

Gravitation modifiée à grande distance et tests dans le système solaire

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et

Peter Wolf, LNE-SYRTE

10 avril 2008

Modified gravity at large distances and solar-system tests

J.-P. Bruneton and **G. Esposito-Farèse**

$\mathcal{GR}\epsilon\mathcal{CO}$, Institut d'Astrophysique de Paris

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Discussions with L. Blanchet, C. Deffayet, B. Fort, G. Mamon,
Y. Mellier, M. Milgrom, R. Sanders, J.-P. Uzan, R. Woodard, *etc.*

April 10th, 2008

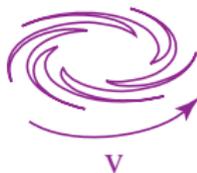
Dark matter and galaxy rotation curves

∃ evidences for dark matter:

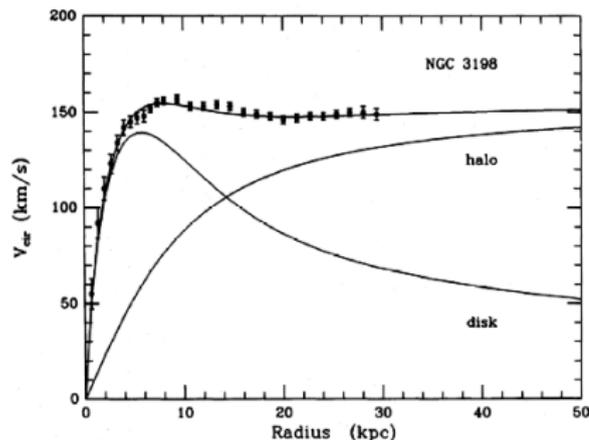
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- **Rotation curves**

of galaxies
and clusters:
almost rigid
bodies



DISTRIBUTION OF DARK MATTER IN NGC 3198



- ∃ many theoretical candidates for dark matter (e.g. from SUSY)
- Numerical simulations of structure formation are successful while incorporating (noninteracting, pressureless) dark matter

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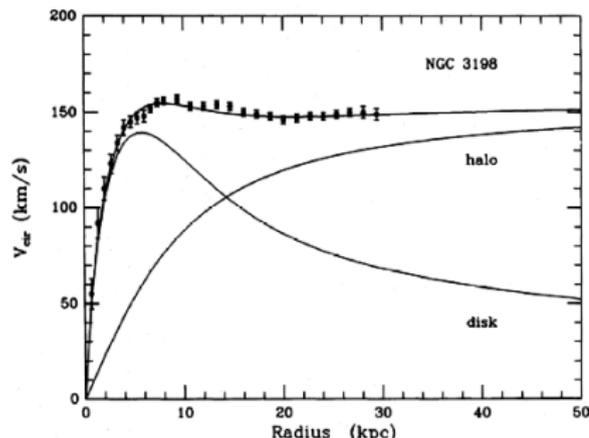
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Milgrom's MOND proposal [1983]

MODified Newtonian Dynamics for small accelerations (i.e., at large distances)

$$a = a_N = \frac{GM}{r^2} \quad \text{if } a > a_0 \approx 1.2 \times 10^{-10} \text{ m.s}^{-2}$$

$$a = \sqrt{a_0 a_N} = \frac{\sqrt{GM a_0}}{r} \quad \text{if } a < a_0$$

- Automatically recovers the Tully-Fisher law [1977]

$$v_\infty^4 \propto M_{\text{baryonic}}$$

- Superbly accounts for galaxy rotation curves
(but clusters still require some dark matter)

[Sanders & McGaugh, Ann. Rev. Astron. Astrophys. 40 (2002) 263]

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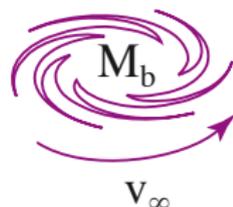
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Consistent field theories of MOND?

- *A priori* easy to predict a force $\propto 1/r$:

If $V(\varphi) = -2a^2 e^{-b\varphi}$, unbounded by below
 then $\Delta\varphi = V'(\varphi) \Rightarrow \varphi = (2/b) \ln(abr)$.

Constant coefficient $2/b$ instead of \sqrt{M} .

Some papers write actions which depend on the galaxy mass M
 \Rightarrow They are actually using a different theory for each galaxy!

