

Realistic Cosmic Strings

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arXiv:1505.07888



Outline

1 Why a realistic model of cosmic strings

2 Theoretical framework

3 Abelian cosmic strings

4 Microscopic and macroscopic structure

Topological defects

- Topological defects

- ▶ Phase transitions with Spontaneous Symmetry Breaking (SSB)

$$G \xrightarrow{\langle \Phi \rangle} G'$$

- ▶ General predictions
- ▶ Topological properties of G/G'

Topological defects

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- ▶ Phase transitions with Spontaneous Symmetry Breaking (SSB)

$$G \xrightarrow{\langle \Phi \rangle} G'$$

- ▶ General predictions
- ▶ Topological properties of G/G'

- Different kinds

- ▶ Monopoles
- ▶ Domain walls
- ▶ Cosmic strings

• Cosmic strings

- ▶ Vortex lines
- ▶ $\pi_1(G/G') \neq 1$

Why a realistic model of cosmic strings

Theoretical framework

Abelian cosmic strings

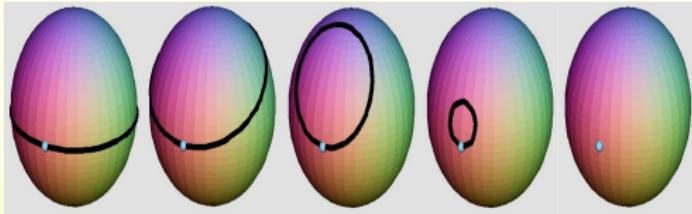
Microscopic and macroscopic structure

Topological defects

Abelian Higgs model

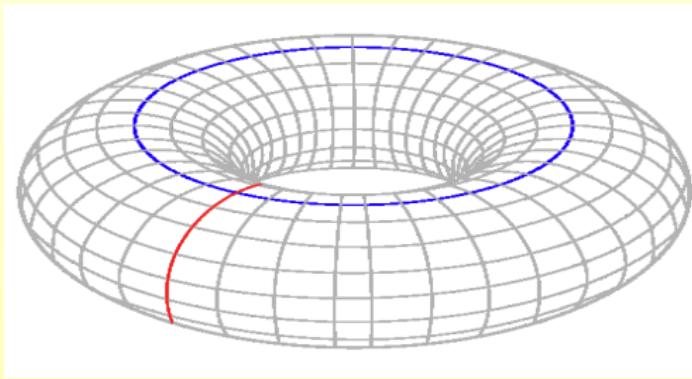
Cosmic strings and cosmology

Necessity of a realistic implementation



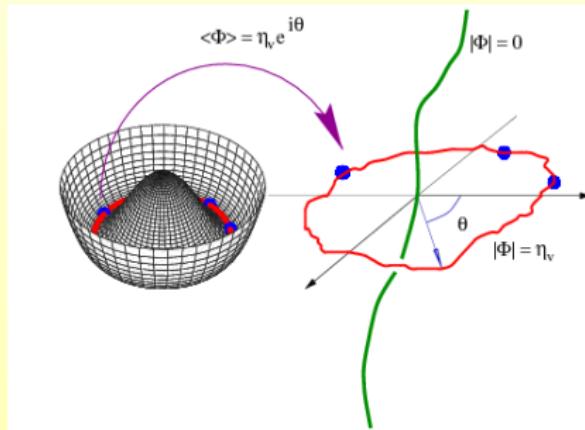
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- ▶ Vortex lines
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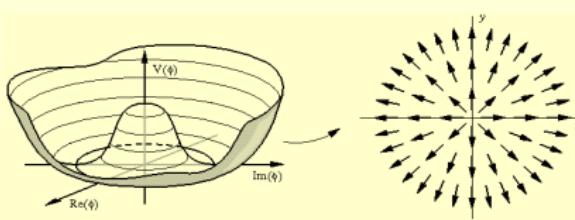
Abelian Higgs model

$$\mathcal{L} = -(\partial_\mu + iqA_\mu)\phi^*(\partial^\mu - iqA^\mu)\phi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \kappa^2(\phi^*\phi - M^2)^2$$



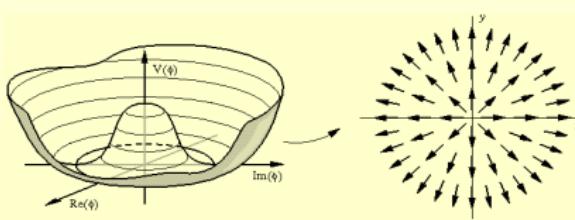
• Ansatz

- ▶ Cylindrical coordinates locally along the string (r, θ, z, t)
- ▶ $\phi = f(r)e^{in\theta}$ with $f \in \mathbb{R}$
- ▶ $A^\mu = A_\theta(r)\tau_{\text{str}}\delta_\mu^\theta$



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- ▶ $\lim_{r \rightarrow \infty} f(r) = M$ and $\lim_{r \rightarrow \infty} A_\theta^{\text{str}}(r) = n/g$
- ▶ $f(0) = 0$ and $A_\theta(0) = 0$

