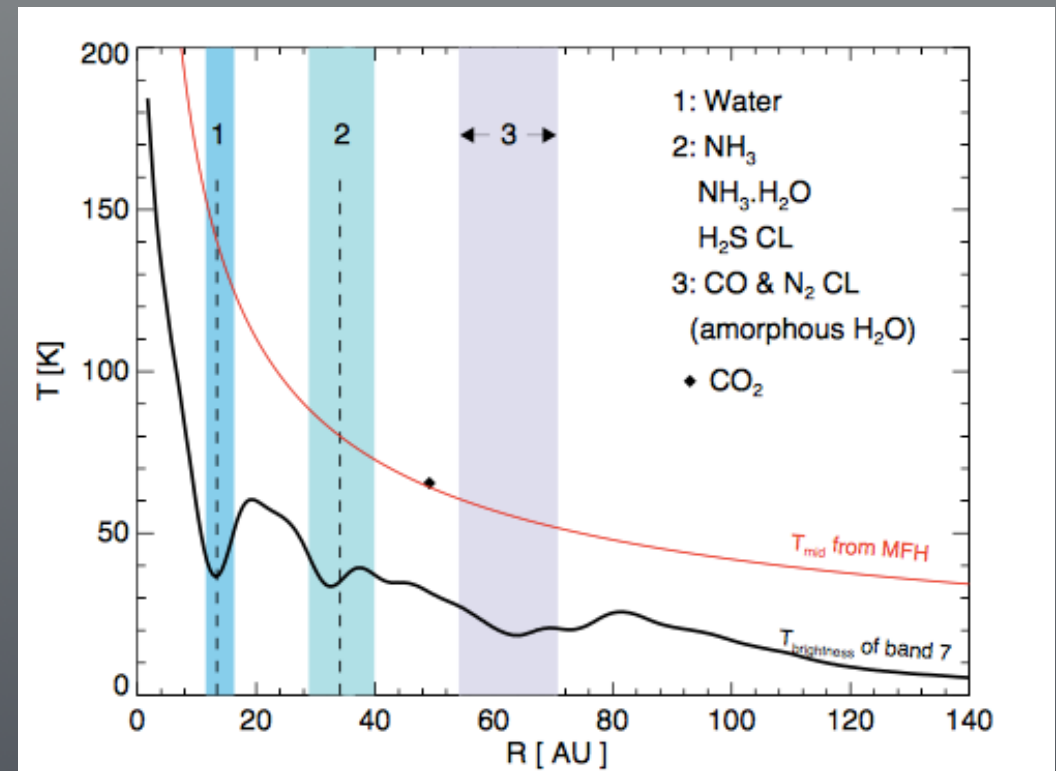
$M_{\star}=1.2 M_{\text{sol}}$ ,  $[M/H]=0.28$   $M_p \sim 1-8 M_{\text{jup}}$  @ 44 AU

