

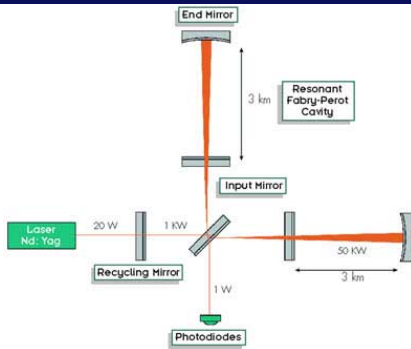
Continuous gravitational waves as probes of neutron stars and dark matter

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1. Overview
2. Neutron stars
3. Dark matter around black holes
4. Planetary-mass primordial black holes
5. Conclusions



- Gravitational waves will change the positions of objects in their paths
- The mirrors move differently, which induce a differential strain on the detector $h(t)$
- Interferometers are most sensitive to the phase of the modulation, as opposed to the amplitude modulation



Sources of gravitational waves

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Overview

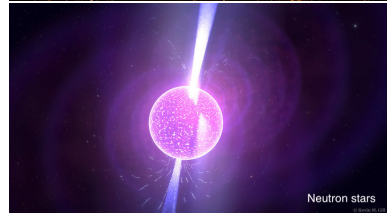
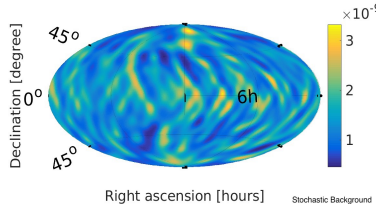
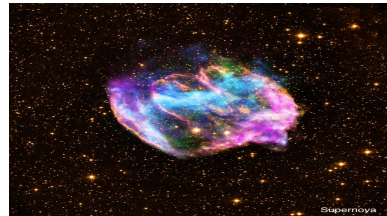
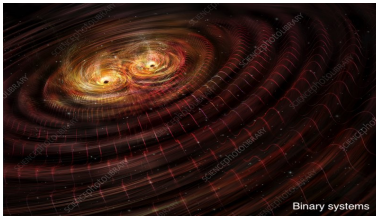
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- Four major sources that LIGO/Virgo search for
- Merging binary systems have already been seen, but the other sources have eluded detection!



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Focus here on continuous waves

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Overview

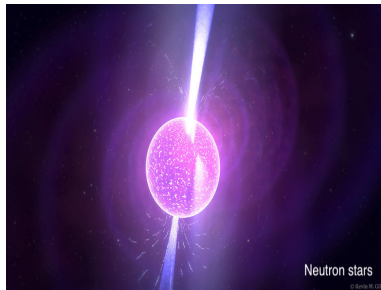
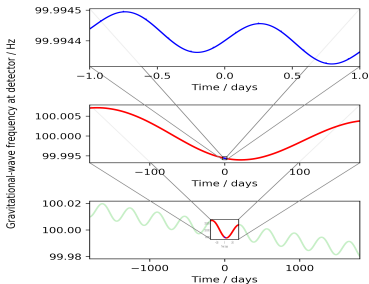
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- Quasi-monochromatic, quasi-infinite duration signals
- Searches for both known and unknown neutron stars are difficult and computationally demanding
- From a data analysis point of view, many signals with different underlying physics follow this signal model

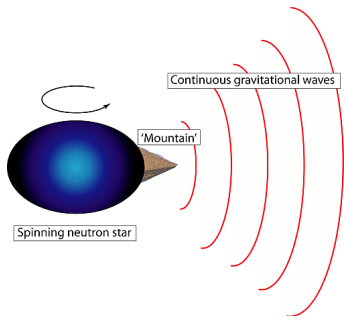


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Deformed neutron stars

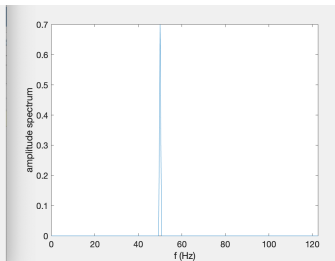
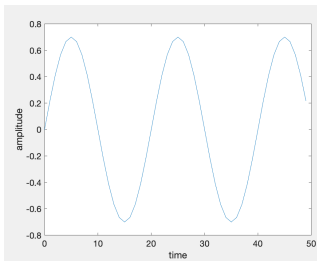
- Small deformation and rotation \rightarrow gravitational waves
- Rotate at $\sim 10 - 1000$ Hz [18]
- Small spin-down or spin-up \dot{f} of $\leq \mathcal{O}(10^{-8})$ Hz/s
- Model frequency evolution as Taylor series expansion to first or second order
- Reasons for deformations:
 - Crustal strain (starquake, formed at birth)
 - Strong *internal* magnetic field buried during accretion [15]