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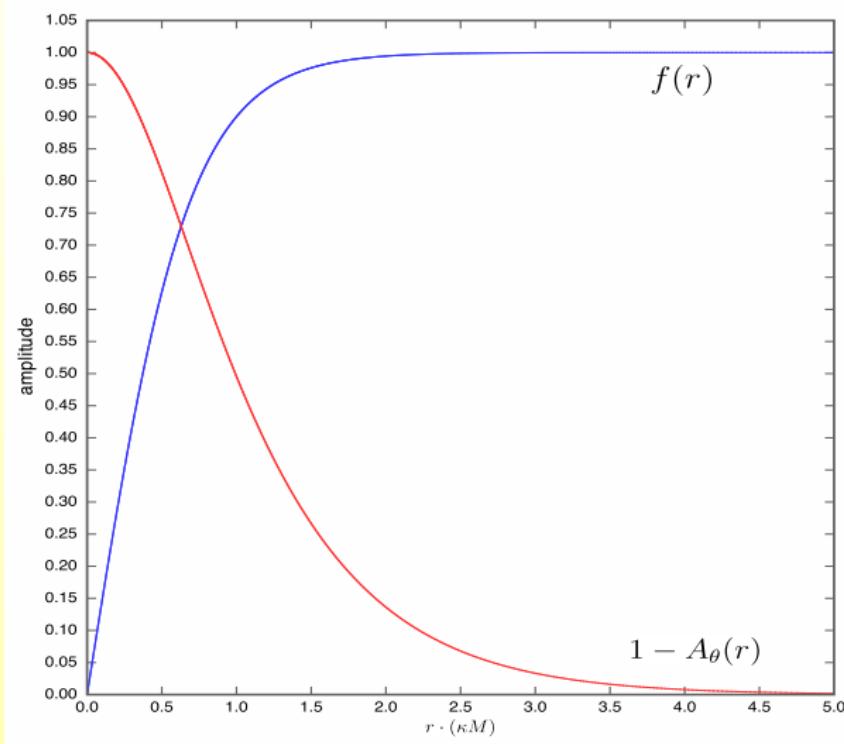
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Macroscopic properties

- From $T^{\mu\nu}$

- ▶ Symmetries $\Rightarrow T^{ab}$ with $a, b \in \{z, t\}$
- ▶ Diagonalization \Rightarrow energy per unit length and tension

$$U = \int 2\pi r dr T^{t't'} \quad T = - \int 2\pi r dr T^{z'z'}$$

- ▶ Equation of state $U = f(T)$

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- **Nambu-Goto String**

- ▶ Lorentz-invariant along the worldsheet
- ▶ $\Rightarrow U = T$

Cosmic strings and cosmology

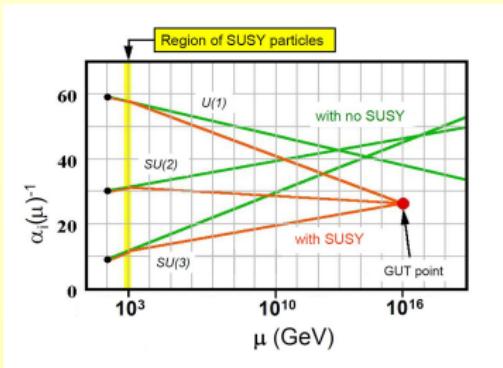
- **Grand Unified Theory (GUT)**

- ▶ Particle physics in one gauge group G_{GUT} , e.g. SO(10) or SU(6)
- ▶ SSB to Standard Model (SM) at $E_{\text{GUT}} \sim (10^{15} - 10^{16}) \text{GeV}$

Cosmic strings and cosmology

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- ▶ SSB to Standard Model (SM) at $E_{\text{GUT}} \sim (10^{15} - 10^{16}) \text{ GeV}$
- ▶ Many theoretical motivations, e.g. $q \in \mathbf{Z}/3$ or coupling constants



● Formation of strings

- ▶ During SSB from the GUT to the SM
- ▶ Formation of strings at the end of inflation

● Formation of strings

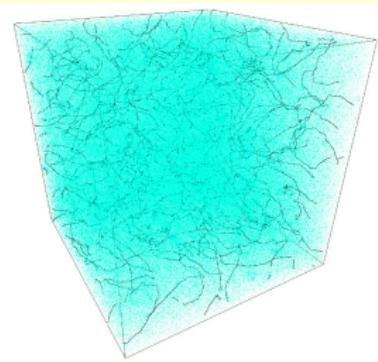
- ▶ During SSB from the GUT to the SM
- ▶ Formation of strings at the end of inflation
- ▶ Almost unavoidable consequence

$$\begin{array}{c}
 \text{SU}(6) \left\{ \begin{array}{l}
 \stackrel{1}{\rightarrow} 5 \ 1_6 \\
 \stackrel{1}{\rightarrow} 4_C \ 2_L \ 1 \\
 \stackrel{1}{\rightarrow} 3_C \ 3_L \ 1
 \end{array} \right. \quad \left. \begin{array}{l}
 \left\{ \begin{array}{l}
 \stackrel{2}{\rightarrow} 5 \qquad \qquad \stackrel{1}{\rightarrow} G_{\text{SM}} \\
 \stackrel{1}{\rightarrow} 3_C \ 2_L \ 1 \ 1 \ \stackrel{2}{\rightarrow} G_{\text{SM}} \\
 \stackrel{1,2}{\rightarrow} G_{\text{SM}}
 \end{array} \right. \\
 \left\{ \begin{array}{l}
 \stackrel{1}{\rightarrow} 3_C \ 2_L \ 1 \ 1 \ \stackrel{2}{\rightarrow} G_{\text{SM}} \\
 \stackrel{2'}{\rightarrow} G_{\text{SM}}
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 \stackrel{1}{\rightarrow} 3_C \ 2_L \ 1 \ 1 \ \stackrel{2}{\rightarrow} G_{\text{SM}} \\
 \stackrel{0}{\rightarrow} 5 \qquad \qquad \stackrel{1}{\rightarrow} G_{\text{SM}} \\
 \stackrel{1}{\rightarrow} G_{\text{SM}}.
 \end{array} \right.
 \end{array}$$

Jeannerot, Rocher, Sakellariadou,
 Phys.Rev., D68:103514, (2003)

• String Network

- ▶ Initial conditions from correlation length
- ▶ Numerical simulations



Ringeval, Sakellariadou, Bouchet, JCAP 0702 (2007)