- Stability

Full Hamiltonian should be bounded by below:
 no tachyon ($m^2 \geq 0$), no ghost ($E_{\text{kinetic}} \geq 0$)

- Well-posed Cauchy problem

Hyperbolic field equations

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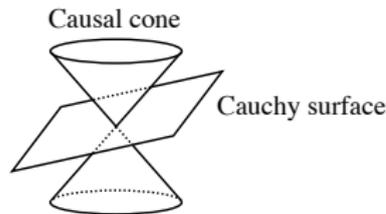


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Most promising framework

Relativistic AQUAdratic Lagrangians
[Bekenstein (TeV), Milgrom, Sanders]

$$S = \frac{c^3}{16\pi G} \int d^4x \sqrt{-g} \left\{ R - 2f(\partial_\mu \varphi \partial^\mu \varphi) \right\} \\ + S_{\text{matter}} \left[\text{matter} ; \tilde{g}_{\mu\nu} \equiv A^2(\varphi) g_{\mu\nu} + B(\varphi) U_\mu U_\nu \right]$$

- A “k-essence” kinetic term can yield the $\frac{\sqrt{GMa_0}}{r}$ MOND force
- Matter coupled to the scalar field
- “Disformal” term (almost) necessary to predict enough lensing

Consistency conditions on $f(\partial_\mu\varphi\partial^\mu\varphi)$

Hyperbolicity of the field equations + Hamiltonian bounded by below

- $\forall x, \quad f'(x) > 0$
- $\forall x, \quad 2xf''(x) + f'(x) > 0$

N.B.: If $f''(x) > 0$, the scalar field propagates faster than gravitons, but still causally
 \Rightarrow no need to impose $f''(x) \leq 0$

These conditions become much more complicated *within matter*

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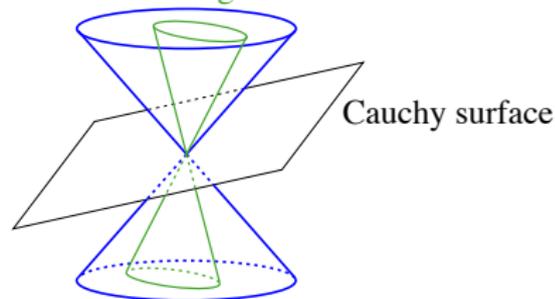
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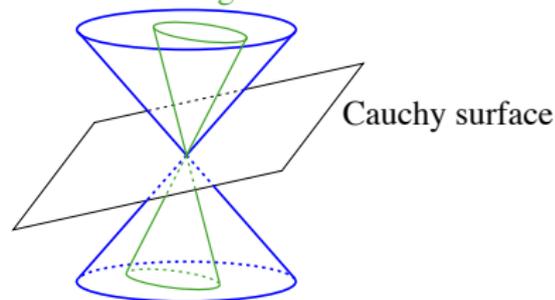
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Difficulties of such models

- Complicated Lagrangians (**unnatural**)
- **Fine tuning** (\approx fit rather than predictive models):
Possible to predict different lensing and rotation curves
- **Discontinuities**: can be cured
- In TeVeS [Bekenstein], gravitons & scalar are slower than photons
 \Rightarrow **gravi-Cerenkov radiation** suppresses high-energy cosmic rays
[Moore *et al.*]
Solution: Accept slower photons than gravitons
- \exists **preferred frame** (ether) where vector $U_\mu = (1, 0, 0, 0)$
Maybe not too problematic if U_μ is dynamical
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- **Post-Newtonian tests** very constraining