Principles of data analysis for known pulsars

- The position, rotational frequency and frequency change (i.e. spin-down) are known from electromagnetic observations
- Correct the data for the known source parameters → signal becomes truly monochromatic
- Take a single Fourier Transform of all data
- Relatively quick for known pulsars, but cannot be applied in the general case





Known pulsar search results

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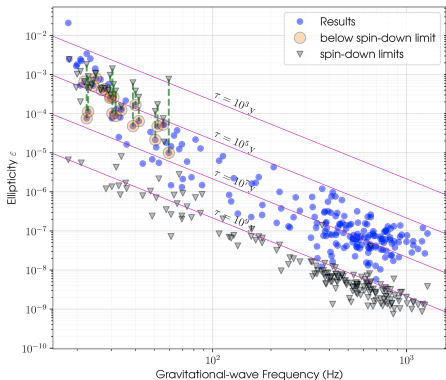
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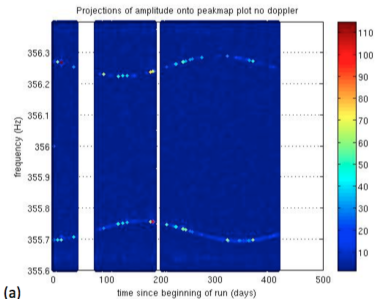
- The goal is to surpass the so-called spin-down limit: probe the case where $< 100\%$ of rotational energy is emitted as gravitational waves
- Results shown for 2nd observing run [1]



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Principles of data analysis for unknown pulsars

- Use as long Fourier Transforms lengths as possible!
- Confine signal power into a single frequency bin
- Analyze as much data as possible
- Primary frequency modulation: Doppler shift ($\mathcal{O}(10^{-4} f)$ Hz). Restricts the Fourier Transform to $\mathcal{O}(1000)$ s
- The noise is not Gaussian
- For each Fourier transform: $h(t) \rightarrow h(f)$. This creates a time/frequency spectrogram [7]





Types of searches

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Search type	Description	Sources
targeted	known $\alpha, \delta, f_0, \dot{f}$	Crab, Vela
directed	known α, δ	galactic center
all-sky	nothing known	any

α : right ascension

δ : declination

f_0 : pulsar rotation frequency at a reference time t_0

\dot{f} : spindown

- We “point” to specific locations when analyzing the data
- We also analyze systems in known and unknown binaries [5], and for post-merger remnants [3]



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Targeted search for pulsar J0537-6910

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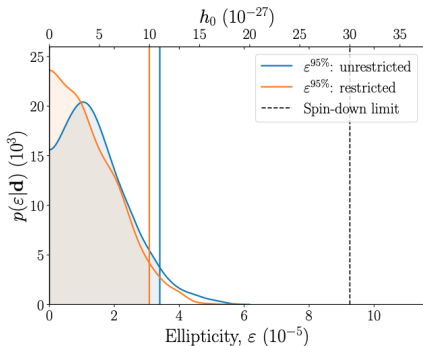
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- These are posteriors that show the probability that a particular signal with a certain parameters exists, given data [4]
- An amplitude of 0 cannot be excluded \rightarrow no signal
- Constraints from current observing run (O3)



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Existing constraints from an all-sky search

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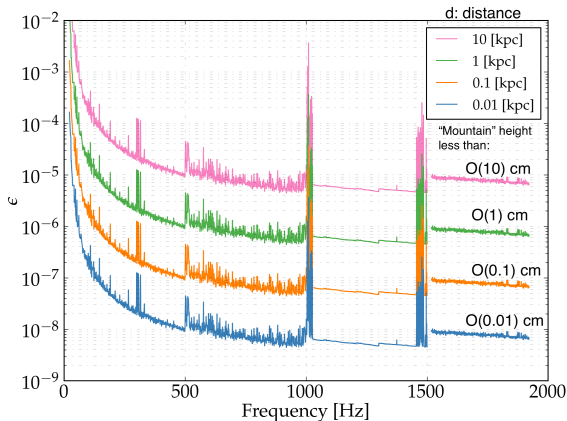
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- These are upper limits: the minimum deformation, or ellipticity ϵ , we could have seen at 95% confidence
- Deformations with smaller ϵ are easier to form
- LIGO/Virgo's 2nd observing run data [2]

