On-shell methods for classical and quantum gravity

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based on the works 1309.0804, 1410.4148 and 1410.7590, 1609.07477 1704.01624

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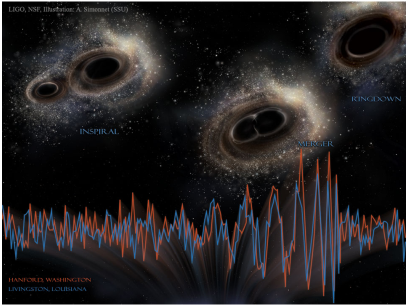
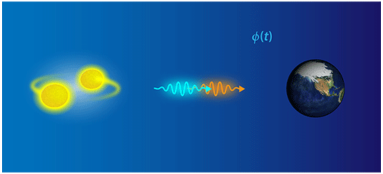
A new observational window on gravitation



The detection of gravitational waves has open a new window on the gravitational physics of our universe

- For the first time detection and test of GR in the strong gravity coupling regime
- For the first time dynamics of Black hole (not just static object curving space-time)

A new observational window on gravitation



[Yunes, Yagi, Pretorius] have listed theoretical implications of GW150914 in particular

GW150914 constrains a number of theoretical mechanisms that modify GW propagation

Quantum gravity as an effective field theory

[Donoghue] has explained that one can evaluate some long-range infra-red contributions in any quantum gravity theory and obtain reliable answers

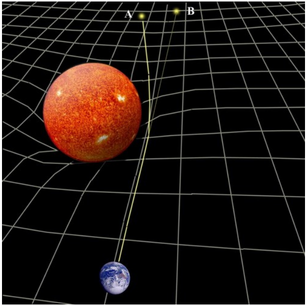
Some physical properties of quantum gravity are *universal* being independent of the UV completion

The one-loop infra-red contributions depend only the structure the low-energy fields and the classical background

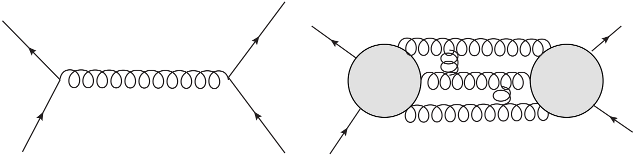
Physics of the effective field theory approach

Using the effective field theory approach to gravity one can compute

- the classical (post-Newtonian) and quantum contributions to the gravitational potential between masses
- Quantum corrections to the bending angle of massless particle by a massive classical object



Classical physics from loops



We will be considering the pure gravitational interaction between massive and massless matter of various spin

$$\mathcal{L}_{\text{EH}} \sim \int d^4x \left(-\frac{2}{\kappa^2}\,\mathcal{R} + \kappa h_{\mu\nu} T^{\mu\nu}_{\text{matter}} \right) ,$$

We will be considering perturbative computations $\kappa^2 = 32\pi G_N$

$$\mathfrak{M} = \frac{1}{\hbar}\,\mathfrak{M}^{\text{tree}} + \hbar^0\,\mathfrak{M}^{1-\text{loop}} + \cdots .$$

Double expansion : classical and quantum parameters

We have two scales in the problem:

- The Schwarzschild radius

$$r_S = \frac{2G_N M}{c^2}$$

- The Compton wave-length

$$\lambda\!\!\!^{-} = \frac{\hbar}{Mc}$$

- Dual with respect to the Planck length

$$r_S \lambda = \frac{2G_N \hbar}{c^3} = 2\, \ell_P^2$$

Double expansion : classical and quantum contributions

Starting from the PPN expansion

$$V^{\text{class}}(r) = \sum_{m\geqslant 0} v_{m,0} \left(\frac{r_S}{r}\right)^m$$

If $\lambda = \hbar/(Mc)$ is the characteristic length of the quantum fluctuations we have at first order

$$\frac{1}{(r \pm \lambdabar)^n} \simeq \frac{1}{r^n} \pm \frac{\lambdabar}{r^{n+1}}$$

leading to the modified potential

$$V(r \pm \lambda) \simeq \sum_m \left(v_{m,0} \frac{r_S^m}{r^m} + v_{m,1} \frac{r_S^m \lambda}{r^{m+1}} \right)$$

Double expansion : classical and quantum contributions

Since

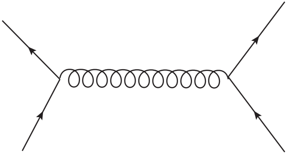
$$r_S\lambda = \frac{2G_N\hbar}{c^3} = 2\,\ell_P^2$$

We have

$$V(r\pm\lambda) \simeq \sum_m \left(v_{m,0}\frac{r_S^m}{r^m} + v_{m,1}\frac{r_S^{m-1}\ell_P^2}{r^{m+1}} \right)$$

This motivates the appearance of the first quantum corrections to the gravitational potential We will use scattering amplitudes to evaluate both the classical and the quantum part of the long range potential.