Could be explained by cm-size pebble acc'n



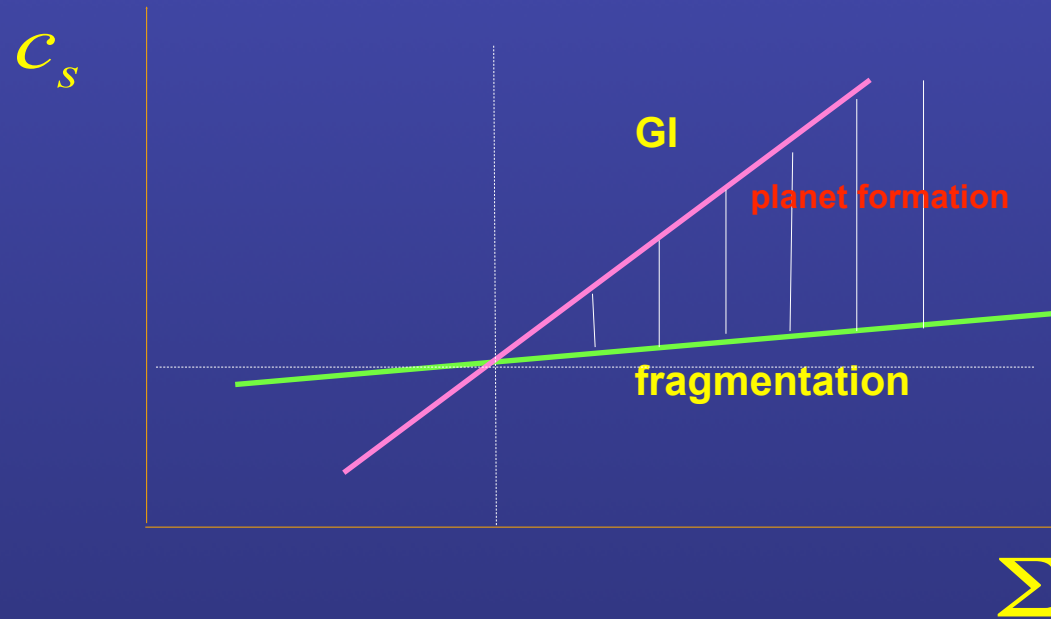
- Zhang et al. 2015

13 AU dip: cond'n H<sub>2</sub>O  
 32 AU dip: cond'n NH<sub>3</sub>+hydrates  
 dcm-size pebbles (Ros & Johansen '13)  
 -> >km-size planetesimals by streaming instabilities (Johansen et al. 2014)





# Gravitational Instability in a disk: Thermodynamical constraints



As a result, **giant planet formation by GI requires**

$$\Sigma > \left[ \frac{\Omega^7}{3\sigma (\pi G)^6} \left( \frac{k}{\mu} \right)^{\frac{4}{j}} \right]^{1/5} = 3 \times 10^5 \text{ g cm}^{-2} a_{AU}^{-21/10} \quad (\sim 100 \text{ MMSN}) !$$

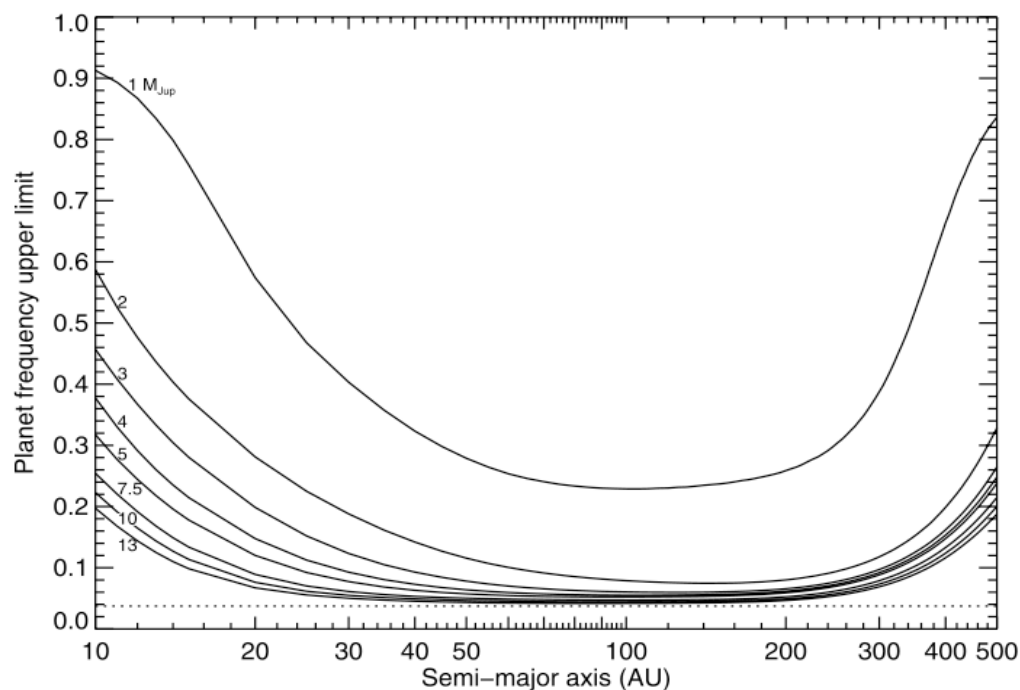
$$T > \left[ \frac{\Omega^2}{3\pi\sigma G} \left( \frac{k}{\mu} \right)^{\frac{3/2}{j}} \right]^{2/5} = 2200 \text{ K } a_{AU}^{-6/5} \quad (\sim T_{Sun}) !!!$$

