

The Origin of Cosmic Fireworks

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Martin Obergaulinger, Re'em Sari

14/7/09

A Tale of Three Explosions

- Gamma Ray Bursts of all sorts
- Radio Flares from Neutron star mergers
- Tidal Disruption Events

I. The origin of GRBs?

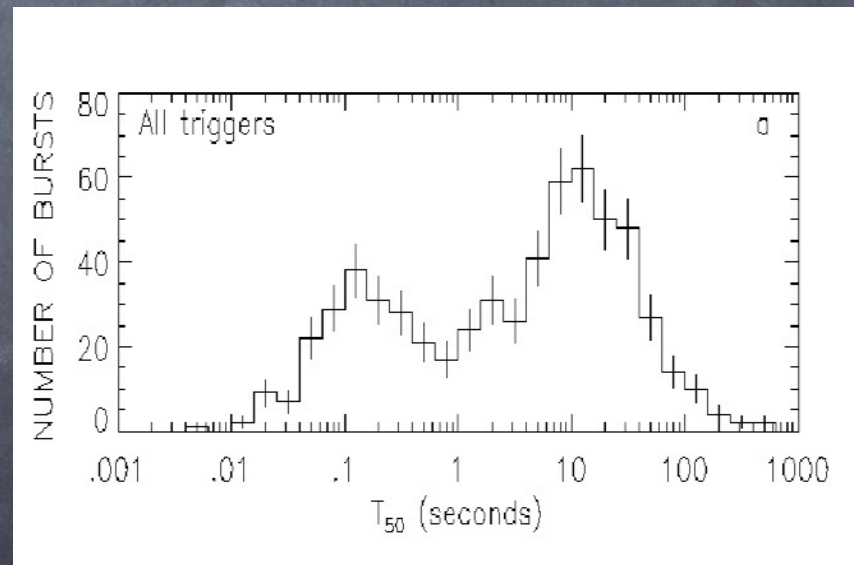
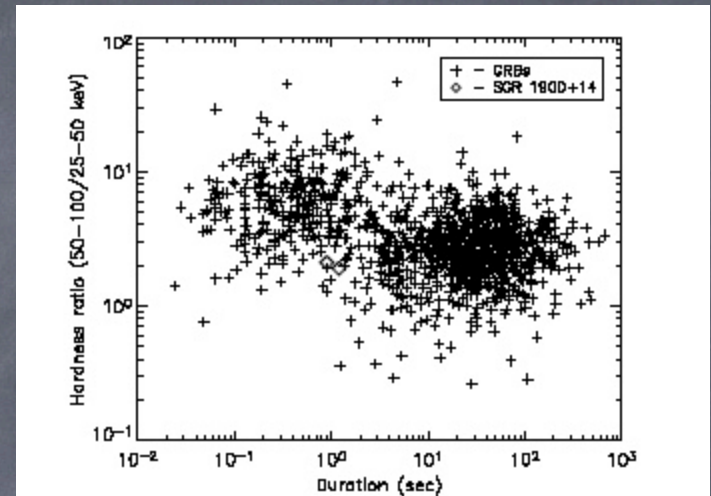
- Bromberg, Nakar & TP, ApJL 2011
Bromberg, Nakar, TP & Sari, ApJ 2011
- Bromberg, Nakar, TP & Sari, arXiv 1111.0969
- Bromberg, Nakar & TP to be submitted 2011
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Properties