Classical physics at tree-level



The tree-level contribution is the 1-graviton exchange giving the classical Newtonian potential in the non-relativistic limit

$$\mathfrak{M}^{\text{tree}} \propto G_N \frac{(m_1 m_2)^2}{\vec{q}^2}$$

The potential is obtained by

$$V(r) = \int \frac{d^3\vec{q}}{(2\pi)^3} \frac{1}{4m_1 m_2} \mathfrak{M}(\vec{q})\, e^{i\vec{q}\cdot\vec{r}}$$

Classical physics from loops

For a 2 body scattering an L-loop gravitational amplitudes has to following dependence in $D=4$ dimensions

$$[\mathfrak{M}^{L-\text{loop}}] = \kappa^{2L+2}\,\hbar^{L-1}\,\Lambda^{2L+2}$$

The classical piece will manifest itself through contributions of the form

$$\mathfrak{M}^{L-\text{loop}}\Big|_{\text{classical}} = m^4\,\kappa^{2L+2}\,\hbar^{L-1}\,(q^2)^{L-1}\,\left(\frac{m}{\hbar c\,\sqrt{q^2}}\right)^L = \frac{1}{\hbar}\frac{m^3}{q^2}\,\left(r_S\,\sqrt{q^2}\right)^L$$

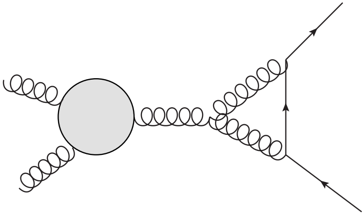
Which after Fourier transform in $D=3$ gives

$$V(r) = \frac{1}{m^2}\int \mathfrak{M}^{L-\text{loop}}\,e^{iq\cdot r}\,d^3q = \frac{1}{\hbar}\,\left(\frac{r_S}{r}\right)^L$$

This how the Schwarzschild radius arises inside quantum loop amplitudes

Classical physics from loops

Let's consider the one-loop contribution for a say a massive scalar of mass m



Putting back the factors of $\hbar$ and c the Klein-Gordon equation reads

$$(\Box - \underbrace{\frac{m^2c^2}{\hbar^2}}_{\lambda\!\!\!^{-}{}^{-2}})\phi = 0$$

Classical physics from loops

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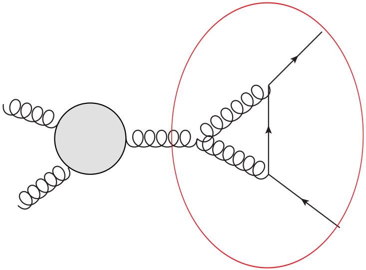

The triangle contribution with a massive leg $p_1^2 = p_2^2 = m^2$ reads

$$\int \frac{d^4\ell}{(\ell+p_1)^2(\ell^2-\frac{1}{\hbar^2})(\ell-p_2)^2}\Bigg|_{\text{finite part}} \sim \frac{1}{m^2}\left(\log(s)+\frac{\pi^2}{\hbar\sqrt{s}}\right)$$

Classical physics from loops

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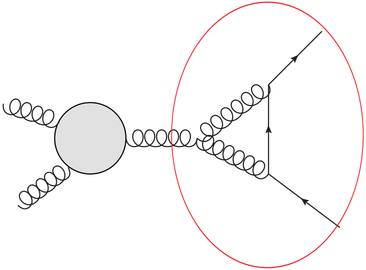


The triangle contribution with a massive leg $p_1^2 = p_2^2 = m^2$ reads

$$\frac{\kappa^2}{2\lambda\sqrt{s}} = \frac{r_S}{\sqrt{s}}$$

Fourier transformed with respect to the non-relativistic momentum transfert $|\vec{q}| = \sqrt{s}$ leads to r_S/r corrections