## Requires rather extreme disk conditions

- Gravitational coupling to the massive disk **quickly migrates** them into the star (Vorobyov & Basu 2005, Machida 2011, Baruteau et al. 2013)

# Statistical constraints from Direct Imaging (with caveats!)

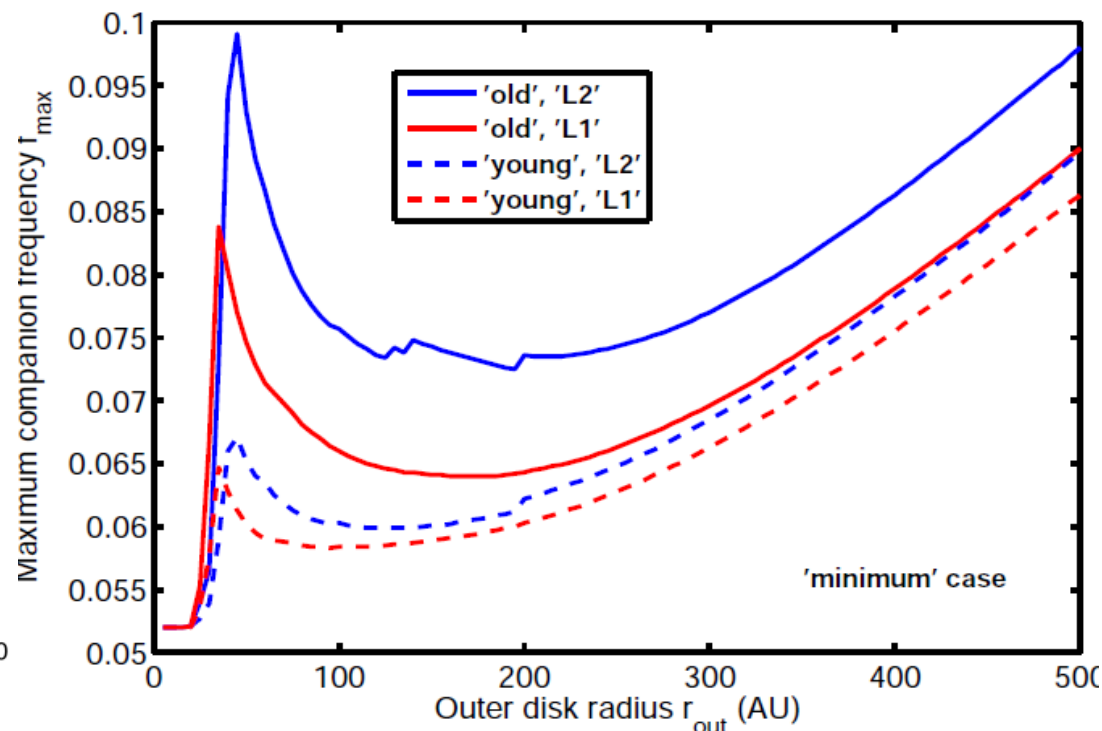
apply both to BD's and planets !