- ◆ Duration 0.01–1000s

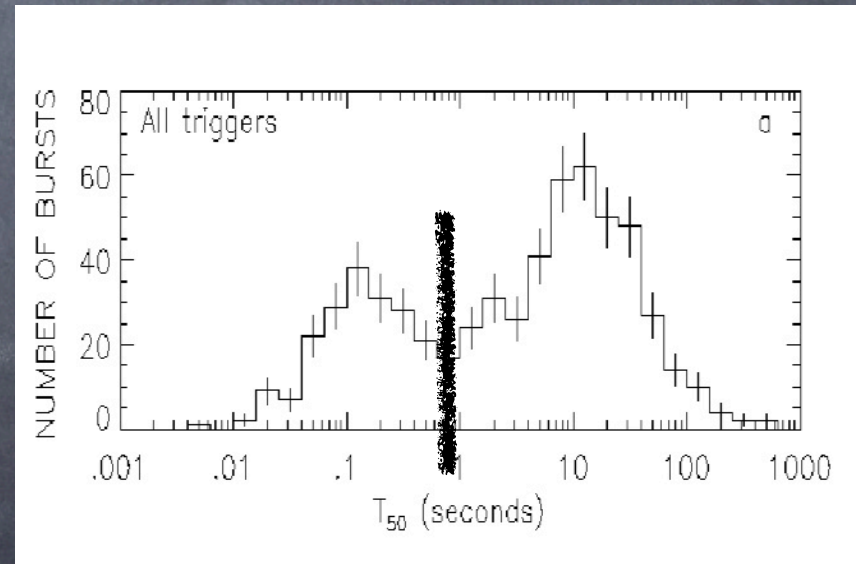
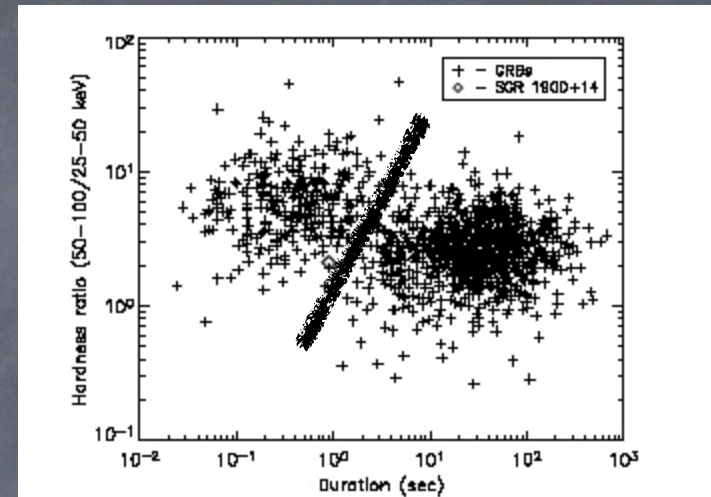
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(non thermal spectrum)
(very high energy tail,
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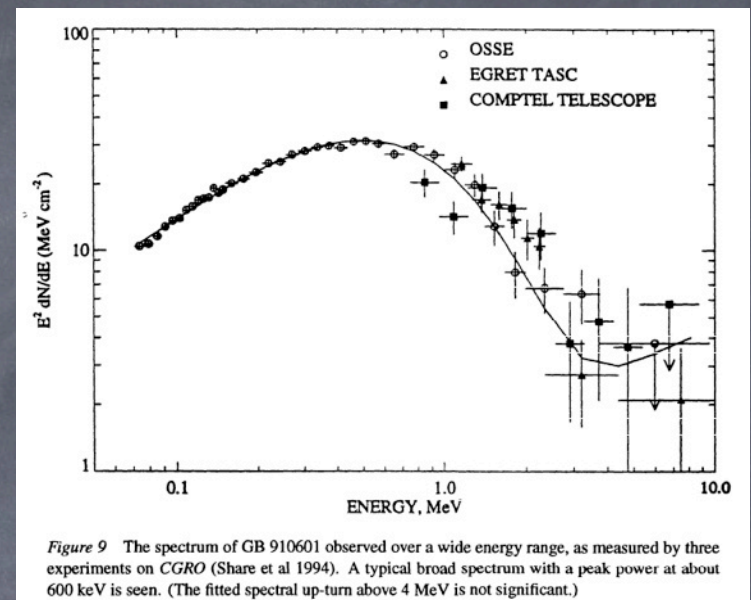
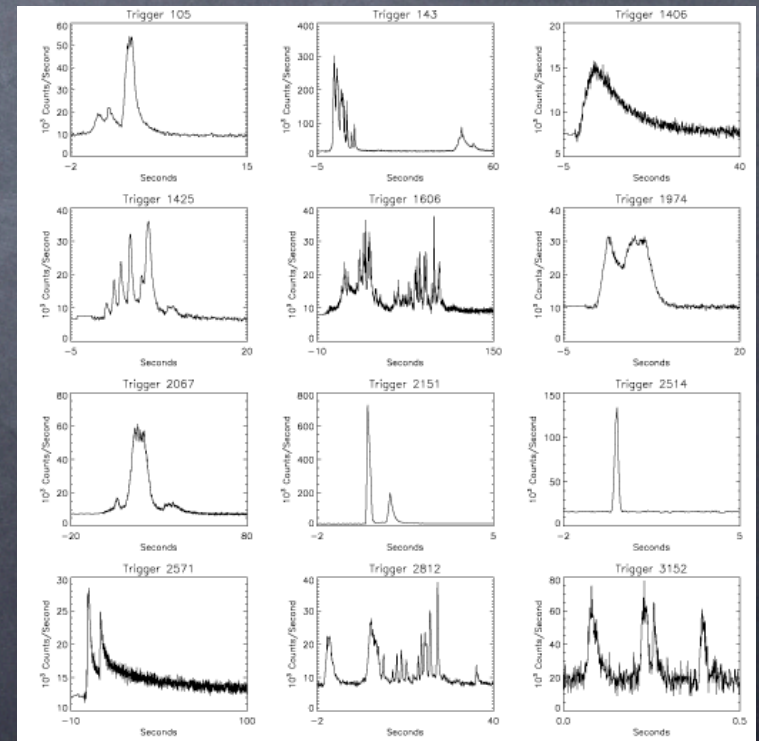


Figure 9 The spectrum of GB 910601 observed over a wide energy range, as measured by three experiments on *CGRO* (Share et al 1994). A typical broad spectrum with a peak power at about 600 keV is seen. (The fitted spectral up-turn above 4 MeV is not significant.)

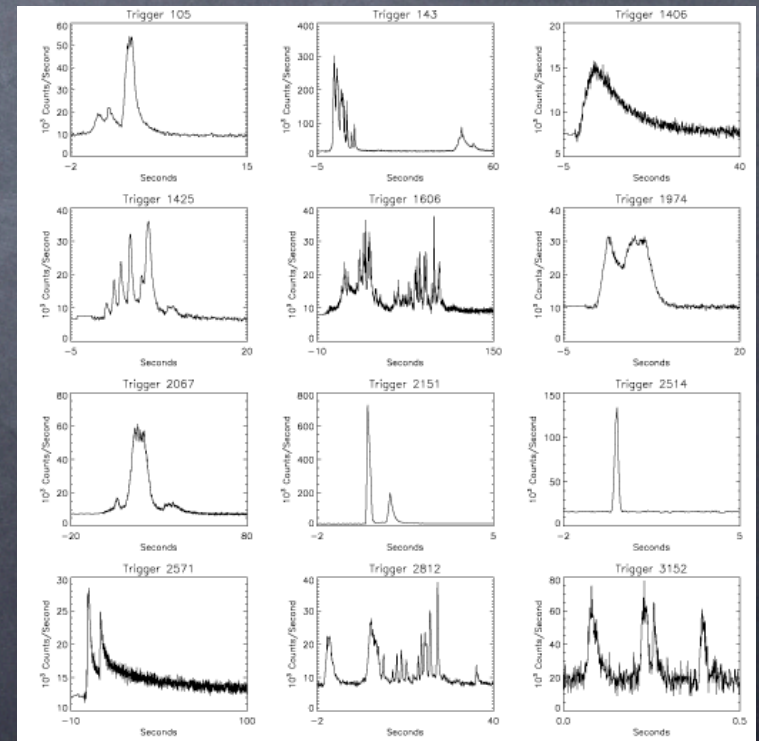
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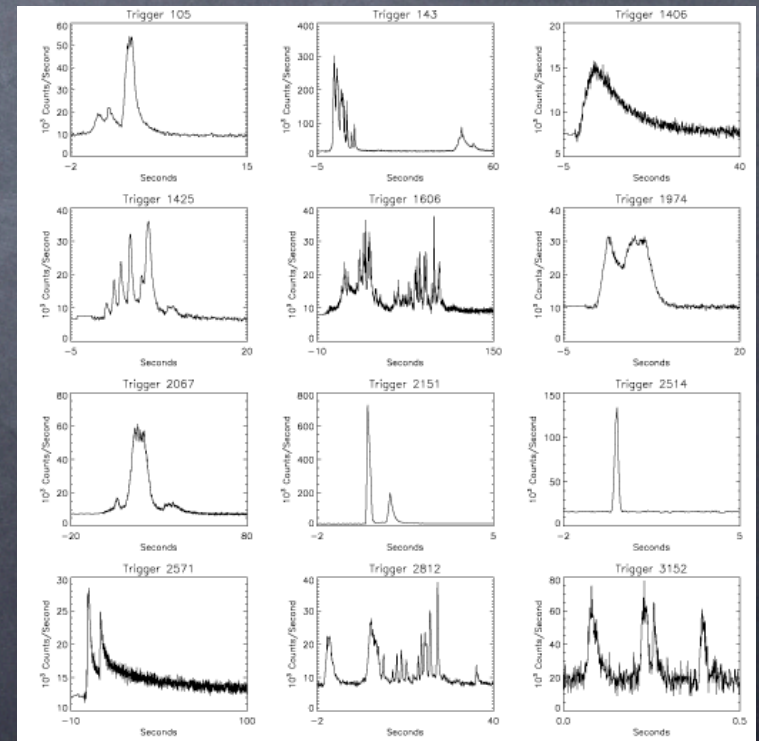
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- Followed by multiwavelength
Afterglow