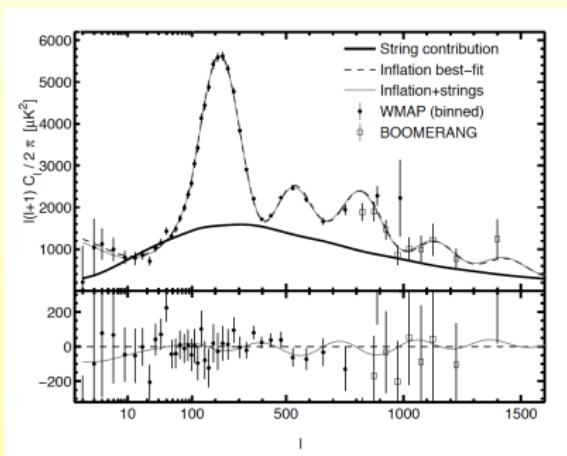
• Evolution from large to small scales

• In the CMB

- ▶ Good candidate before inflation to seed the LSS formation
- ▶ Power spectrum different from inflationary prediction

• In the CMB

- ▶ Good candidate before inflation to seed the LSS formation
- ▶ Power spectrum different from inflationary prediction
- ▶ No significant observations



Bevis, Hindmarsh, Kunz, Urrestilla, Phys. Rev. Lett. 100 (2008)

Necessity of a realistic implementation

- Toy model strings

- ▶ Only scale of formation appears
- ▶ Does not encode the GUT

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- Realistic implementation

- ▶ U as function of the GUT parameters
- ▶ New structure and properties
- ▶ Could differ from toy-models

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SUSY GUT

- **GUT**

- ▶ Gauge group $\text{SO}(10)$
- ▶ 45 gauge field A_μ^a
- ▶ Higgs field content : Σ in **126**, $\bar{\Sigma}$ in **$\overline{126}$** and Φ in **210**

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- ▶ Gauge group $\text{SO}(10)$
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- ▶ Higgs field content : Σ in **126**, $\bar{\Sigma}$ in **$\overline{126}$** and Φ in **210**

- **Valid with particle physics**

- ▶ Recover the SM at low energy
- ▶ Coupling constant, mass spectrum, proton decay...

• Supersymmetry (SUSY)

- ▶ Bosonic part of the supermultiplets
- ▶ F -term theory $\Rightarrow D$ -term identically zero

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- ▶ F -term theory $\Rightarrow D$ -term identically zero

- **Potential from F -terms**

- ▶ Superpotential W , singlet, X^3 ($\Rightarrow X^4 \in V$)
- ▶ Potential from F -terms

$$F_X = \frac{\partial W}{\partial X} \quad \text{and} \quad V_X = F^\dagger F$$

- With $\text{SO}(10)$

$$W = \frac{m}{2} \Phi^2 + m_\Sigma \Sigma \bar{\Sigma} + \frac{\lambda}{3} \Phi^3 + \eta \Phi \Sigma \bar{\Sigma}$$

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$$F_\Phi = m\Phi + \lambda(\Phi^2)_\Phi + \eta(\Sigma \bar{\Sigma})_\Phi$$

$$\begin{aligned} V_\Phi = & m^2 \Phi \Phi^\dagger + |\eta|^2 (\Sigma \bar{\Sigma})_\Phi (\Sigma \bar{\Sigma})_\Phi^\dagger + |\lambda|^2 (\Phi^2)_\Phi (\Phi^2)_\Phi^\dagger \\ & + \left[\lambda \eta^* (\Phi^2)_\Phi (\Sigma \bar{\Sigma})_\Phi^\dagger + m \lambda^* \Phi (\Phi^2)_\Phi^\dagger + m \eta^* \Phi (\Sigma \bar{\Sigma})_\Phi^\dagger \right] + \text{h.c.} \end{aligned}$$

Hybrid inflation

- **Hybrid inflation (HI)**

- ▶ Inflationary potential from GUT physics
- ▶ Inflaton S , SO(10) singlet
- ▶ Additional term in W

$$W_{\text{HI}} = \kappa S (\Sigma \bar{\Sigma} - M^2)$$

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- **Generalization**

- ▶ SUSY, HI, SO(10) = example

$$V_S = \kappa^2 (\Sigma \bar{\Sigma} - M^2)^2$$

$$\begin{aligned} V_{\bar{\Sigma}} = & m_\Sigma^2 \Sigma \Sigma^\dagger + |\eta|^2 (\Phi \Sigma)_\Sigma (\Phi \Sigma)^\dagger_\Sigma + \kappa^2 S S^* \Sigma \Sigma^\dagger + \left[m_\Sigma \kappa S^* \Sigma \Sigma^\dagger \right. \\ & \left. + \eta \kappa S^* (\Phi \Sigma)_\Sigma \Sigma^\dagger + m_\Sigma \eta^* \Sigma (\Phi \Sigma)^\dagger_\Sigma \right] + \text{h.c.} \end{aligned}$$

$$V_\Sigma = V_{\bar{\Sigma}}(\Sigma \longleftrightarrow \bar{\Sigma})$$

$$V_S = \kappa^2 (\Sigma \bar{\Sigma} - M^2)^2$$

$$\begin{aligned} V_{\bar{\Sigma}} = & m_{\Sigma}^2 \Sigma \Sigma^\dagger + |\eta|^2 (\Phi \Sigma)_{\Sigma} (\Phi \Sigma)_{\Sigma}^\dagger + \overbrace{\kappa^2 S S^* \Sigma \Sigma^\dagger}^{\propto |S^2|} + \left[m_{\Sigma} \kappa S^* \Sigma \Sigma^\dagger \right. \\ & \left. + \eta \kappa S^* (\Phi \Sigma)_{\Sigma} \Sigma^\dagger + m_{\Sigma} \eta^* \Sigma (\Phi \Sigma)_{\Sigma}^\dagger \right] + \text{h.c.} \end{aligned}$$

$$V_{\Sigma} = V_{\bar{\Sigma}}(\Sigma \longleftrightarrow \bar{\Sigma})$$

Cosmological evolution

- SSB scheme

$$\mathrm{SO}(10) \xrightarrow{\langle\Phi\rangle} G' \xrightarrow{\langle\Sigma\rangle} G_{\mathrm{SM}} \times \mathbf{Z}_2$$