Lafreniere et al. (2007)

<23% of stars have >2 M<sub>J</sub> planets at 25-450 AU

<9% of stars have >5 M<sub>J</sub> planets at 25-450 AU



Janson et al. (2012)

<10% of stars host ~Jupiter-mass objects formed by disk instability

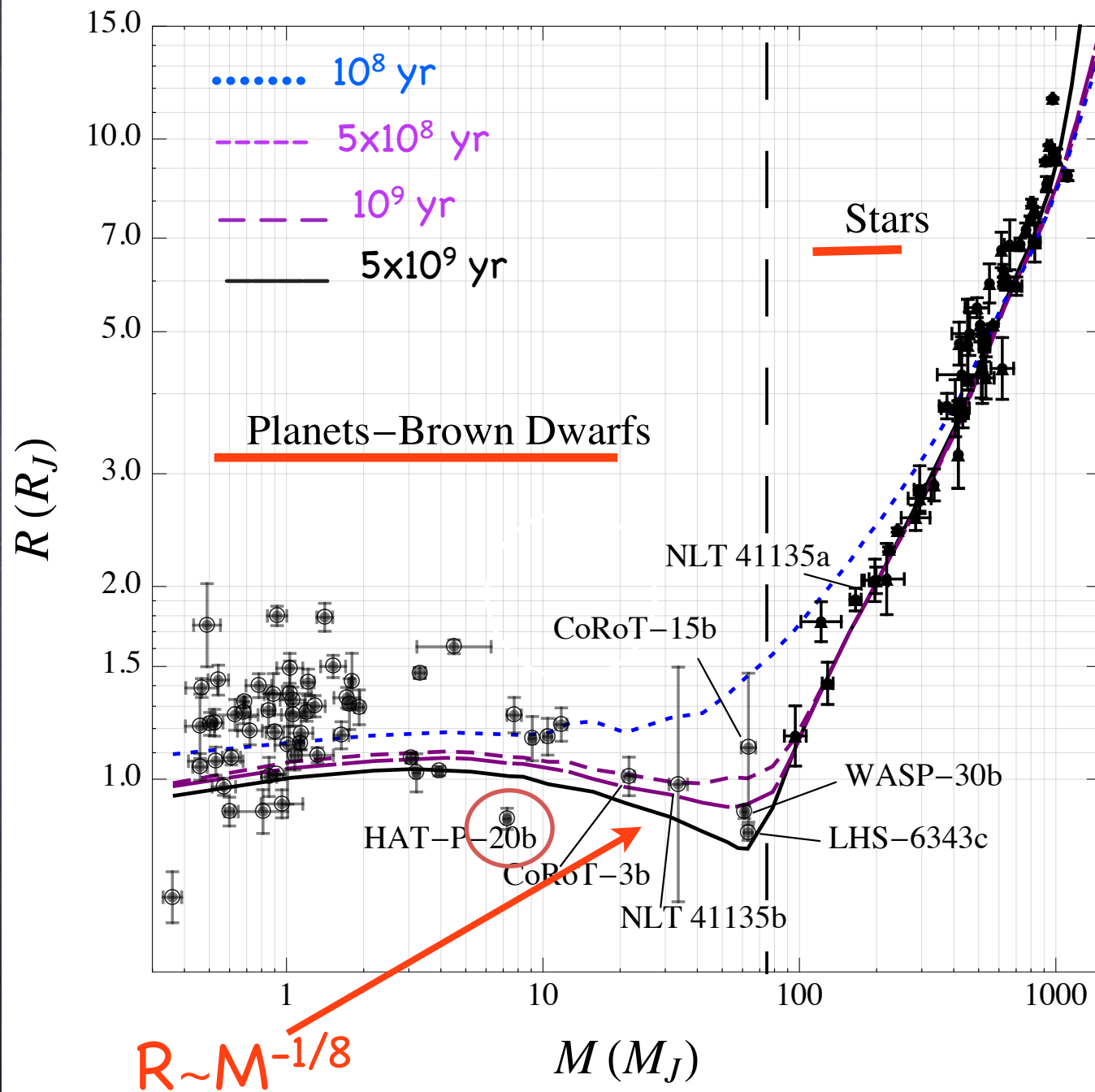
Brandt et al. (2015)

combined sample (SEEDS, GDPS, NICI MG) of 250 stars over a wide range of ages and Sp

<5% of stars have substellar 5-70 M<sub>Jup</sub> comp'ns between 10-100 AU @ 95% C.L.