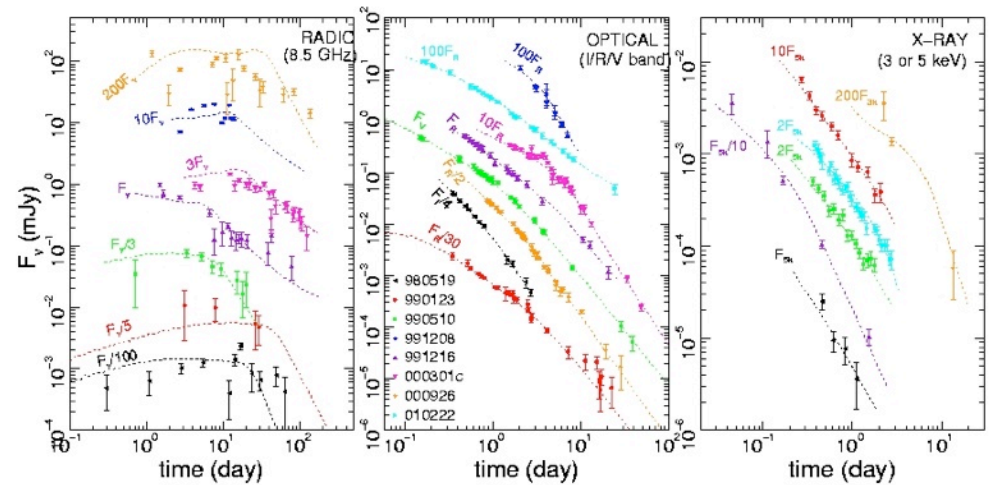
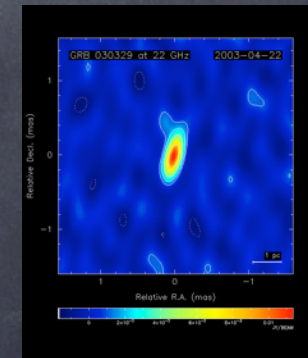
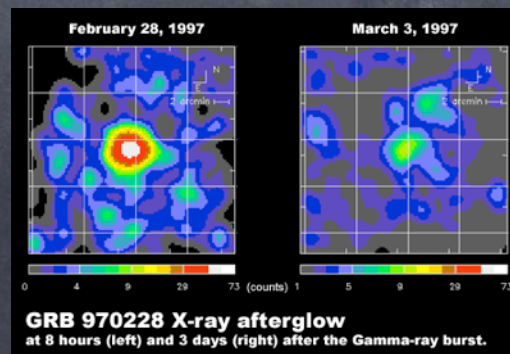
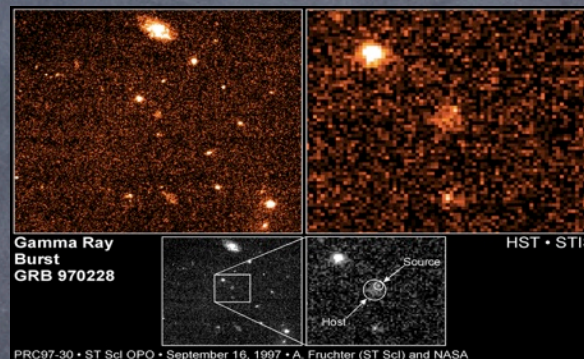