Classical physics from loops

The $1/\hbar$ term at one-loop contributes to the *same* order as the classical tree term [Donoghue; Bjerrum-Bohr, Donoghue, Holstein; Donoghue, Holstein; Bjerrum-Bohr, Donoghue, Vanhove]

$$\mathfrak{M} = \frac{1}{\hbar}\left(\frac{G_N(m_1m_2)^2}{\vec{q}^2} + \frac{G_N^2(m_1m_2)^2(m_1+m_2)}{|\vec{q}|} + \cdots\right) + \hbar^0 G_N^2\, O(\log(\vec{q}^2)) + \cdots$$

For the scattering between a massive matter of mass m and massless matter of energy E one gets

$$\mathfrak{M} \sim \frac{1}{\hbar}\left(G_N\frac{(mE)^2}{\vec{q}^2} + G_N^2\frac{m^3E^2}{|\vec{q}|}\right) + \hbar\, G_N^2 O\left(\log(\vec{q}^2), \log^2(\vec{q}^2)\right).$$

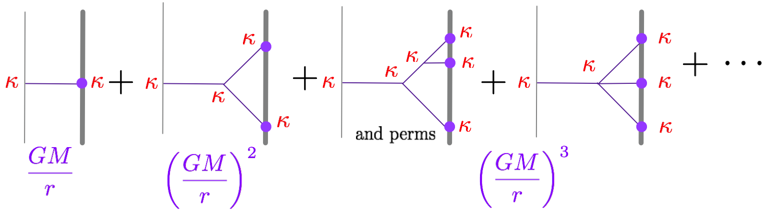
The mechanisms generalizes to higher loop-order amplitudes to leads to the higher order post-Newtonian corrections

Classical metric [Bjerrum-Bohr, Damgaard, Festuccia, Planté, Vanhove; to appear]

By considering the graviton emission $\langle p_1 | T_{\mu\nu} | p_2 \rangle$ one can obtain the metric (in harmonic gauge) by extracting the classical pieces from higher loop amplitudes

$$g_{00} = \frac{1 - \frac{GM}{r}}{1 + \frac{GM}{r}} = 1 + 2 \sum_{k \geqslant 1} (-1)^k \left(\frac{GM}{r} \right)^k$$

The kth order term is a $k-1$-loop term



The tree skeleton graphs are the one computed by `[Duff, PRD(1973)]`

Corrections to Newton's potential

One-loop corrections to Newton's potential can be calculated using effective field theory approach to gravity [Donoghue; Bjerrum-Bohr, Donoghue, Holstein; Bjerrum-Bohr, Donoghue, Vanhove]

$$V(r) = -\frac{G_N m_1 m_2}{r}\left(1 + C\,\frac{G_N(m_1+m_2)}{r} + Q\frac{G_N\hbar}{r^2}\right) + Q' G_N^2 m_1 m_2 \delta^3(\vec{x})$$

- C is the classical correction and Q and Q' are quantum corrections

Corrections to Newton's potential

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- C is the classical correction and Q and Q' are quantum corrections
- Q in the potential $V(r)$ is ambiguous but $V(r)$ is not observable

The coefficients of $1/\sqrt{-q^2}$ and $\log(-q^2)$ in the amplitude are unambiguously defined and depend on the long range physics

Corrections to Newton's potential

One-loop corrections to Newton's potential can be calculated using effective field theory approach to gravity [Donoghue; Bjerrum-Bohr, Donoghue, Holstein; Bjerrum-Bohr, Donoghue, Vanhove]

$$\mathfrak{M}^{1-\text{loop}}(q^2) = \frac{G_N(m_1m_2)^2}{q^2} + C\frac{G_N^2(m_1m_2)^2(m_1+m_2)}{|q|} + \hbar\left(QG_N^2(m_1m_2)^2\log(-q^2) + Q'G_N^2(m_1m_2)^2 Q'G_N^2(m_1m_2)^2\right)$$

- Q' is the short distance UV divergences of quantum gravity: need to add the R^2 term ['t Hooft-Veltman]
$$S = \int d^4x|-g|^{\frac{1}{2}}\left[\frac{2}{32\pi G_N}\mathcal{R} + c_1\mathcal{R}^2 + c_2 R_{\mu\nu}R^{\mu\nu} + \cdots\right]$$