<11% of stars have substellar 10-70 M<sub>Jup</sub> comp'ns between 30-1600 AU @ 95% C.L.

## 3) THE DIAGNOSTIC(S)

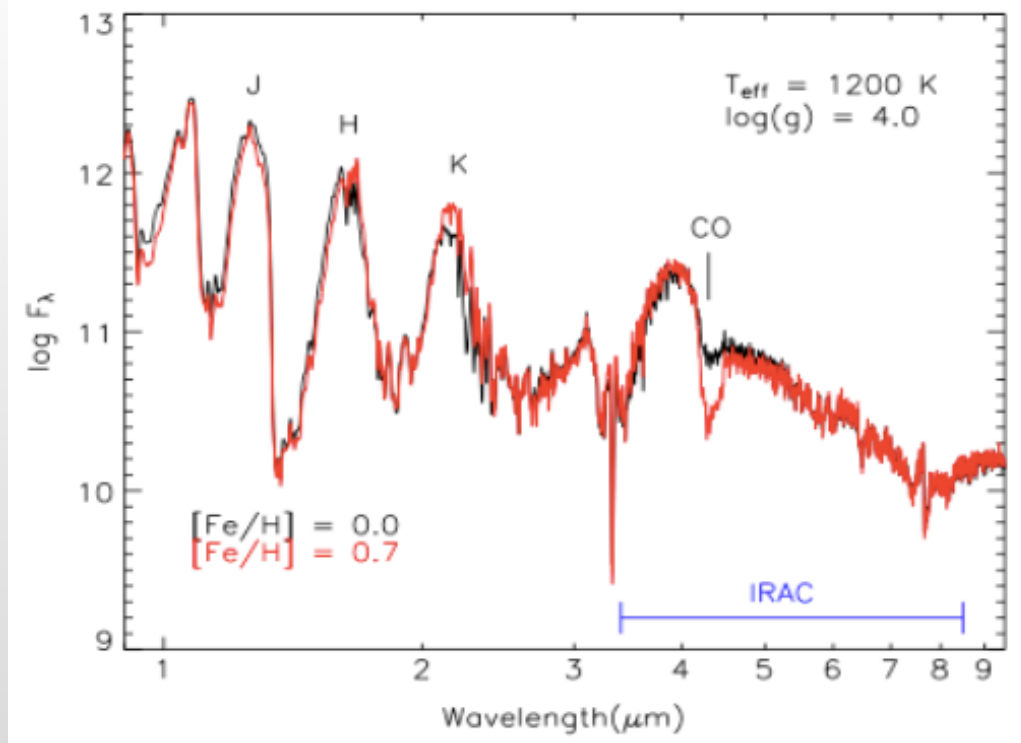


(Chabrier & Baraffe 2000)