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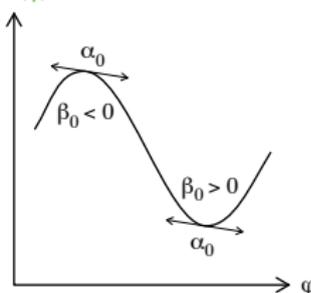
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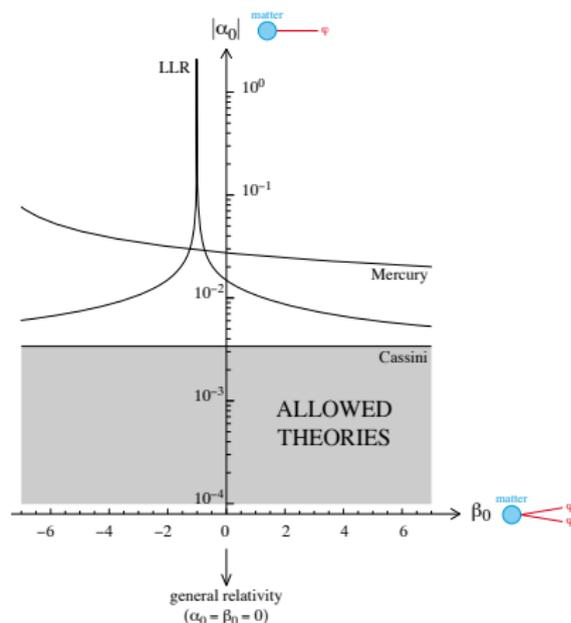
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- Solar-system tests \Rightarrow matter *a priori* **weakly** coupled to φ
- TeVeS *tuned* to pass them even for strong matter-scalar coupling
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matter-scalar
coupling function
 $\ln A(\varphi)$



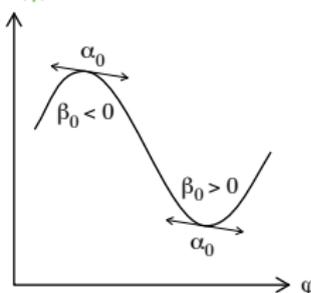
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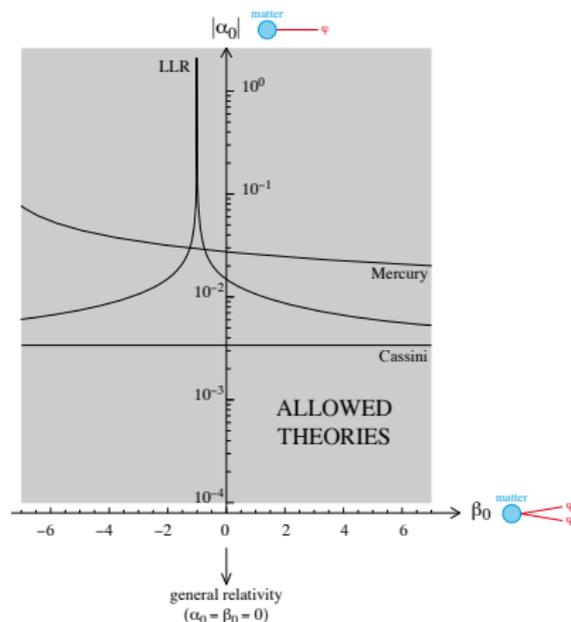
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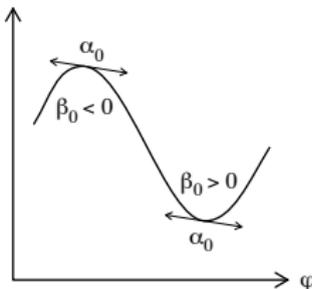
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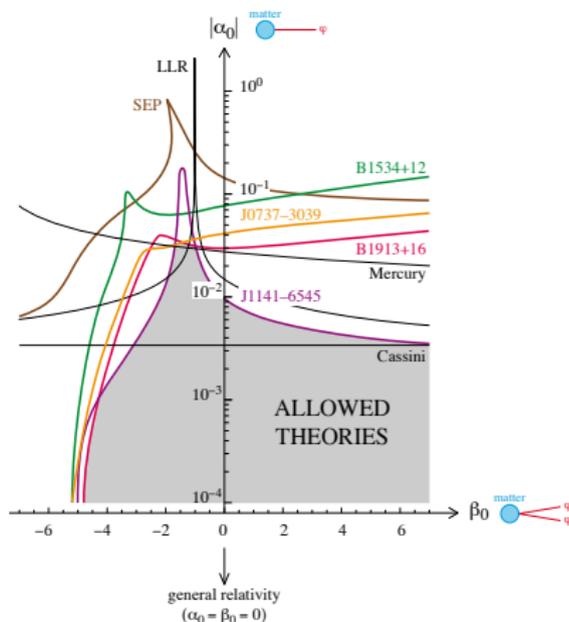
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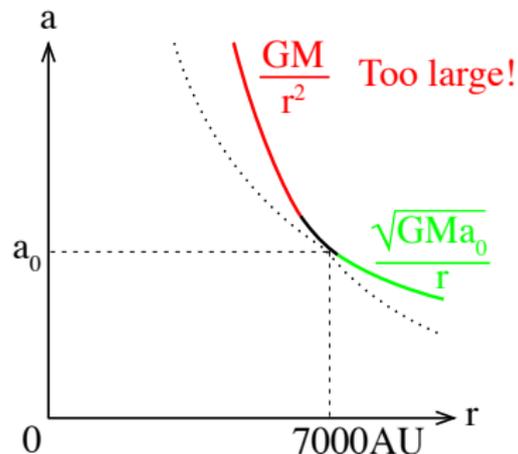
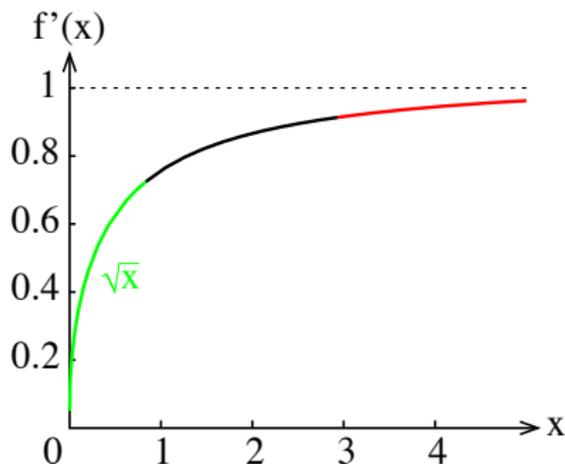