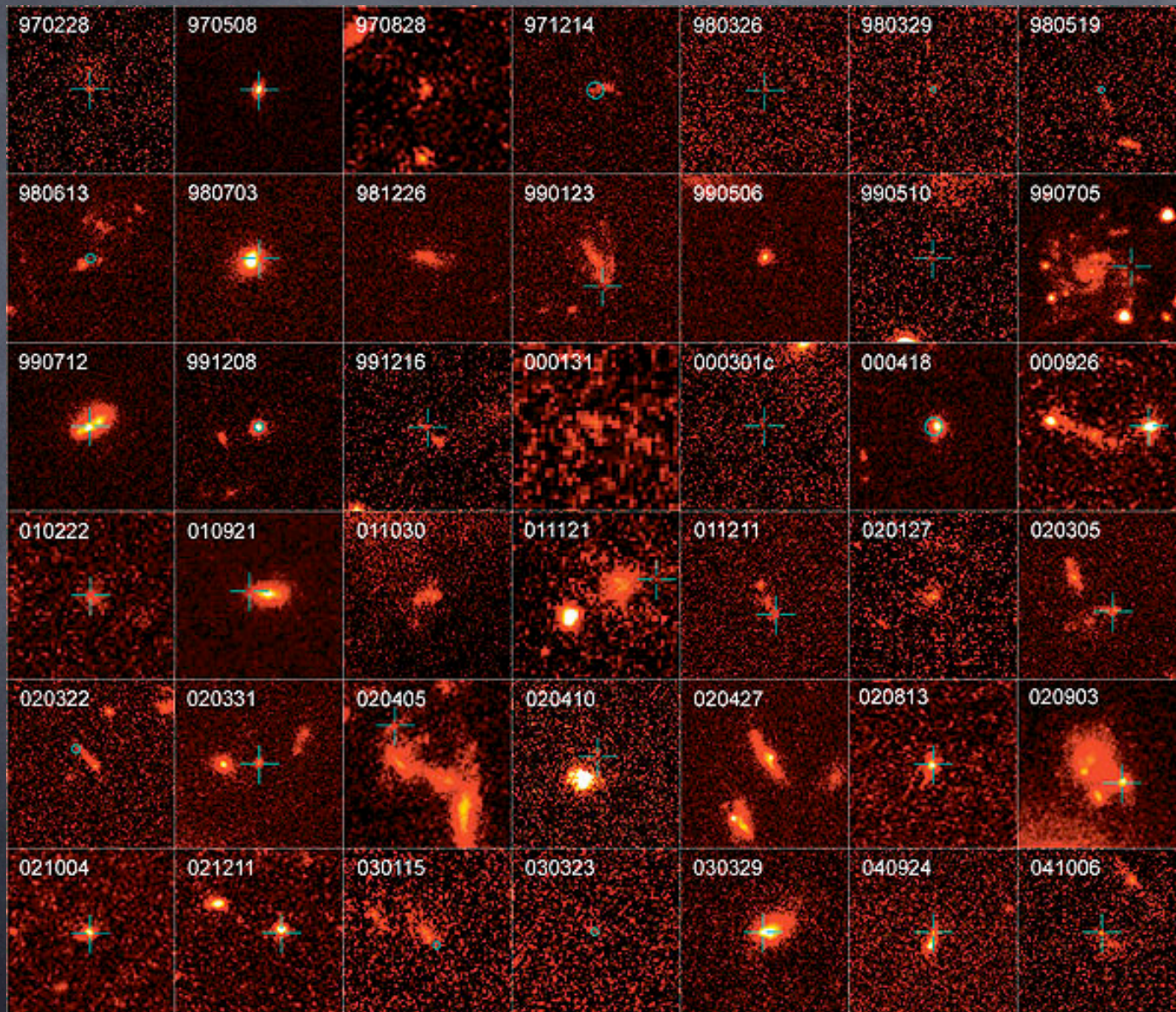
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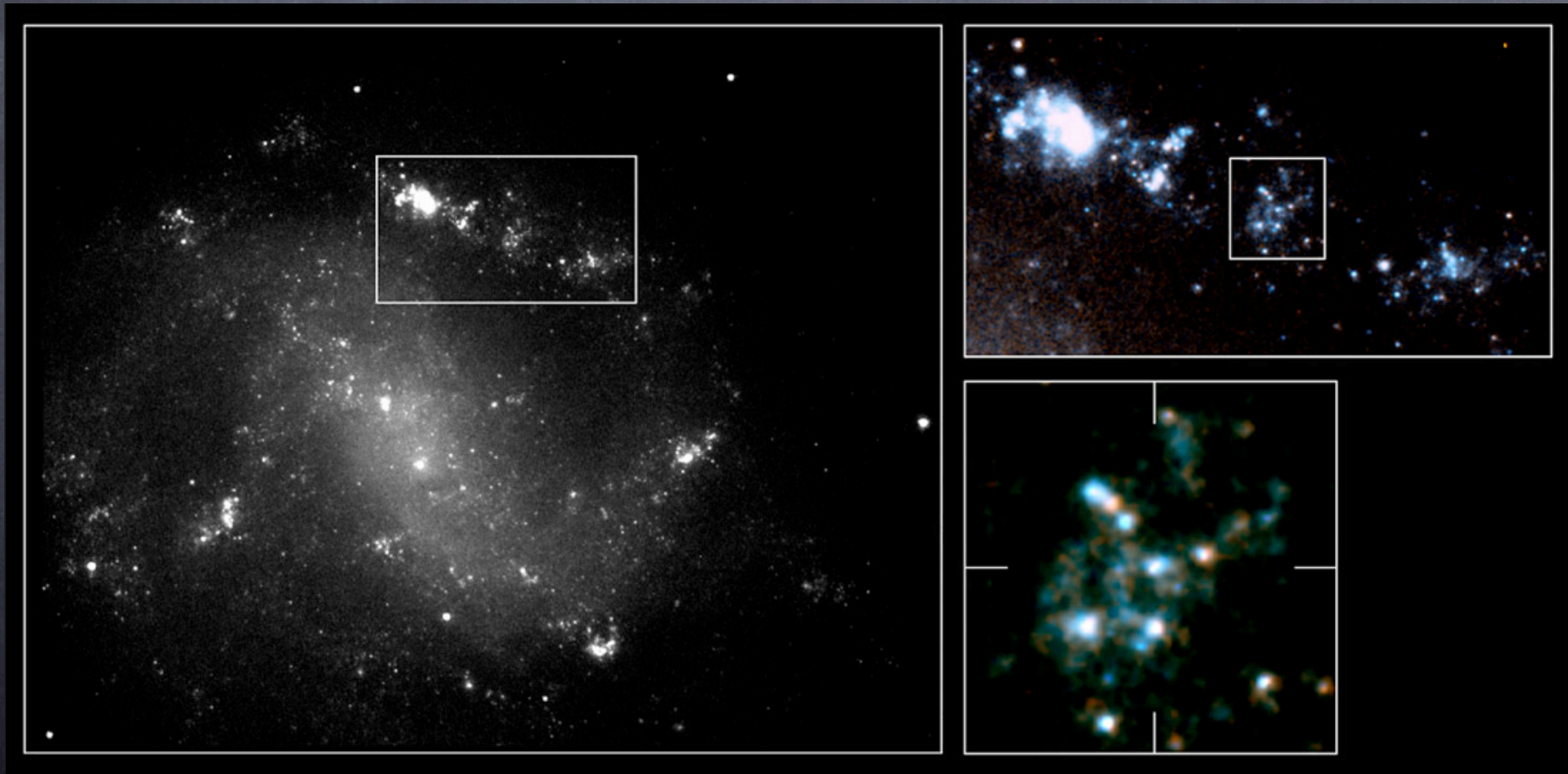
Energy, Energy, Energy + Time

- $E \approx 10^{51}\text{--}10^{52}$ ergs \approx the binding energy of a compact stellar mass object.
- $0.01\text{--}100$ sec + $E \approx 10^{51}\text{--}10^{52}$ ergs
 \Rightarrow a newborn stellar mass compact object.
- \Rightarrow Collapsing stars or mergers.