Cosmological evolution

- **SSB scheme**

$$\text{SO}(10) \xrightarrow{\langle\Phi\rangle} G' \xrightarrow{\langle\Sigma\rangle} G_{\text{SM}} \times \mathbf{Z}_2$$

- ▶ Monopoles form at first step (before inflation)
- ▶ Inflation begins after first step, ends at second step
- ▶ Cosmic strings form at second step (after inflation)

- First step

- ▶ $V_\Sigma \sim |\kappa S \Sigma|^2 \Rightarrow \Sigma = 0$
- ▶ $V_\Phi = 0 \Rightarrow$ several solutions, associated to different G'
- ▶ Inflaton : $V = \kappa^2 M^4 + \text{quant. corr.} \Rightarrow$ inflation begins

- **First step**

- ▶ $V_\Sigma \sim |\kappa S \Sigma|^2 \Rightarrow \Sigma = 0$
- ▶ $V_\Phi = 0 \Rightarrow$ several solutions, associated to different G'
- ▶ Inflaton : $V = \kappa^2 M^4 + \text{quant. corr.} \Rightarrow$ inflation begins

- **Second step**

- ▶ $S \longrightarrow S_{\text{crit}}$
- ▶ $V = 0 \Rightarrow$ several solutions, different gauge symmetries

• Requirement

- ▶ Recover the SM at low energy
- ▶ Stability of inflationary valley

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- **Two possible schemes**

- ▶ $G'_1 = 3_C 2_L 2_R 1_{B-L}$
- ▶ $G'_2 = 3_C 2_L 1_R 1_{B-L}$

Description with singlet of the SM

- Proper description

- ▶ Vacuum Expectations Values (VEVs) at the end of SSB
- ▶ No symmetry restoration
- ▶ Simplify the model

- **Branching rules**

- ▶ Description on a (maximal) subalgebra
- ▶ $SU(4) \subset SU(3) \times U(1)$
- ▶ Restricted representations w.r.t. the subalgebra

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- ▶ Restricted representations w.r.t. the subalgebra

$$\left\{ \begin{array}{l} \mathbf{1} = \mathbf{1}(0) \\ \mathbf{10} = \mathbf{1}(-6) + \mathbf{3}(-2) + \mathbf{6}(2) \\ \mathbf{15} = \mathbf{1}(0) + \mathbf{3}(4) + \overline{\mathbf{3}}(-4) + \mathbf{8}(0) \end{array} \right.$$

- **SO(10) case**

- ▶ Σ in **126**, $\bar{\Sigma}$ in **$\overline{126}$** and Φ in **210**
- ▶ Pati-Salam group : $2_L 2_R 4_C$

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$$\begin{aligned}\Phi_p &= \Phi(1, 1, 1), & \sigma &= \Sigma(1, 3, \overline{10}), \\ \Phi_a &= \Phi(1, 1, 15), & \overline{\sigma} &= \overline{\Sigma}(1, 3, 10), \\ \Phi_b &= \Phi(1, 3, 15), & S &= S(1, 1, 1).\end{aligned}$$

- **SSB scheme**

- ▶ $V(\Phi_a, \Phi_b, \Phi_p, \sigma, \bar{\sigma}, S)$
- ▶ Properly describe the SSB scheme

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$$\text{SO}(10) \xrightarrow{\langle\Phi\rangle} G' \xrightarrow{\langle\Sigma\rangle} G_{\text{SM}} \times \mathbf{Z}_2$$

- ▶ $G'_1 = 3_C 2_L 2_R 1_{B-L}$: $\Phi_a \neq 0$
- ▶ $G'_2 = 3_C 2_L 1_R 1_{B-L}$: $\Phi_a \neq 0$, $\Phi_b \neq 0$ and $\Phi_p \neq 0$

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Cosmic strings

$$\mathrm{SO}(10) \xrightarrow{\langle\Phi\rangle} G' \xrightarrow{\langle\Sigma\rangle} G_{\mathrm{SM}} \times \mathbf{Z}_2$$

- **String-forming step**

- ▶ From G' to $G_{\mathrm{SM}} \times \mathbf{Z}_2$
- ▶ At the end of inflation
- ▶ At least one $\mathrm{U}(1)$ is broken

Cosmic strings

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- **String-forming step**

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- ▶ At the end of inflation
- ▶ At least one $\mathrm{U}(1)$ is broken

- **Abelian string**

- ▶ Non abelian case left aside
- ▶ No current
- ▶ $\mathrm{U}(1)_{\mathrm{str}}$ and τ_{str}
- ▶ $\mathrm{U}(1)_{\mathrm{str}} = \mathrm{U}(1)_{\mathrm{B-L}}$

● Ansatz for the string

- ▶ Cylindrical coordinates locally along the string (r, θ, z, t)
- ▶ $\Sigma = \bar{\Sigma} = 0$ in the center of string
- ▶ VEVs defining the SM at infinity

- **Ansatz for the string**

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- **Coupling with the gauge field**

- ▶ Only Σ and $\bar{\Sigma}$ charged under $U(1)_{\text{str}}$
- ▶ Phase in $e^{\pm in\theta}$

- **Use of SM singlets**

- ▶ (1) Potential at least quadratic in the non singlet terms

$$V = \varphi\varphi\varphi\tilde{\varphi} \quad \Rightarrow \text{not possible}$$

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- **Use of SM singlets**

- ▶ (1) Potential at least quadratic in the non singlet terms

$$V = \varphi\varphi\varphi\tilde{\varphi} \quad \Rightarrow \text{not possible}$$