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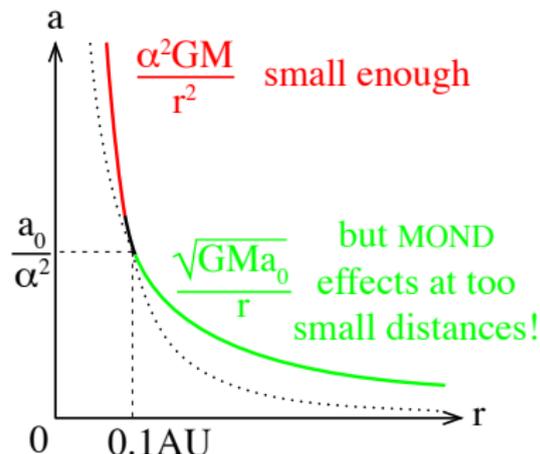
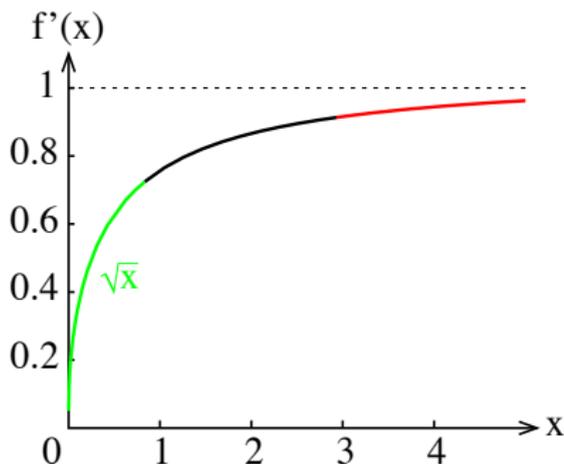
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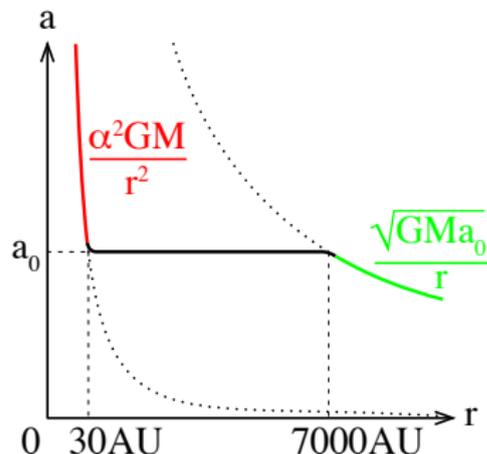
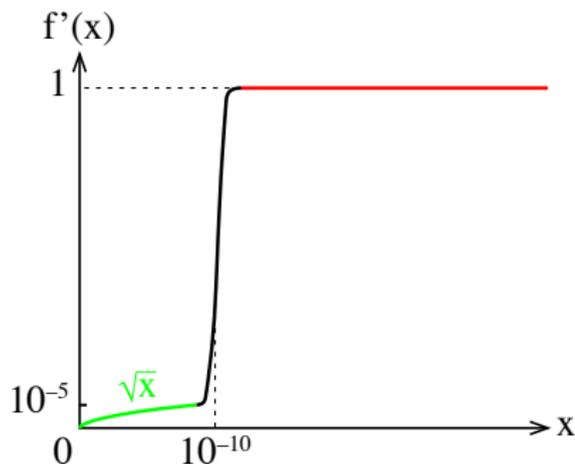
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Quite unnatural! (and not far from being experimentally ruled out)