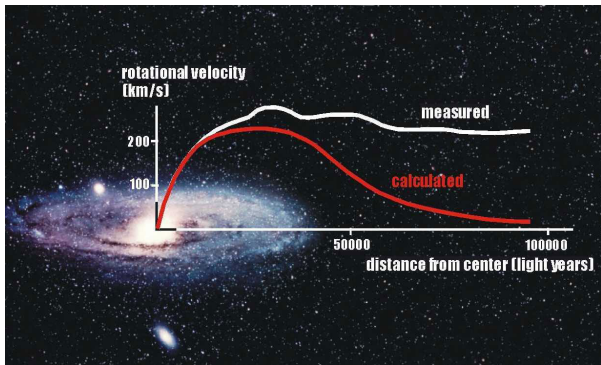


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The dark matter problem

- Experiments hint at the existence of dark matter, but its nature remains elusive
- Gravitational waves can probe ultralight dark matter, $\mathcal{O}(10^{-14} - 10^{-11} \text{ eV}/c^2)$, and ultra-heavy ($\mathcal{O}(M_{\odot})$)





The dark matter context

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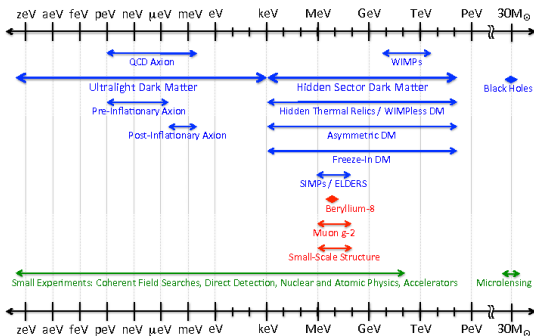
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- Various experiments are sensitive to different dark matter mass regimes
- The frequency sensitivity of ground-based detectors, 10-2000 Hz, fixes the boson masses we can look for

Dark Sector Candidates, Anomalies, and Search Techniques

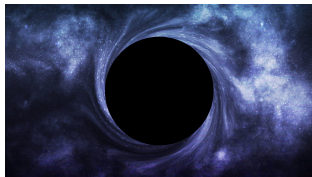


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Black hole superradiance

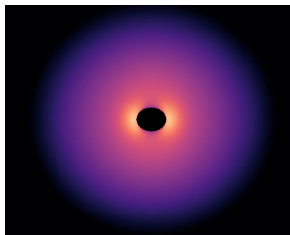
- Near a black hole, quantum fluctuations \rightarrow bosons pop into existence
- Many bosons fall in, but if the Compton wavelength is comparable to the radius of the black hole, bosons can scatter off it
- Greater effect for black holes with higher spins χ





Black hole superradiance

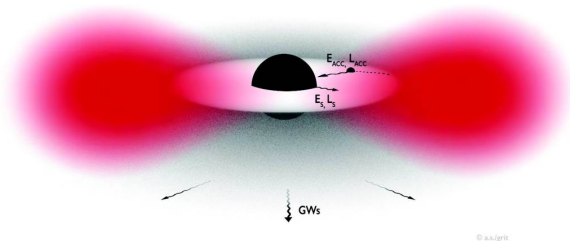
- Energy (mass/spin) extracted from the black hole by scattering bosons → outgoing boson amplitude boosted
- Unlike photons, bosons are massive, so they tend to be bound to black hole → successive scatterings [6]
- A boson “cloud” [13]





Growth of boson clouds

- Clouds are formulated as solutions to Schrodinger-like equations for a scalar field in the Kerr metric: “gravitational Hydrogen atom”
- The lowest, fastest growing state is $l = 1, m = 1$



- Superradiance (instability) condition: $\omega_{\text{boson}} < m\Omega_{\text{BH}}$
- No limit on the number of bosons in each state [9]





Depletion of scalar boson clouds

- Assume bosons couple to gravity and annihilate into gravitons [8]
- Gravitational waves are emitted from one energy level at a time \rightarrow monochromatic up to small spinup due to classical self-gravity
- Timescale of depletion \gg timescale of cloud growth
- Consider boson mass range $[10^{-14}, 10^{-11}]$ eV

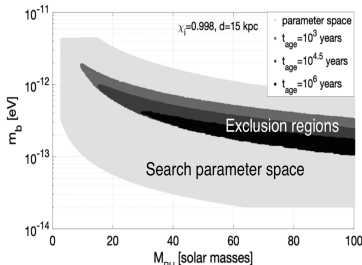
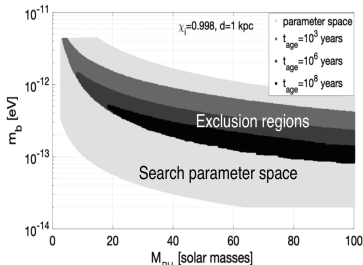
We expect *continuous* gravitational waves!