The (long) GRB-Supernova connection

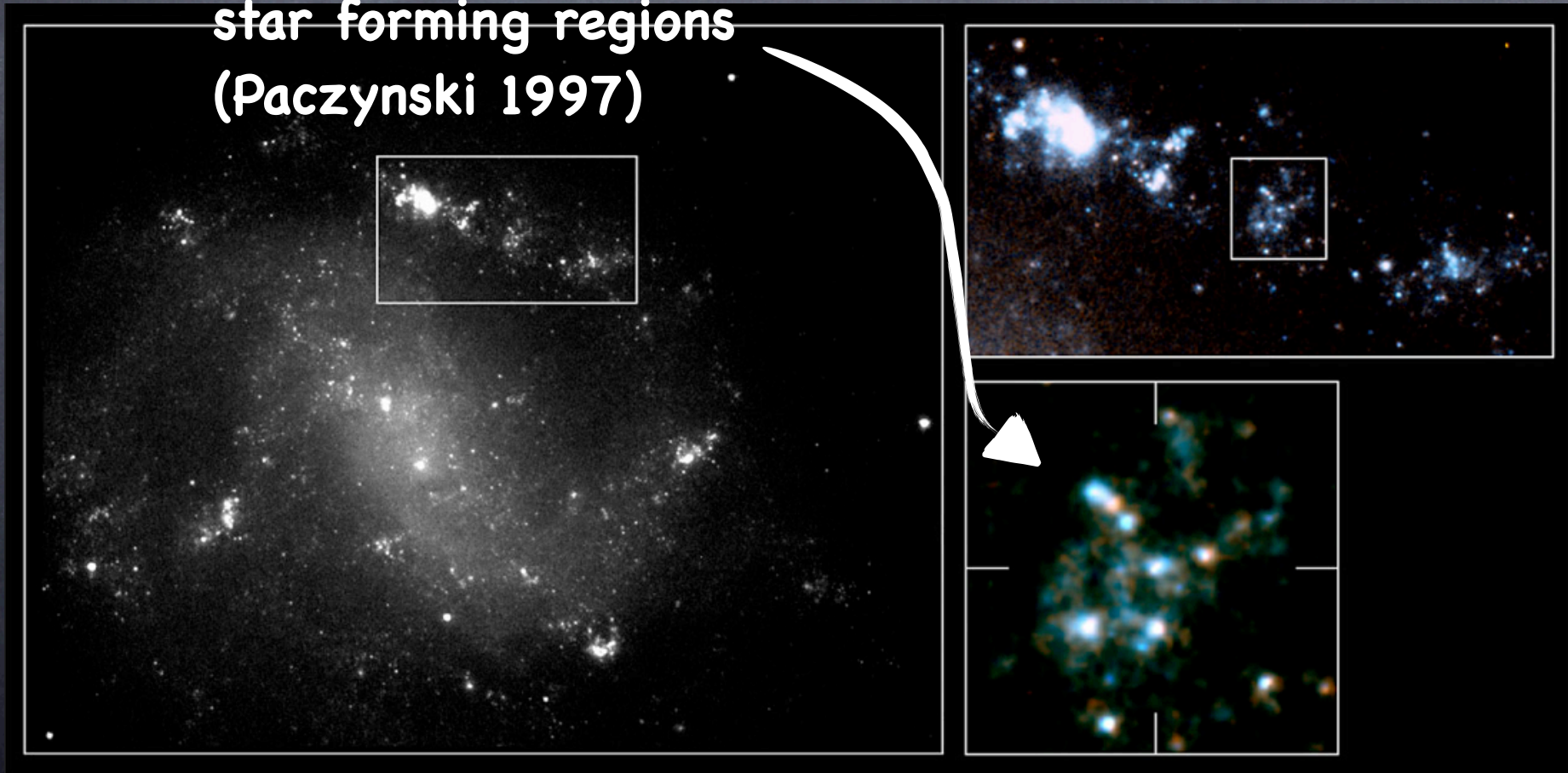


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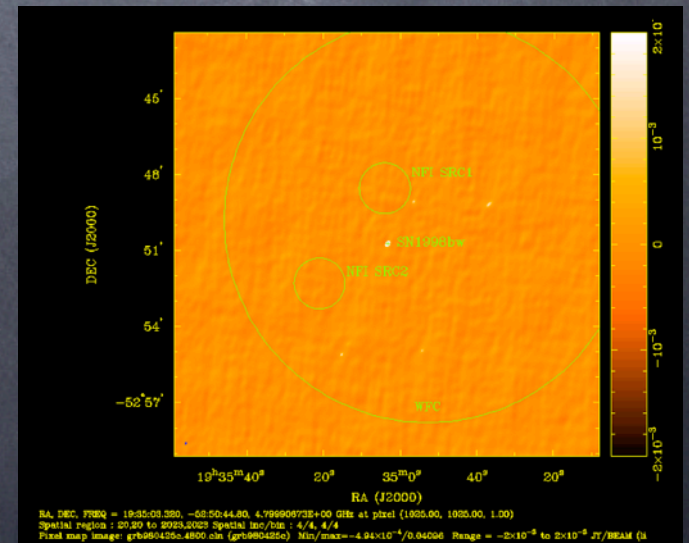
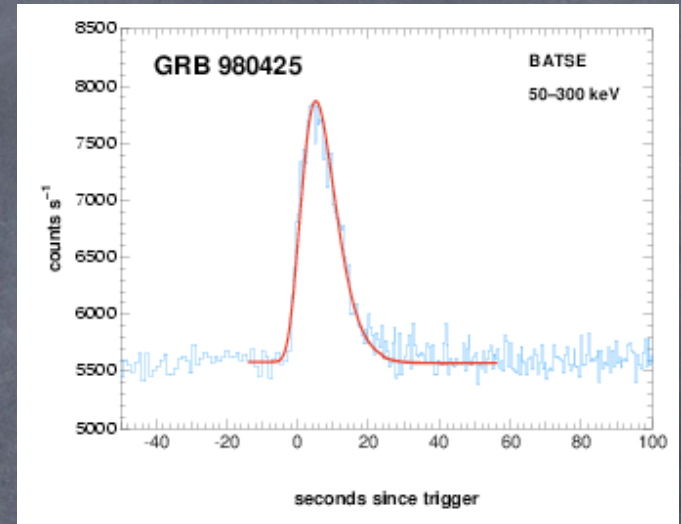
The (long) GRB-Supernova connection

- Observational indications
 - Long GRBs arise in star forming regions (Paczynski 1997)