Effective field theory and gravity

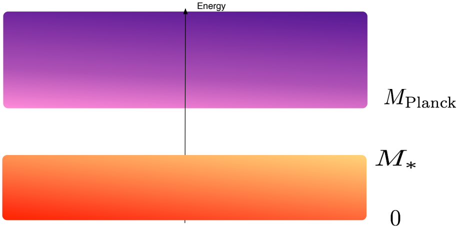
We are working in the context of an effective field theory assuming :

- standard QFT (local, unitary, lorentz invariant, . . .)
- The low-energy DOF: graviton, usual matter fields
- Standard symmetries: General relativity as we know it

We will allow in the lagrangian only the low energy DOF with higher derivative term suppressed by some high energy scale M_*

$$S_{eff} = \int_{\mathcal{M}_4} d^4x \, \sqrt{g} \, \left(\Lambda^4 + \frac{M_{Pl}^2}{2} \mathcal{R}_{(4)} + c_0 \, R^2 + \sum_{k \geqslant 1} d_k \left(\frac{R}{M_*^2} \right)^{2+k} \right)$$

Effective field theory and gravity



The parameters may be derived from a fundamental microscopic theory but in EFT treatment we can work without knowing the explicit relations

Recovering classical General relativity

The classical contribution from the matter scattering *must* reproduce the post-Newtonian contributions to the gravitational potential derived from General relativity

$$V^{\rm GR}(r) = -\frac{G_N m_1 m_2}{r}\left(1 + 3G_N\frac{m_1+m_2}{r}\right)$$

When scattering massless matter we can consider the bending angle on the massive object of Schwarzschild radius $r_S = 2G_N m/c^3$

$$\theta^{\rm GR} = \frac{2r_S}{b} + \frac{15\pi}{16}\left(\frac{r_S}{b}\right)^2$$

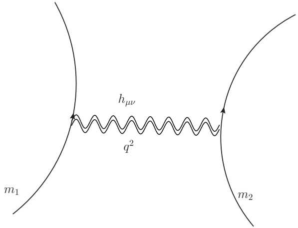
We need to see that the classical pieces from the loop match the general relativity answer (i.e. satisfy the equivalence principle)

We can determine quantum corrections to gauge invariant quantities

Perturbative technics

Classical Newton's potential is obtained in the non-relativistic limit

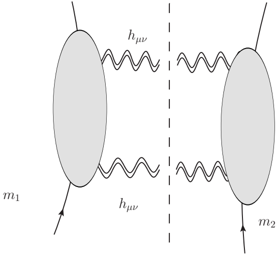
$$V(|\vec{q}|) = \frac{G_N m_1 m_2}{\vec{q}^2} \qquad V(r) = -\frac{G_N m_1 m_2}{r}$$



is derived by a tree-level graph exchanging a graviton

Loop amplitude

Since we are only interested in the long range (low energy) graviton exchange, it is enough to just evaluate the gravitons cut

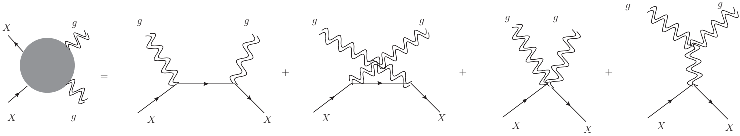


we need to know the gravitational Compton amplitudes on a particle of spin s with mass m

$$X^{s,m} + \text{graviton} \to X^{s,m} + \text{graviton}$$

Gravitational compton scattering

Gravitational Compton scatting off a massive particle of spin $s = 0, \frac{1}{2}, 1$



using Feynman rules and DeWitt or Sannan's 3- and 4-point vertices this is a big mess but this will be simplified using the momentum kernel formalism to gravity amplitude

The Momentum Kernel formalism Gravity amplitude

The KLT relation allow to express the field theory multi-particle tree-level amplitudes as bilinear of color ordered Yang-Mills amplitudes

$$\mathfrak{M}_n^{\text{tree}} = (-1)^{n-3} \sum_{\sigma,\gamma \in \mathfrak{S}_{n-3}} \mathcal{S}[\gamma(2,\dots,n-2)|\sigma(2,\dots,n-2)]_{k_1}$$
$$\times \mathcal{A}_n(1,\sigma(2,\dots,n-2),n-1,n)\widetilde{\mathcal{A}}_n(n-1,n,\gamma(2,\dots,n-2),1)$$