Constraints from an all-sky search

- Reinterpretation of upper limits on ϵ from slide 10: $\epsilon \propto h_0$, the gravitational-wave amplitude [17]
- $h_0, f \rightarrow$ constraints on boson mass m_b and black hole mass M_{BH} with assumptions on spins χ , distances d , and ages t_{age} of black holes



- More combinations of m_b and M_{BH} excluded for younger systems (small t_{age}) than older ones
- Darker colors are constraints on older systems



Directed search for Cygnus X-1

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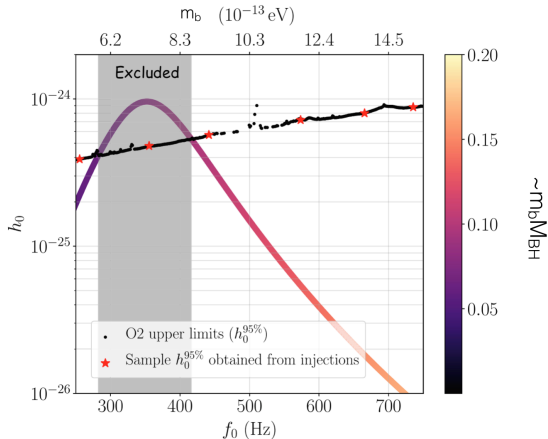
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- Binary parameters and mass/spin known
- Viterbi method used to find the optimal signal path [21]



Primordial black holes (PBHs)

- Black holes that formed in the early universe that can take on a wide range of masses depending on when they formed [10]
- Not predicted by standard model → detection would mean new physics
- PBHs could be linked to dark matter:
 - Dark matter may support cooling mechanisms that allow densely populated regions in space to collapse into PBHs [20]
 - Dark matter may interact with neutron stars and collect inside of them, causing a collapse [14]





How do PBHs form?

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- Early universe contained inhomogeneities that stopped too dense regions from expanding, causing collapses [11]
- Quantum fluctuations in various inflation theories [12]
- A binary could form if two PBHs form independently but are kept from merging by the gravitational pull of other nearby PBHs [19]
- Also, formation could occur through capture in a PBH halo [22]
- ... many more



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$$\dot{f}_{\text{gw}} \propto \mathcal{M}^{5/3} f_{\text{gw}}^{11/3} = k f_{\text{gw}}^{11/3} \quad (4.1)$$

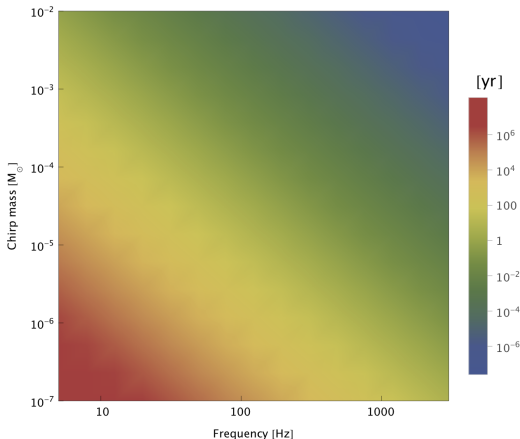
$$f_{\text{gw}}(t) = f_0 \left(1 + \frac{8}{3} k f_0^{n-1} (t - t_0) \right)^{-\frac{3}{8}} \quad (4.2)$$

- The inspiral portion of a merger can be modelled as a quasi-circular orbit that loses energy due to gravitational-wave emission
- The approximation holds at distances $> R_{\text{isco}}$
- For chirp masses $\mathcal{M} = [10^{-5}, 10^{-3}] M_{\odot}$, $f_{\text{isco}} > O(\text{MHz})$
- Timescale of emission varies from $O(\text{yrs})$ to $O(s)$
- Continuous-wave methods applicable
- Methods also applicable to binaries that could be seen by future detectors





How long until PBHs merge?