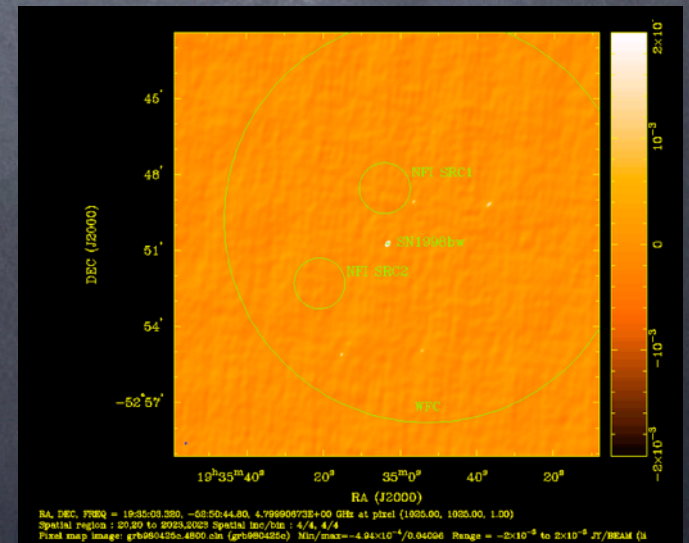
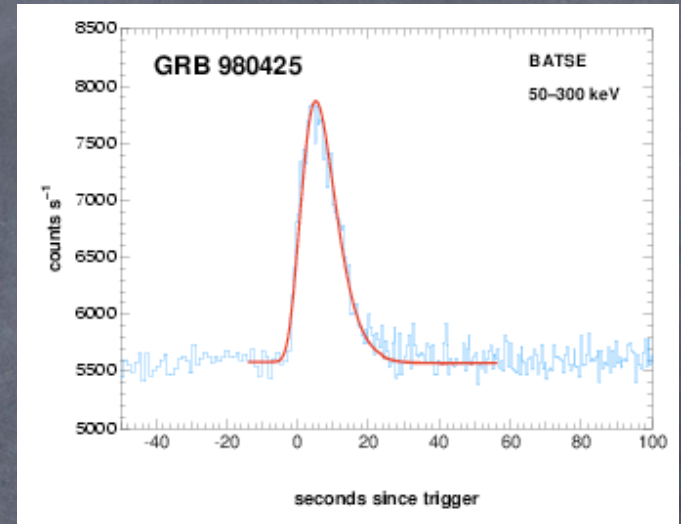
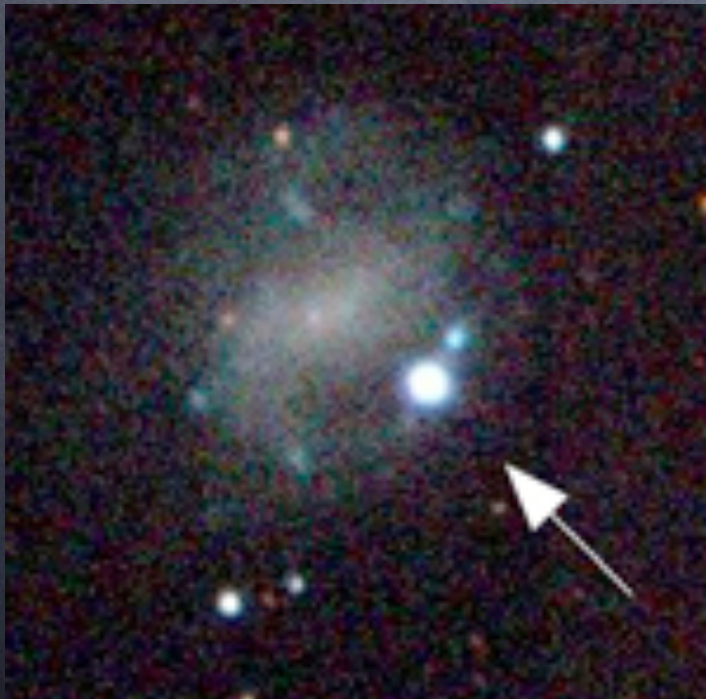
Supernova 1998bw-GRB980425

$$E_p \approx 67 \pm 40 \text{ keV}$$
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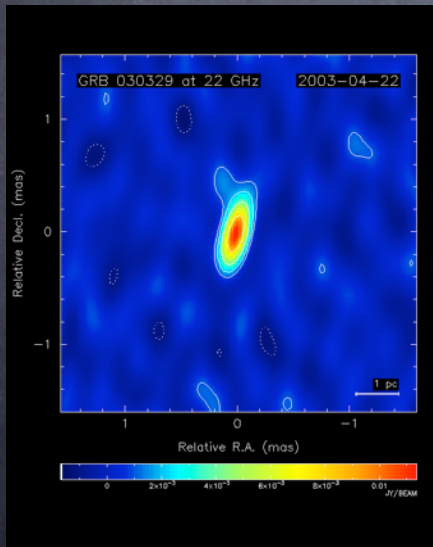
The Smoking Gun

GRB030329-SN 2003dh - a regular GRB with a 98bw like supernova.



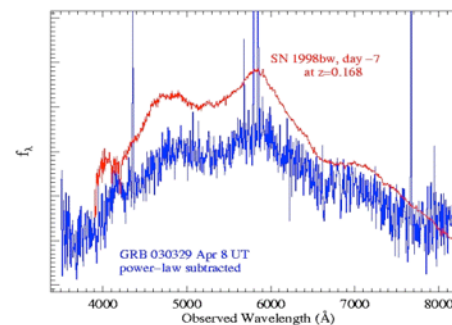
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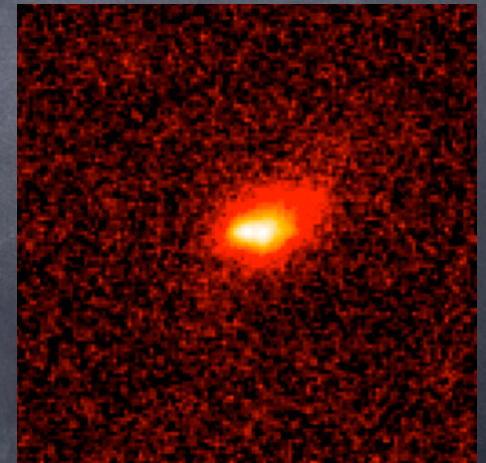