The color ordered Yang-Mills amplitudes satisfy the annihilation relation $\forall \beta \in \mathfrak{S}_{n-2}$

$$\sum_{\sigma \in \mathfrak{S}_{n-2}} \mathcal{S}(\sigma(2,\dots,n-1)|\beta(2,\dots,n-1))|_{k_1} \mathcal{A}(1,\sigma(2,\dots,n-1),n) = 0$$

[Bern, Carrasco, Johansson] [Kawai,Lewellen, Tye; Tye, Zhang;Bjerrum-Bohr, Damgaard, Feng, Søndergaard; Bjerrum-Bohr, Damgaard, Søndergaard, Vanhove; Stieberger]

The Momentum kernel in field theory

The $\alpha' \to 0$ limit of the monodromy relations between string theory amplitudes lead to an object named momentum kernel $\mathcal{S}$

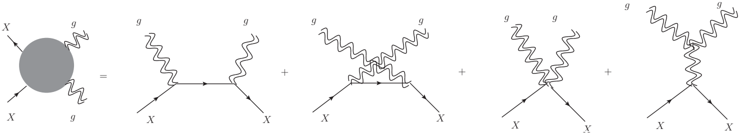
$$\mathcal{S}[i_1,\ldots,i_k|j_1,\ldots,j_k]_p := \prod_{t=1}^{k}\left(p\cdot k_{i_t} + \sum_{q>t}^{k}\theta(t,q)\, k_{i_t}\cdot k_{i_q}\right)$$

$\theta(t,q) = 1$ if $(i_t - i_q)(j_t - j_q) < 0$ and 0 otherwise

[Bern, Carrasco, Johansson; Bjerrum-Bohr, Damgaard, Vanhove; Stieberger; Mafra, Schlotterer]

[Bjerrum-Bohr, Damgaard, Feng, Søndergaard; Bjerrum-Bohr, Damgaard, Søndergaard, Vanhove]

Gravitational compton scattering

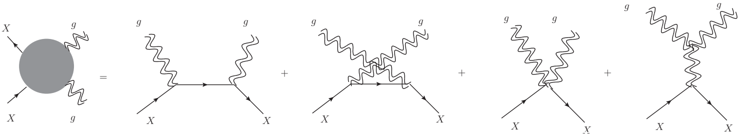


We express the gravity Compton scattering as a product of two Yang-Mills amplitudes [Kawai, Lewellen, Tye], [Bern, Carrasco, Johansson]

$$\mathfrak{M}(X^s g \to X^s g) = G_N \times (p_1 \cdot k_1)\, \mathcal{A}_s(1234)\tilde{\mathcal{A}}_0(1324)$$

$\mathcal{A}_s(1234)$ is the color ordered amplitudes scattering a gluon off a massive spin s state $X^s g \to X^s g$

Gravitational compton scattering

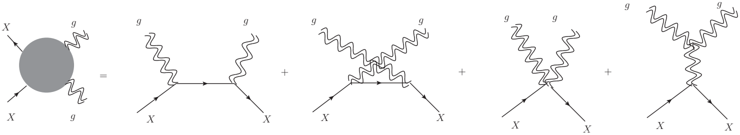


We express the gravity Compton scattering as a product of two QED Compton amplitudes using the monodromy relations [Bjerrum-Bohr, Donoghue, Vanhove]

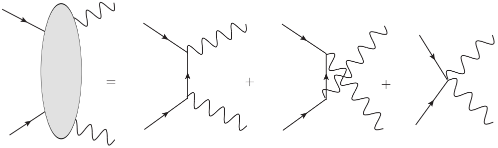
$$(k_1 \cdot k_2)\, \mathcal{A}_s(1234) = (p_1 \cdot k_2)\, \mathcal{A}_s(1324)$$

$$\mathfrak{M}(X^s g \to X^s g) = G_N \frac{(p_1 \cdot k_1)(p_1 \cdot k_2)}{k_1 \cdot k_2}\, \mathcal{A}_s(1324) \tilde{\mathcal{A}}_0(1324)$$