- ▶ (2) They all vanish at infinity
- ▶ (1) + (2) \Rightarrow configuration where they all vanish permitted by the EOM

Tensor formulation

- **Fields**

- ▶ Σ : (**126**) fifth rank anti-symmetric tensor Σ_{ijklm} , self dual

$$\Sigma_{ijklm} = \frac{i}{5!} \epsilon_{ijklmabcde} \Sigma_{abcde},$$

- ▶ $\bar{\Sigma}$ (**126**): fifth rank anti-symmetric tensor $\bar{\Sigma}_{ijklm}$, anti-self-dual

$$\bar{\Sigma}_{ijklm} = -\frac{i}{5!} \epsilon_{ijklmabcde} \bar{\Sigma}_{abcde},$$

- ▶ Φ (**210**) fourth rank anti-symmetric tensor Φ_{ijkl}
- ▶ S singlet: complex function

$$W = \frac{1}{2}m\Phi_{ijkl}\Phi_{ijkl} + m_\Sigma\Sigma_{ijklm}\bar{\Sigma}_{ijklm} + \frac{1}{3}\lambda\Phi_{ijkl}\Phi_{klmn}\Phi_{mnij} \\ + \eta\Phi_{ijkl}\Sigma_{ijmno}\bar{\Sigma}_{klmno} + \kappa S(\Sigma_{ijklm}\bar{\Sigma}_{ijklm} - M^2)$$

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• ***F*-terms computation**

- ▶ Taking into account that the components are not independent
- ▶ F_X in the conjugate representation of X

$$(F_\Phi)_{ijkl} = m\Phi_{ijkl} + \lambda\Phi_{[ij|ab}\Phi_{ab|kl]} + \eta\Sigma_{[ij|abc}\bar{\Sigma}_{|kl]abc}$$

- Group theoretical approach

$$\Phi \Sigma \bar{\Sigma} = 1_{210 \times 126 \times \bar{126}}$$

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$$\Phi \Sigma \bar{\Sigma} = \mathbf{1}_{\mathbf{210} \times \mathbf{126} \times \overline{\mathbf{126}}}$$

$$\mathbf{210} \times \overline{\mathbf{126}} = \mathbf{10} + \mathbf{120} + \overline{\mathbf{126}} + \mathbf{320} + \dots$$

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$$\mathbf{210} \times \overline{\mathbf{126}} = \mathbf{10} + \mathbf{120} + \overline{\mathbf{126}} + \mathbf{320} + \dots$$

$$\mathbf{1}_{\mathbf{210} \times \mathbf{126} \times \overline{\mathbf{126}}} \ni \mathbf{126} \times (\mathbf{210} \times \overline{\mathbf{126}})_{\overline{\mathbf{126}}}$$

- Group theoretical approach

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$$\mathbf{210} \times \overline{\mathbf{126}} = \mathbf{10} + \mathbf{120} + \overline{\mathbf{126}} + \mathbf{320} + \dots$$

$$\mathbf{1}_{\mathbf{210} \times \mathbf{126} \times \overline{\mathbf{126}}} \ni \mathbf{126} \times (\mathbf{210} \times \overline{\mathbf{126}})_{\overline{\mathbf{126}}}$$

$$\Rightarrow \frac{\partial(\Phi \Sigma \bar{\Sigma})}{\partial \Sigma} = (\Phi \bar{\Sigma})_{\bar{\Sigma}}$$

- **SM singlet formulation**

$$\Phi_a = a(x^\mu) \langle \Phi_a \rangle_0 \quad \text{with} \quad \langle \Phi_a \rangle_0 \langle \Phi_a \rangle_0^\dagger = 1$$

- **SM singlet formulation**

$$\Phi_a = a(x^\mu) \langle \Phi_a \rangle_0 \quad \text{with} \quad \langle \Phi_a \rangle_0 \langle \Phi_a \rangle_0^\dagger = 1$$

$$\begin{cases} \Sigma = \langle \bar{\Sigma} \rangle^\dagger = \sigma(x^\mu) \langle \sigma \rangle_0, \\ \Phi = a(x^\mu) \langle \Phi_a \rangle_0 + b(x^\mu) \langle \Phi_b \rangle_0 + p(x^\mu) \langle \Phi_p \rangle_0, \end{cases}$$

$$\left\{ \begin{array}{l} \frac{p}{\sqrt{4!}} = \Phi_{1234} \\ \frac{\sqrt{4!}a}{3} = \Phi_{5678} = \Phi_{5690} = \Phi_{7890} \\ \frac{b}{\sqrt{4!6}} = \Phi_{1256} = \Phi_{1278} = \Phi_{1290} = \Phi_{3456} = \Phi_{3478} = \Phi_{3490} \\ \frac{1}{\sqrt{5!2^5}}(i)^{(-a-b+c+d+e)}\sigma = \Sigma_{a+1,b+3,c+5,d+7,e+9} \\ \frac{1}{\sqrt{5!2^5}}(-i)^{(-a-b+c+d+e)}\bar{\sigma} = \bar{\Sigma}_{a+1,b+3,c+5,d+7,e+9} \end{array} \right.$$