New simpler models?

Nonminimal metric coupling

$$S = \frac{c^3}{16\pi G} \int d^4x \sqrt{-g} R \quad \text{pure G.R. in vacuum}$$

$$+ S_{\text{matter}} \left[\text{matter} ; \tilde{g}_{\mu\nu} \equiv f(g_{\mu\nu}, R^\lambda{}_{\mu\nu\rho}, \nabla_\sigma R^\lambda{}_{\mu\nu\rho}, \dots) \right]$$

Can reproduce MOND, but Ostrogradski [1850] \Rightarrow unstable within matter

Nonminimal scalar-tensor model

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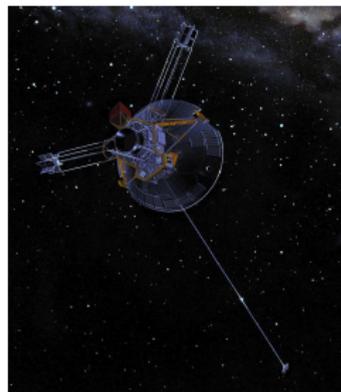
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Pioneer 10 & 11 anomaly

- **Extra acceleration** $\sim 8.5 \times 10^{-10} \text{ m.s}^{-2}$ towards the Sun between 30 and 70 AU
- **Simpler problem** than galaxy rotation curves ($M_{\text{dark}} \propto \sqrt{M_{\text{baryon}}}$), because we do not know how this acceleration is related to M_{\odot}
- \Rightarrow *several* stable & well-posed solutions



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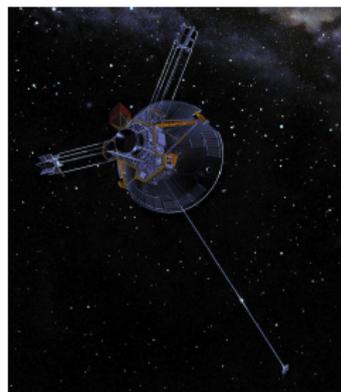
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- $\lambda \approx \alpha^3 (10^{-4} \text{ m})^2$ to fit Pioneer anomaly

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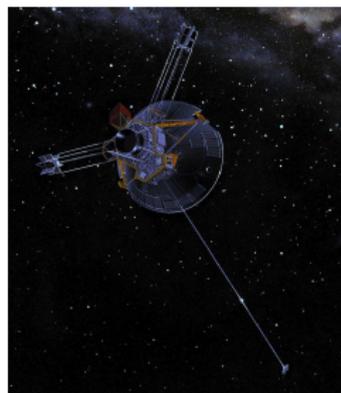
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Conclusions

A consistent field theory should satisfy different kinds of constraints:

- **Mathematical:** stability, well-posedness of the Cauchy problem, no discontinuous nor adynamical field
- **Experimental:** solar-system & binary-pulsar tests, galaxy rotation curves, gravitational lensing by “dark matter” haloes, CMB
- **Esthetical:** natural model, rather than fine-tuned *fit* of data

Best present candidate: TeVeS [Bekenstein–Sanders], but it has still some mathematical *and* experimental difficulties

∃ simpler models, useful to exhibit the generic difficulties of all MOND-like field theories

By-product of our study: a consistent class of models for the Pioneer anomaly (but *not* natural!)

Nonlocal models? [Work in progress with Cédric Deffayet & Richard Woodard]

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Best present candidate: TeVeS [Bekenstein–Sanders], but it has still some **mathematical** *and* **experimental** difficulties

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