Gravitational compton scattering

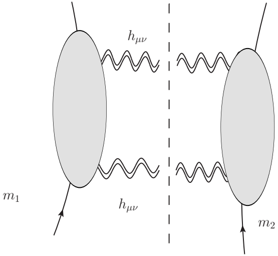


The gravity Compton scattering is expressed as the square of QED (abelian) Compton amplitudes [Bjerrum-Bohr, Donoghue, Vanhove]



$$\mathfrak{M}(X^s g \to X^s g) = G_N \frac{(p_1 \cdot k_1)(p_1 \cdot k_2)}{k_1 \cdot k_2} \mathcal{A}_s(1324)\tilde{A}_0(1324)$$

The one-loop amplitude between massive particles



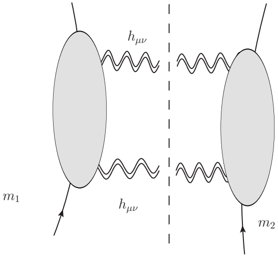
We are only interested in the $1/\sqrt{-q^2}$ and $\log(-q^2)$ terms since the terms of $(q^2)^n/\sqrt{-q^2}$ and $(q^2)^n \log(-q^2)$ are negligible in the non-relativistic limit.
Only the massless graviton cut is enough.

The cut contributions

$$\mathfrak{M}^{1-\text{loop}}|_{\text{singlet cut}} = \int \frac{d^{4-2\epsilon}\ell}{\ell_1^2 \ell_2^2 \prod_{i=1}^4 \ell_1 \cdot p_i}$$

$$\mathfrak{M}^{1-\text{loop}}|_{\text{non-singlet cut}} = \int d^{4-2\epsilon}\ell \frac{\Re\text{e}\left(\text{tr}_-(\not\ell_1 \not p_1 \not\ell_2 \not p_2)\right)^4}{\ell_1^2 \ell_2^2 \prod_{i=1}^4 \ell_1 \cdot p_i}$$

The one-loop amplitude between massive particles



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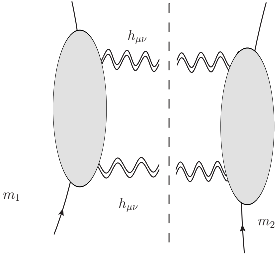
Only the massless graviton cut is enough.

- In the non-relativistic limit the amplitude decomposes

$$\mathfrak{M}^{1-\text{loop}} \simeq G_N^2 (m_1m_2)^4 (I_4(s,t) + I_4(s,u)) + G_N^2 (m_1m_2)^3 s (I_4(s,t) - I_4(s,u)) + G_N^2 (m_1m_2)^2 (I_3(s,m_1) + I_3(s,m_2)) + G_N^2 (m_1m_2)^2 I_2(s)$$

The one-loop amplitude between massive particles



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Only the massless graviton cut is enough.

The result is given by

$$\mathfrak{M}^{1-\text{loop}} \simeq G_N^2 (m_1 m_2)^2 \left(\underbrace{6\pi}_{C} \frac{m_1+m_2}{\sqrt{-q^2}} \underbrace{-\frac{41}{5}}_{Q} \log(-q^2) \right)$$

Universality of the result

In the case of scattering of particles of different spin S_1 and S_2 the non-relativistic potential reads

$$\mathfrak{M}^{1-\text{loop}}(q^2) \simeq G_N^2 (m_1 m_2)^2 \left(C\, \frac{(m_1+m_2)}{\sqrt{-q^2}} + Q\hbar \log(-q^2) \right)$$

C and Q have a *spin-independent* and a *spin-orbit* contribution

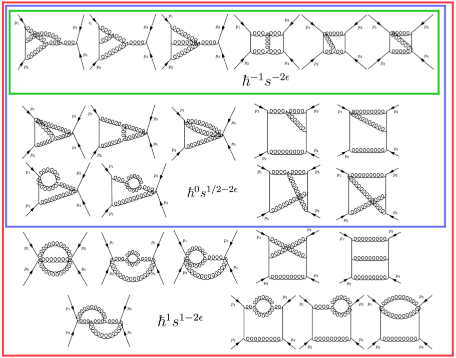
$$C, Q = C, Q^{S-I} \langle S_1|S_1\rangle \langle S_2|S_2\rangle + C, Q^{S-O}_{1,2} \langle S_1|S_1\rangle\, \vec{S}_2 \cdot \frac{p_3 \vec{\times} p_4}{m_2} + (1 \leftrightarrow 2)$$