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Cosmic strings
Tensor formulation
Ansatz for the string

$$\langle \Phi \rangle \langle \Phi \rangle^\dagger = pp^* + aa^* + bb^*$$

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$$\langle \Phi \Phi \rangle_\Phi \langle \Phi \rangle^\dagger = \underbrace{\frac{1}{9\sqrt{2}}}_{\sqrt{N}} a^2 a^* + \left[\frac{1}{6\sqrt{6}} (2b^* b p + b^2 p^*) + \frac{1}{9\sqrt{2}} (2b^* b a + b^2 a^*) \right]$$

$$\langle \Phi \rangle \langle \Phi \rangle^\dagger = pp^* + aa^* + bb^*$$

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$$\langle \Phi \Phi \rangle_\Phi \langle \Sigma \bar{\Sigma} \rangle_\Phi^\dagger = \underbrace{\frac{a^2 \sigma \sigma^*}{180}}_N + \frac{b^2 \sigma \sigma^*}{120} - \frac{ab \sigma \sigma^*}{45\sqrt{2}} - \frac{bp \sigma \sigma^*}{30\sqrt{6}},$$

$$K = - \left| (\nabla_\mu - igA_\mu^{\text{str}}) \sigma \right|^2 - |(\nabla_\mu p)|^2 - |(\nabla_\mu a)|^2 \\ - |(\nabla_\mu b)|^2 - |(\nabla_\mu S)|^2 - \frac{\text{Tr}(\tau_{\text{str}}^2)}{4} F_{\mu\nu}^{\text{str}} F^{\mu\nu \text{ str}}$$

Ansatz for the string

- **Ansatz**

$$\begin{cases} p = p(r) \\ a = a(r) \\ b = b(r) \\ \sigma = f(r)e^{in\theta} \\ A_\mu = A_\theta^{\text{str}}(r)\tau^{\text{str}}\delta_\mu^\theta \\ S = S(r) \end{cases}$$

- **Boundary conditions**

► $r \rightarrow \infty:$ $\begin{cases} f, a, b, p, S \rightarrow \text{SM VEVs} \\ A_\theta = n \end{cases}$

- **Boundary conditions**

► $r \rightarrow \infty:$ $\begin{cases} f, a, b, p, S \rightarrow \text{SM VEVs} \\ A_\theta = n \end{cases}$

► $r = 0:$ $\begin{cases} f = A_\theta = 0 \\ \frac{dx}{dr} = 0 \text{ for } a, b, p \text{ and } S \end{cases}$

Outline

- 1 Why a realistic model of cosmic strings**
- 2 Theoretical framework**
- 3 Abelian cosmic strings**
- 4 Microscopic and macroscopic structure**

Numerical solution

- **Numerical Implementation**
 - ▶ Successive overrelaxed method
 - ▶ On a finite range lattice
 - ▶ Around 2000 configurations studied

- Range of GUT parameters

$$W = \frac{m}{2} \Phi^2 + m_\Sigma \Sigma \bar{\Sigma} + \frac{\lambda}{3} \Phi^3 + \eta \Phi \Sigma \bar{\Sigma} + \kappa S (\Sigma \bar{\Sigma} - M^2).$$

$$\text{with } \begin{cases} \frac{m}{M} = \frac{m_\Sigma}{M} \in \{1, 20\} \\ \kappa \in \{10^{-2}, 30\} \\ \eta \in \{10^{-2}, 10\} \\ \lambda \in \{10^{-2}, 1\} \end{cases}$$

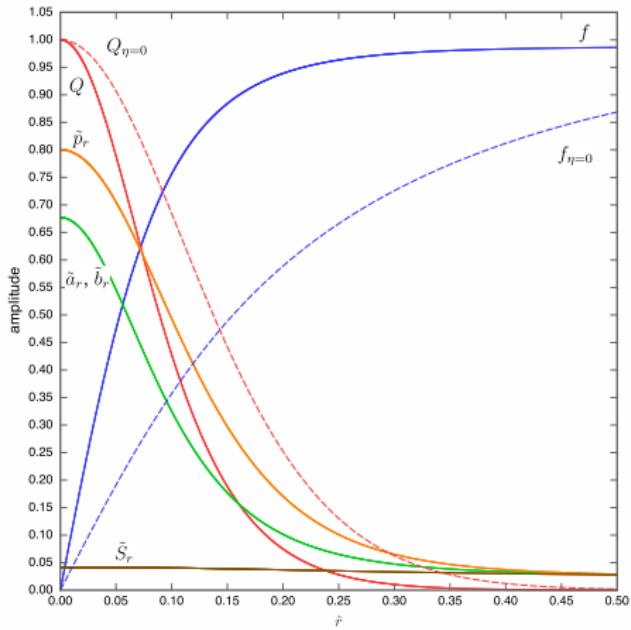
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- Toy-model limit: $\eta \rightarrow 0$

Microscopic structure



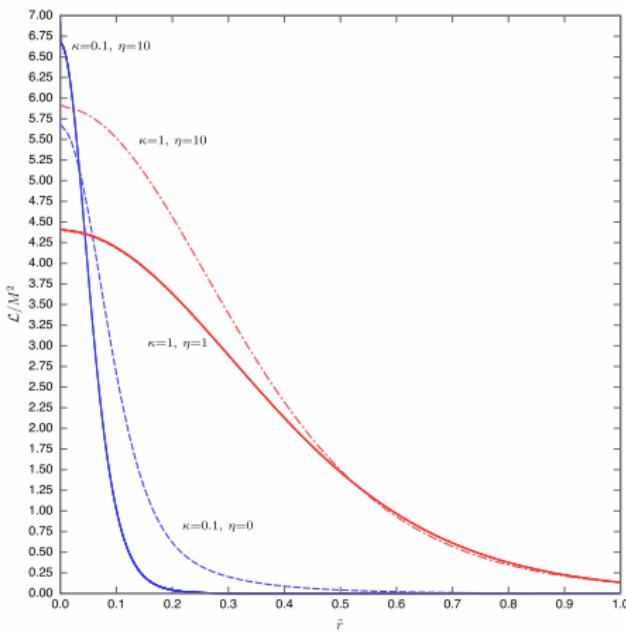
$X(r)$

- $\eta = 10$, $\kappa \in \{10^{-2}, 30\}$

$$W \supset \eta \Phi \Sigma \bar{\Sigma} + \kappa S(\Sigma \bar{\Sigma} - M^2)$$