Supernova Spectrum Emergence

GRB 030329 is now also SN2003dh

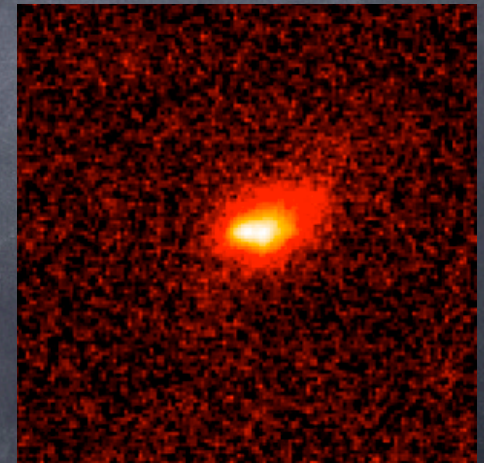
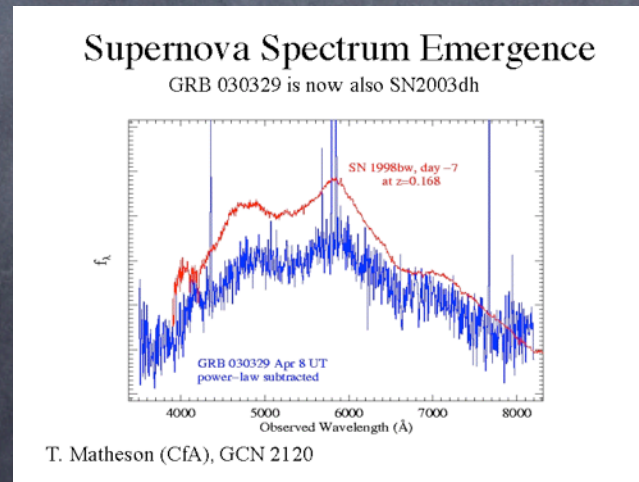
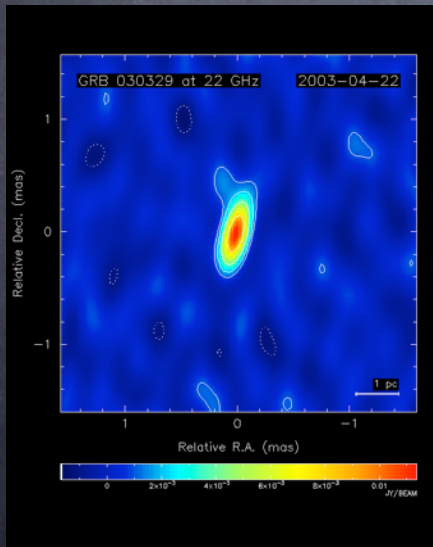


T. Matheson (CfA), GCN 2120

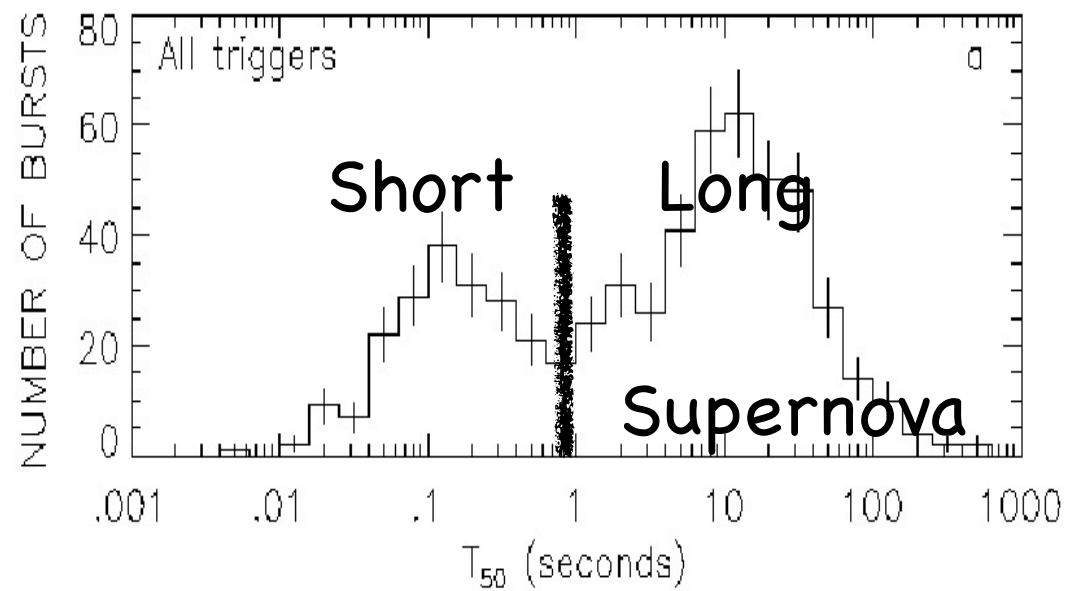


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● Recently we have also GRB101219B - SN 2010ma



Prologue

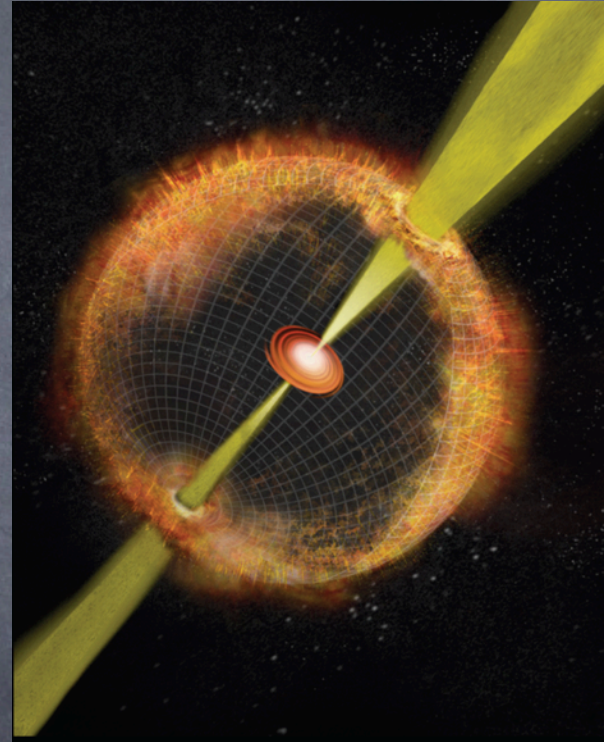
Several times during the short history of GRBs just when we thought we understood something Nature showed us to be wrong.