This expression is generic for all type of matter

the numerical coefficients are the same for all matter type
The universality of the coefficients with respect to the spin of the external states is a consequence of

- The reduction to the product of QED amplitudes
- the low-energy theorems of [Low, Gell-Mann, Goldberger] and [Weinberg]

The two-loop amplitude [Planté, Phd thesis]

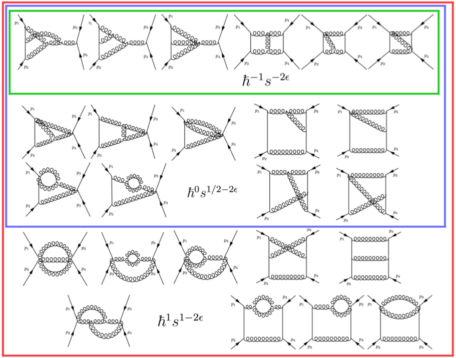


Only the green box contributes to the classical 2PN effective potential

$$\delta V^{2-\text{loop}} = -\frac{17G^3 m_1 m_2 (m_1^2 + m_2^2)}{4r^3} + \frac{3G^3 m_1^2 m_2^2}{2r^3}$$

The blue box to the first quantum correction and the red box to two-loop quantum corrections

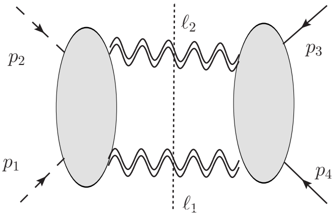
The two-loop amplitude [Planté, Phd thesis]



The green box gives the complete classical 2PN effective radial action integral for a test mass $\frac{S}{G_N m_1 m_2} = -Et + h\varphi + \int dr\,\sqrt{R(r,E,h)}$

$$R(r,E,0) = 2E + \frac{8EG_NM}{c^2r} + \frac{(6c^2+15E)(G_NM)^2}{c^4r^2} + \frac{17(G_NM)^3}{2c^4}\frac{1}{r^3} + O(E)$$

The one-loop amplitude for massless particles



We consider the gravitational one-loop amplitude between a massless particle of spin S and a massive scalar

$$\begin{aligned}\kappa^{-4}\, i\mathfrak{M}_S^{1-\text{loop}} &= bo^S(s,t)\, I_4(s,t) + bo^S(s,u)\, I_4(s,u)\\ &+ t_{12}^S(s)\, I_3(s,0) + t_{34}^S(s)\, I_3(s,M^2)\\ &+ bu^S(s,0)\, I_2(s,0)\,.\end{aligned}$$

The coefficients satisfy interesting BCJ relations

$$\frac{bo^S(s,t)}{t-M^2} + \frac{bo^S(s,u)}{u-M^2} = t_{12}^S(s)$$

The amplitude

The low-energy approximation

$$
\begin{aligned}
i\mathfrak{M}_S^{\text{tree}+1-\text{loop}} &= \frac{\mathcal{N}^{(S)}}{\hbar}\Big[\kappa^2\,\frac{(2M\omega)^2}{16q^2} \\
&+ \hbar\frac{\kappa^4}{16}\Big(4(M\omega)^4(I_4(t,s)+I_4(t,u))+3(M\omega)^2 sI_3(t) \\
&- \frac{15}{4}(M^2\omega)^2 I_3(t,M)+bu^S(M\omega)^2 I_2(t)\Big)\Big]
\end{aligned}
$$

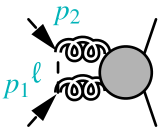
For photon scattering only the amplitudes with helicity $(++)$ and $(--)$ are non-vanishing.