- Lagrangian density

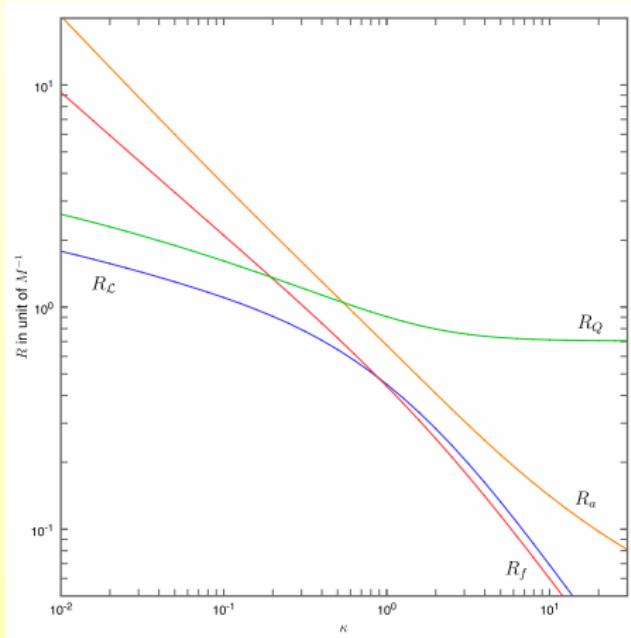
$$W \supset \eta \Phi \Sigma \bar{\Sigma} + \kappa S (\Sigma \bar{\Sigma} - M^2)$$



$$\mathcal{L}(r)$$

- Characteristic radiiuses

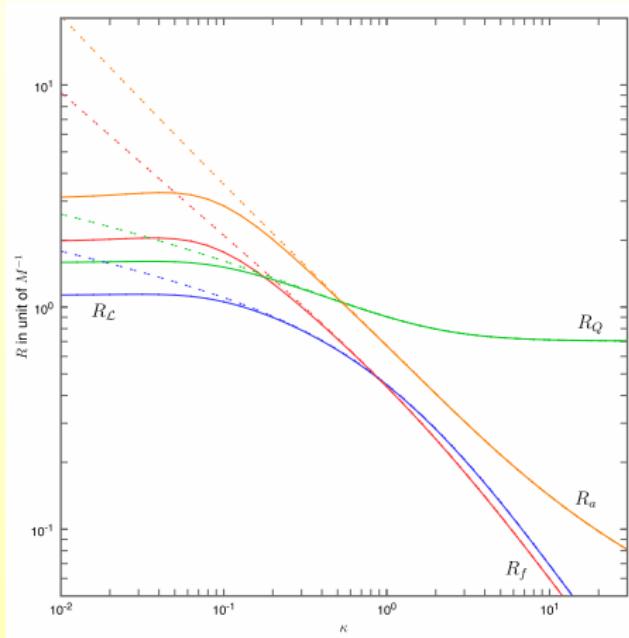
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R_X

- Characteristic radiiuses

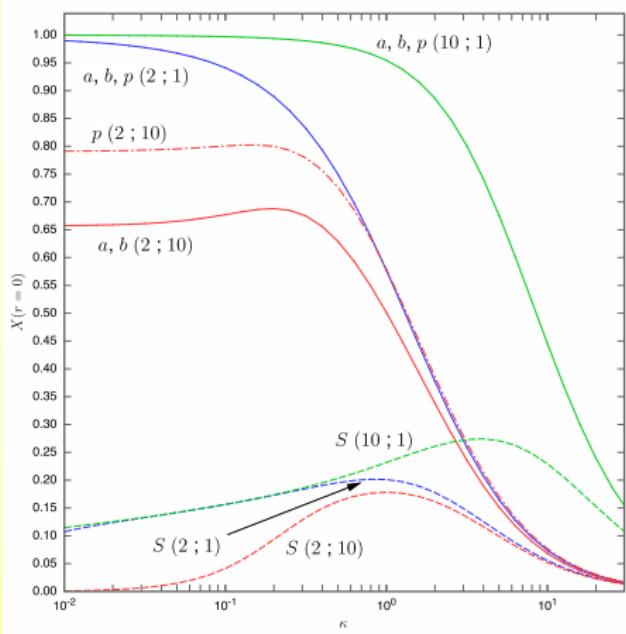
$$W \supset \eta \Phi \Sigma \bar{\Sigma} + \kappa S (\Sigma \bar{\Sigma} - M^2)$$



R_X

- Condensation in the string

$$W \supset \eta \Phi \Sigma \bar{\Sigma} + \kappa S(\Sigma \bar{\Sigma} - M^2)$$

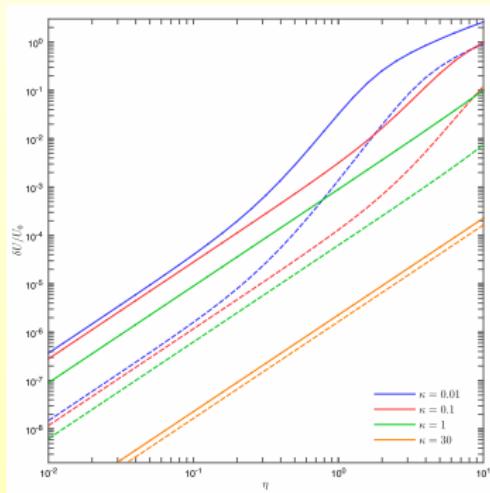


$$X(r = 0) \text{ for } (m/M, \eta)$$

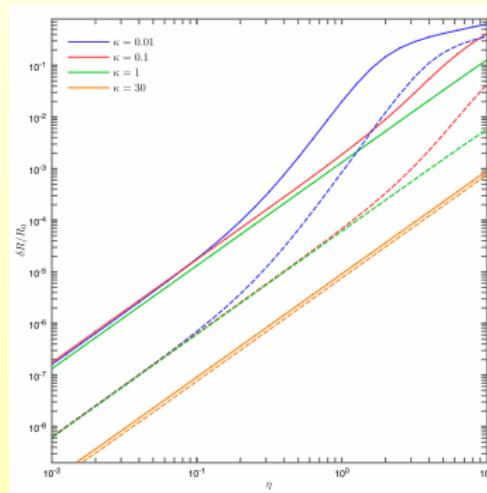
Macroscopic structure

- Influence of η

$$W \supset \eta \Phi \Sigma \bar{\Sigma} + \kappa S (\Sigma \bar{\Sigma} - M^2)$$