- Within this frequency range, lighter equal-mass PBHs take longer to merge [16]
- Larger equal-mass PBHs would have formed later





$$\dot{f}_{\text{gw}} \propto \mathcal{M}^{5/3} f_{\text{gw}}^{11/3} = k f_{\text{gw}}^{11/3} \quad (4.3)$$

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Simulated signal

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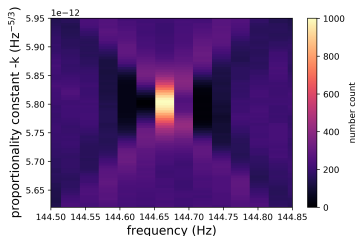
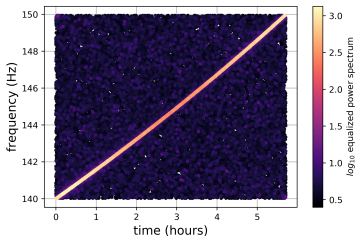
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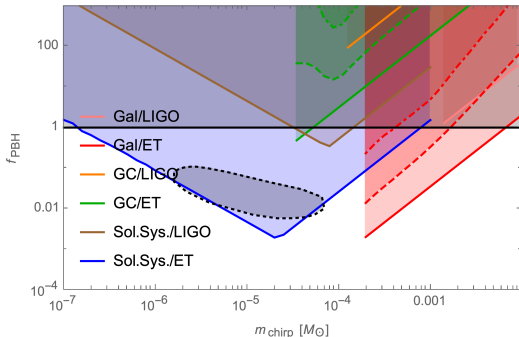
- Hough Transform maps points in the time/frequency plane to lines in a plane that relates to the frequency and spin-up of the source
- Injection parameters: $h_0 = 4 \times 10^{-23}$, $f_0 = 144.669$ Hz, $\mathcal{M} = 10^{-3} M_{\odot}$, $k = 5.73 \times 10^{-12} \text{ Hz}^{-5/3}$



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Projected constraints in the current and advanced detector era



- Projected constraints on the fraction of dark matter PBHs could compose for current and future detectors at different distances [16]
- Agnostic and thermal mass functions used
- Oval-dotted lines are constraints from gravitational lensing experiments (OGLE)





Conclusions

- Many interesting physical systems are sources of continuous gravitational waves
- Similar analysis techniques can be employed for a variety of sources
- Current and future searches for boson clouds could be performed over the whole sky, or targeting the remnants of binary mergers



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Backup slides



General search scheme

- Choose a Fourier Transform length
- Fourier Transform the strain data, and perform a matched filter (in the case of fully coherent searches), or make a time/frequency spectrogram (in semi-coherent searches)
- Apply a method robust to noise disturbances on the spectrogram to estimate parameters





- Look for coincident candidates in each detector
- At this point, many candidates exist and must be discarded in a “follow-up” stage
- Demodulate the signal with the known parameters, and increase the Fourier Transform length
- In the event of no detection, put upper limits with simulations or estimations of theoretical sensitivity





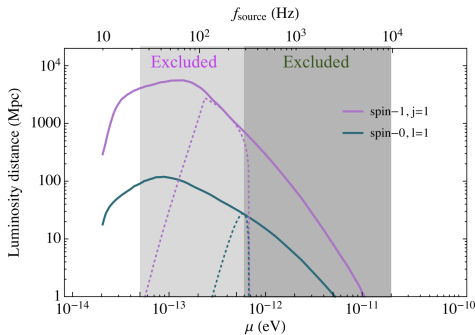
Vector bosons

- Emit gravitational waves with higher amplitudes, but on shorter timescales, than those from scalar bosons
- For shorter signals, spinup becomes important
- Parameter space mostly composed of “transient” continuous wave signals
- Possible targets: merger remnants, x-ray binaries
- Interplay between instability and depletion timescales important





Distance reach



- Assumes monochromatic signal
- Black hole mass chosen as a function of particle mass to give the strongest gravitational-wave signal
- Dotted line: $M_{BH} = 64M_{sun}$





Ellipticity:

$$\epsilon \equiv \frac{|I_{xx} - I_{yy}|}{I_{zz}}, \quad (5.1)$$

Amplitude of continuous wave:

$$h_0 = \frac{16\pi^2 G}{c^4} \frac{I_{zz} \epsilon f_{\text{rot}}^2}{d}, \quad (5.2)$$

Spindown limit:

$$h_{0,\text{sd}} = \frac{1}{d} \left(\frac{5GI_{zz}}{2c^3} \frac{|\dot{f}_{\text{rot}}|}{f_{\text{rot}}} \right)^{1/2}, \quad (5.3)$$





References I

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