Therefore there is no birefringence effects to contrary to case with electrons loops contributing to the interaction [Drummond, Hathrell;Berends, Gastmans]

The amplitude

$$
\begin{aligned}
i\mathfrak{M}_S^{\text{tree}+1-\text{loop}} &\simeq \frac{\mathcal{N}^{(S)}}{\hbar}\,\frac{(M\omega)^2}{4}\\
\times\quad & \Big[\frac{\kappa^2}{q^2} + \kappa^4\frac{15}{512}\frac{M}{\sqrt{-q^2}}\\
+\quad & \hbar\kappa^4\frac{15}{512\pi^2}\,\log\left(\frac{-q^2}{M^2}\right) - \hbar\kappa^4\,\frac{bu^S}{(8\pi)^2}\,\log\left(\frac{-q^2}{\mu^2}\right)\\
+\quad & \hbar\kappa^4\frac{3}{128\pi^2}\,\log^2\left(\frac{-q^2}{\mu^2}\right) + \kappa^4\,\frac{M\omega}{8\pi}\frac{i}{s}\log\left(\frac{-q^2}{M^2}\right)\Big]
\end{aligned}
$$

The last line contains the infrared divergences



$$\propto \int_0 \frac{d^{4-2\epsilon}\ell}{\ell^2\,2\ell\cdot p_1\,2\ell\cdot p_2} \sim \frac{(t/\mu^2)^{-\epsilon}}{\epsilon^2\,t}$$

The bending angle via Eikonal approximation

$$i\mathcal{M}(\boldsymbol{b}) \simeq 2(s - M^2)\left[e^{i(\chi_1+\chi_2)} - 1\right]$$

$\chi_1(\boldsymbol{b})$ is the Fourier transform of the one graviton (tree-level) exchange

$$\chi_1(\boldsymbol{b}) = \frac{1}{2M2E}\int \frac{d^2\boldsymbol{q}}{(2\pi)^2}\, e^{-i\boldsymbol{q}\cdot\boldsymbol{b}}\, \mathfrak{M}_S^{(1)}(\boldsymbol{q}) \simeq 4G_N ME\left[\frac{1}{d-2} - \log(b/2) - \gamma_E\right]$$

$\chi_2(\boldsymbol{b})$ is the Fourier transform of the two gravitons (one-loop) exchange

$$\chi_2(\boldsymbol{b}) = \frac{1}{2M2E}\int \frac{d^2\boldsymbol{q}}{(2\pi)^2}\, e^{-i\boldsymbol{q}\cdot\boldsymbol{b}}\, \mathfrak{M}_X^{(2)}(\boldsymbol{q})$$

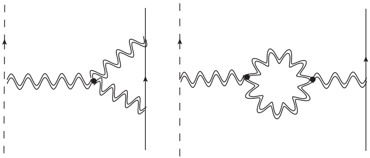
$$\simeq -G_N^2 M^2 E\frac{15\pi}{4b} - \frac{G_N^2 M^2 E}{2\pi b^2}\left(8bu^S + 9 - 48\log\frac{b}{2b_0}\right).$$

The bending angle

The bending angle $\theta_S \simeq -\frac{1}{E}\frac{\partial}{\partial b}\left(\chi_1(b)+\chi_2(b)\right)$ is

$$\theta_S \simeq \frac{4GM}{b} + \frac{15}{4}\frac{G^2M^2\pi}{b^2} + \frac{8bu^S + 9 - 48\log\frac{b}{2b_0}}{\pi}\,\frac{G^2\hbar M}{b^3}\,.$$

- The classical contribution including the 1rst Post-Newtonian correction is correctly reproduced
- The quantum corrections are new: not only from a quantum corrected metric



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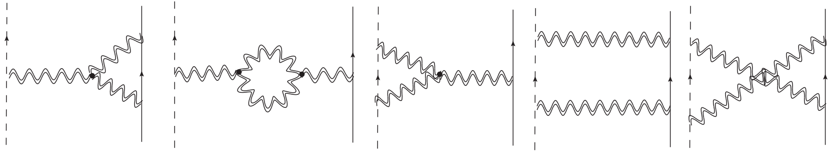

The difference between the bending angle for a massless photon and massless scalar

$$\theta_\gamma - \theta_\varphi = \frac{8(bu^\gamma - bu^\varphi)}{\pi}\,\frac{G^2\hbar M}{b^3}\,.$$

Outlook

Recent progresses from string theory technics, on-shell unitarity, double-copy formalism simplifies a lot perturbative gravity amplitudes computations

- The amplitudes relations discovered in the context of massless supergravity theories extend to the pure gravity case with massive matter
- The use of quantum gravity as an effective field theory allows to compute universal contributions from the long-range corrections
- We can reproduce the classical GR post-Newtonian corrections to the potential and understand some generic properties using low-energy theorems: hope to be able to simplify the computation of PPN corrections.