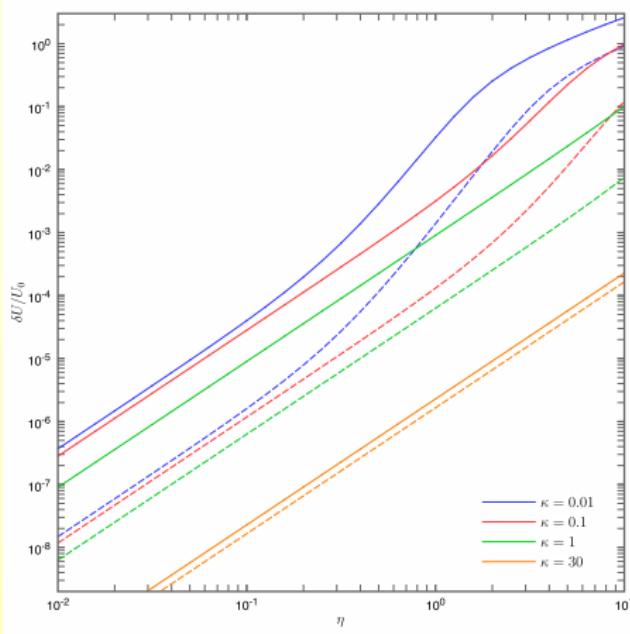
$$\delta U/U_0$$



$$\delta R/R_0$$

• Influence of η

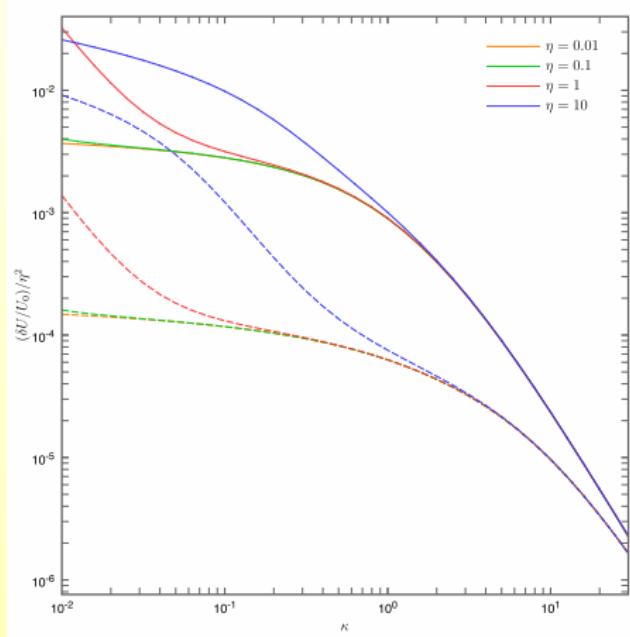
$$W \supset \eta \Phi \Sigma \bar{\Sigma} + \kappa S (\Sigma \bar{\Sigma} - M^2)$$



$$\delta U/U_0$$

- **Influence of κ**

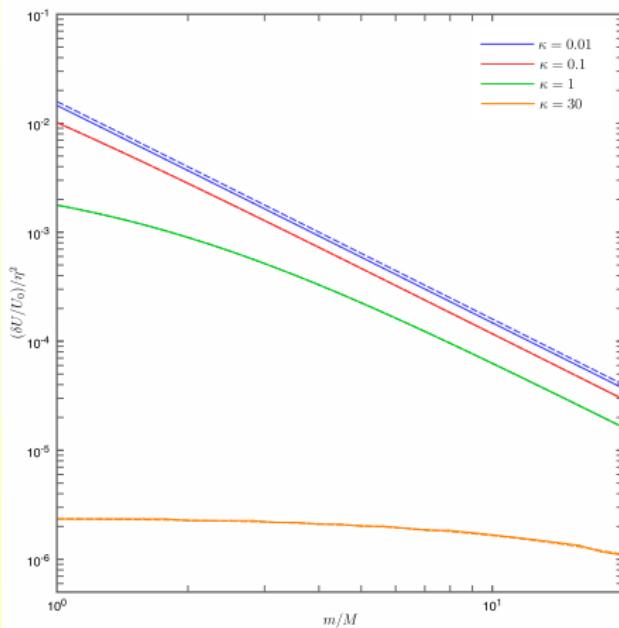
$$W \supset \eta \Phi \Sigma \bar{\Sigma} + \kappa S (\Sigma \bar{\Sigma} - M^2)$$



$$\delta U / U_0 / \eta^2$$

- Influence of $\frac{m}{M} = \frac{m_\Sigma}{M}$

$$W \supset \eta \Phi \Sigma \bar{\Sigma} + \kappa S (\Sigma \bar{\Sigma} - M^2)$$



$$\delta U / U_0 / \eta^2$$

Conclusion

• Results

- ▶ Realistic description of cosmic strings
- ▶ Structure and energy per unit length function of the GUT parameters
- ▶ Perturbative description \Rightarrow two different classes of model

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- ▶ Bosonic currents
- ▶ Intercommutation and reconnection \Rightarrow network evolution
 \Rightarrow CMB observation

Conclusion

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- **Thanks for your attention !**