Chabrier, Leconte & Baraffe 2010

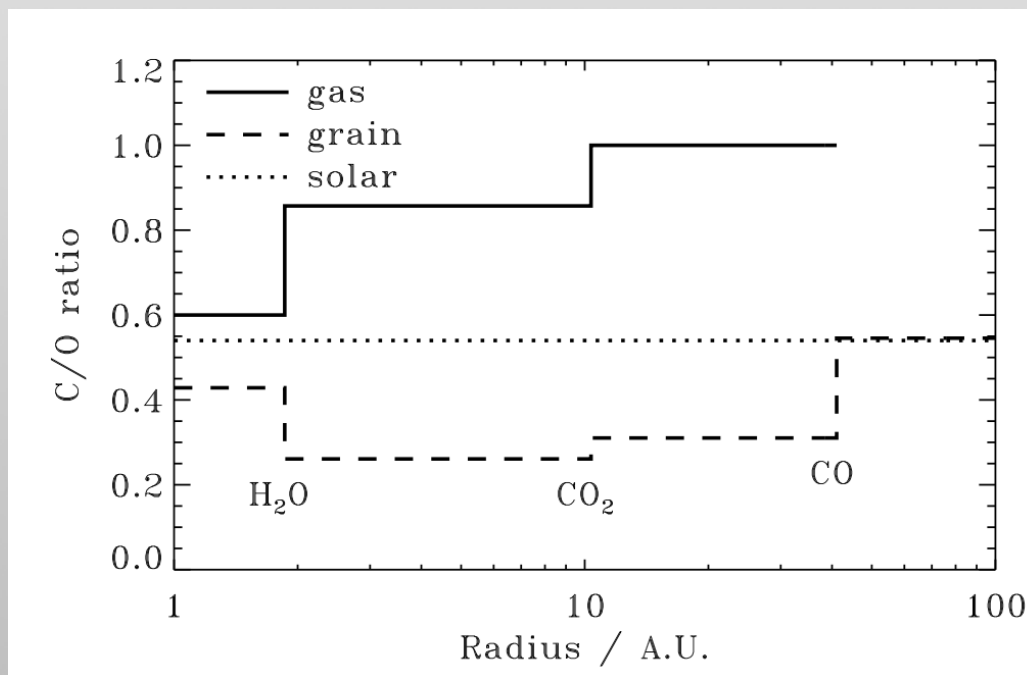
# Distinguishing (low-mass) BDs from (massive) GPs

Effect of an increase of metallicity (factor 5) on spectra (*Chabrier et al. 2007*)



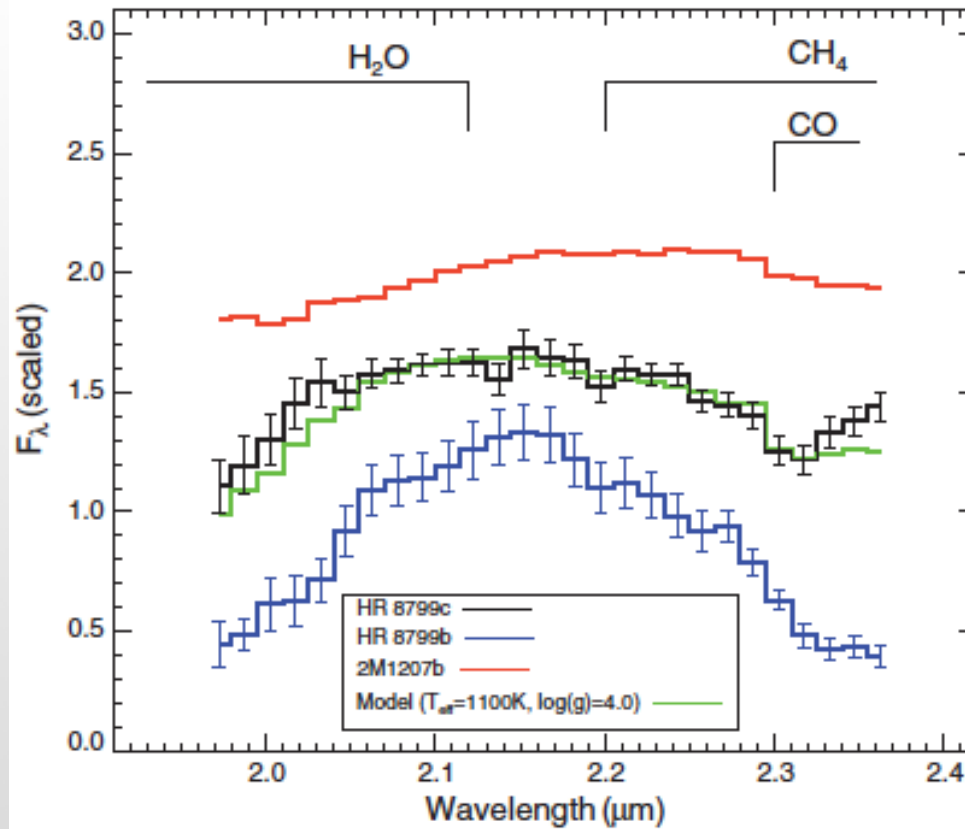
- deeper CO absorption
  - enhanced K-band flux
- Fortney et al. 2008*

