Differentially rotating neutron stars: A perturbative study

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# Outline

# Astrophysical motivation Outline of perturbative method Results Discussion-Conclusions

# Astrophysical motivation

Neutron stars are born with differential rotation but as they cool differential rotation (DR) is smoothed out to uniform Window with DR that can be astrophysically interesting

#### **Newtonian studies**

□ Hansen et.al, ApJ, 217, 151, 1977, effect of DR on mode splitting

□ Karino & Eriguchi, ApJ, **578**, 413, 2002, GR reaction & f-mode

#### **Full GR studies**

□ Komatsu et.al, MNRAS, **239**, 153, 1989, Construct DR models, F = f(M,J,K,n) parameter to classify equilibrium models.

□ Stergioulas et.al MNRAS, **352**, 1089, 2004, splitting of f-mode

Dimmelmeier et.al. MNRAS, **368**, 1609, 2006, Axisymmetric modes CFC

Perturbative GR studies missing....

# Low T/W instability

Newtonian simulations have shown that stars with high degree of differential rotation, show a dynamical instability at low values of  $T/W \sim 0.08$  to 0.14.

- $\Box$ J. Centrella et.al, ApJ, **550**, L196, 2001 $\rightarrow$  N=3.33 polytrope, T/W ~ 0.14
- □Shibata et.el MNRAS, 334, L27, 2002  $\rightarrow$  N=1 polytrope T/W ~ 0.01

□A. Watts et.al ApJ, 618, L37, 2005  $\rightarrow$  due to f-mode entering into corotation band

□Saio & Yoshida, MNRAS, **368**, 1429, 2006 → diagnosis with canonical angular momentum.

 $\Box$ Ou & Tohline, ApJ, 651, 1068, 2006  $\rightarrow$  not necessarily rapidly rotating

**D**New & Shapiro, ApJ, **548**, 439, 2001  $\rightarrow$  for supermassive neutron stars

GW emitted could be detected by LISA in super massive NS

# Perturbative method

### Assumptions

 $\checkmark$  Slow rotation (star is spherical), 0< $\epsilon_{\rm e}$  =  $\Omega_{\rm e}$  /  $\Omega_{\rm K}$ <1

✓ J-constant differential rotation law

$$\Omega(r,\theta) = \frac{A^2 \Omega_c + e^{-2\nu} \omega(r,\theta) r^2 \sin^2 \theta}{A^2 + e^{-2\nu} r^2 \sin^2 \theta}$$

✓ Relativistic polytrope,  $p=Kρ^{\Gamma}$ , ε = ρ + p / (Γ-1)

- Cowling approximation (neglect spacetime perturbations)
- ✓ Non barotropic perturbations  $\Gamma \neq \Gamma_1$

Differential rotation background TOV equations + equation for frame dragging  $\omega'' - \left[4\pi(\varepsilon + p)re^{2\lambda} - \frac{4}{r}\right]\omega' - \left[16\pi(\varepsilon + p) + \frac{\Lambda - 2}{r^2}\right]e^{2\lambda}\omega = -16\pi(\varepsilon + p)e^{2\lambda}\Omega$ Expand terms  $\omega, \Omega$  in spherical harmonics,  $\omega_1, \omega_3, \Omega_1, \Omega_3$ 0.012 - A = 14.15 km $-\cdot A = 25 \text{ km}$ **C** 0.008  $\cdots A = 50 \text{ km}$ -- A = 75 km  $-\cdots$  A = 100 km 0.004 0.0006 0.0004 0.0002 0 2 0 4 6 10 12 14 km

# Comparison with non-linear backgrounds

RNS code by Stergioulas



### Perturbative Equations & Method

Perturbed conservation of energy-momentum  $\delta (T_{\mu\nu}; \mu) = 0$ 

Expand all variables in spherical harmonics  $A(t,r,\theta,\phi) = R(r,t) Y_{lm}(\theta,\phi)$ Integrate over solid angle to get PDEs of (r,t)

> 5 independent variables 4 polar (f,p,g modes) 1 axial (r,inertial modes,CS)

Infinitely coupled hyperbolic system of PDEs  $P_{lm} + I m (P_{lm} + A_{l\pm 1} + P_{l\pm 2}) + A_{l\pm 1} + A_{l\pm 3} = 0$   $A_{lm} + I m (A_{lm} + P_{l\pm 1} + A_{l\pm 2}) + P_{l\pm 1} + P_{l\pm 3} = 0$ To solve them we have to truncate for  $I_{max}$ 

# Results

# Effects of rotation to stellar modes

□ Mixes the character of the modes

polar-led, polar in non rotating limit

axial-led, axial in non rotating limit

□ Splitting of modes (like Zeeman effect in atomic physics)

 $\sigma^{lm} = \sigma_0^{lm} \pm \alpha(l,m,A) \varepsilon_e$ 

Study the effects with respect to the three parameters

- Maximum number of couplings I<sub>max</sub>
  Azimuthal index m
  Decree of differential rotation A
  - Degree of differential rotation A

# Effect of $I_{max}$ to quasi radial frequencies

max	F (kHz)	H <sub>1</sub> (kHz)	
0	2.687	4.551	Radial
1	2.710	4.571	Dipole
2	2.712	4.575	Non-radial I=2
3	2.712	4.575	Non-radial I=3

Frequencies have converged for  $I_{max}=2$ 

# Axis-symmetric perturbations Comparison with non-linear results

Model	F (kHz)	H <sub>1</sub> (kHz)	f <sub>2</sub> (kHz)	p <sub>2</sub> (kHz)
B0	2.706	4.547	1.846	4.100
B1	2.702 (2%)	4.555 (2%)	1.895 (1%)	4.117 (1%)
B3	2.735 (4%)	4.578 (4%)	1.915 (1%)	4.124 (2%)
B6	2.797 (6%)	4.624 (4%)	1.944 (1%)	4.134 (7%)
B9	2.885 (8%)	4.686 (6%)	1.974 (8%)	4.147 (14%)

#### Madac colitting



## Dependence of splitting to compactness



Passamonti, A.S, Kokkotas, PRD accepted

# Dependence of splitting to degree of differential rotation



#### Modes and corotation

Pattern speed of the mode :  $\sigma/m$ 

Corotation band :  $\Omega_{\rm e} < \Omega < \Omega_{\rm s}$ 

If the pattern speed of the mode is equal to the local angular velocity of the star we have a corotation mode

### f - mode and Corotation band



## Corotation points for different models



# Epilogue Discussion

# Drawbacks of our approximation

 $\Box$  Definition of rotational parameter  $\epsilon = \Omega_e / \Omega_K$  with respect to T/W

□ Cold polytropic EoS

□ Real code, cannot see damping/growth rate of modes

On going work and extensions :

□ Use hot EoS for nascent neutron stars to study g-modes

□Use relativistic canonical energy (??) to study the T/W instability.

Long term goal : GW asteroseismology, estimate stellar parameter from the detected GW signal.

Based on papers:

A.Stavridis, A. Passamonti, K.D. Kokkotas, PRD, 75, 064019, 2007A. Passamonti, A.Stavridis, K.D. Kokkotas, PRD, accepted..