*spectrum calculated by T. Barman*



C/O enhanced in regions of gas giant formation

*Öberg et al. 2008*



**HR 8799c:** enhanced C/O ratio and depleted C and O levels tend to favor a formation by C.A. (Konopacky et al. 2013)

Not so simple!....

- non-equilibrium chemistry (see talk by P. Tremblin)
- depends on the region of formation (different condensation lines)
- migration ....
- HOWEVER, **stellar C/O expected for planets formed by G.I.**

## Conclusion 1

- Brown dwarfs form **dominantly** like stars, by gravoturbulent frag'n of a parent cloud/core (emerging population of isolated pre- and proto-BD cores)
- Overwhelming majority of discovered **exoplanets** consistent w/ predictions from **Core Accretion**
- Formation by **G.I. requires very peculiar (extreme ?) conditions**  
Adding migration makes the problem even more acute (requires some «fine tuning»)  
Even if G.I. occurs, survival of the fragment is still an issue
- G.I. might be operational at large orbital distances.  
However:
  - \* scarcity of objects at such distances (constraints from direct imaging)
  - \* C.A. might work («pebble» accretion)
  - \* other possibilities (outward migration, planet scattering,...)... stay tuned ! Future projects (SPHERE, GPI, CHARIS) will help
  - G.I. might produce some objects (eg [Delorme et al. '13](#))
- **Diagnostics to distinguish BDs from GPs:**
  - **M-R** ex.: Hat P 2b ( $9 M_{\text{Jup}}$ ), Hat P 20b ( $7 M_{\text{Jup}}$ ),  $M_{\text{core}} > 200 M_{\text{Earth}}$  ([Leconte et al. '09, '11](#))
  - **spectro** : C/O, C/H expected to be non-stellar if formation by C.A.

## Conclusion 2

- **BD f'n extends down to the opacity limit fragmentation (~1-5 M<sub>Jup</sub>)**  
(D-burning and non D-burning BDs, the same way PP and CNO stars)
  - **GP f'n seems to extend at least to ~7-9 M<sub>Jup</sub> (Hat P 2b, Hat P 20b)**
  - \* K And b : **~10-20 M<sub>Jup</sub>**, 40 AU companion to a B9 star w/ **q=M2/M1<1%**
  - \* 2M 1207 b : **~4 M<sub>Jup</sub>**, ≥ 55 AU companion to a BD w/ **q=20%**
- => existence of non D-burning BD's**  
**/ D-burning planets not unlikely** (Baraffe et al. '08, Mollière & Mordasini '12)

*The very definition of a brown dwarf or a giant planet is intrinsically, tightly linked to its **formation mechanism**.*

- Free floating objects down to a few M<sub>Jup</sub> can rather unambiguously be identified as genuine (non D-burning) **brown dwarfs**
  - **planets** are necessarily companions of a central, significantly more massive object. Should have a non-stellar mean metal abundance
- ? {  
- genuine (ejected) free floating planets  
- ~M<sub>Jup</sub> objects formed by G.I.

**BD AND GP MASS DOMAIN OVERLAP**  
**DEUTERIUM BURNING PLAYS NO ROLE**  
**IAU DEFINITION SCIENTIFICALLY MEANINGLESS**