This may be one of these cases...

Or maybe not?

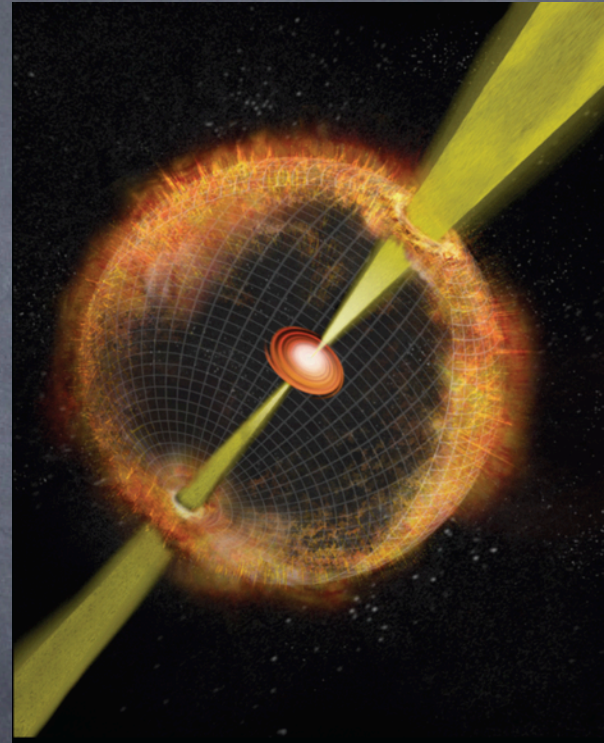
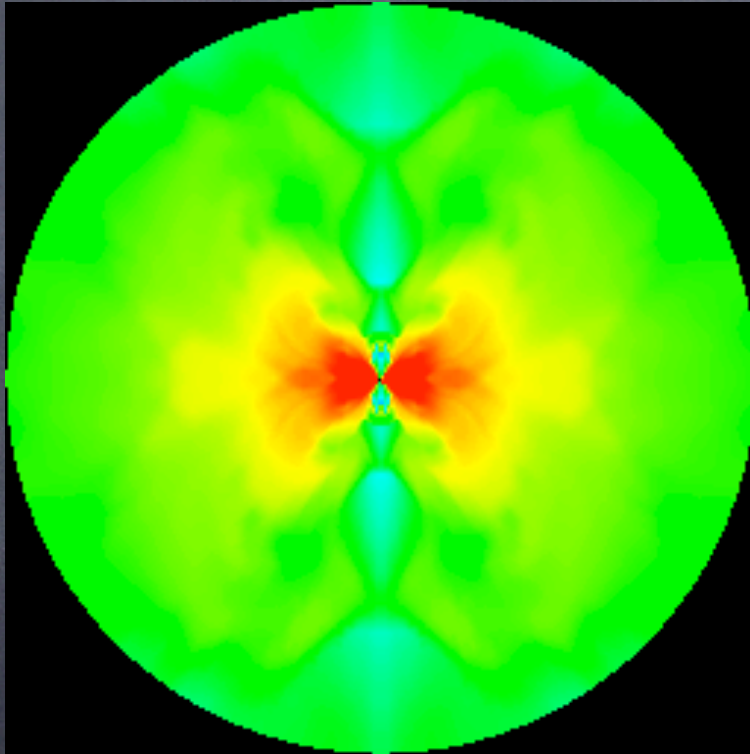
The Collapsar Model

(MacFadyen & Woosley 1998)

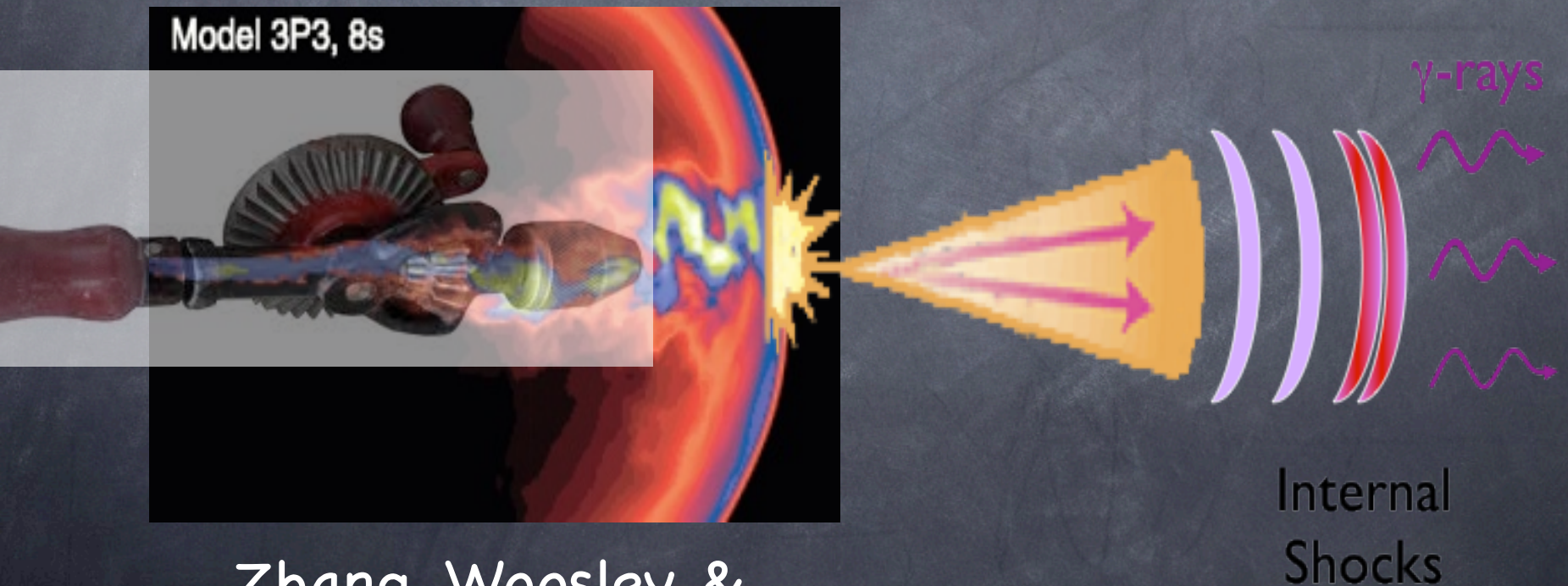


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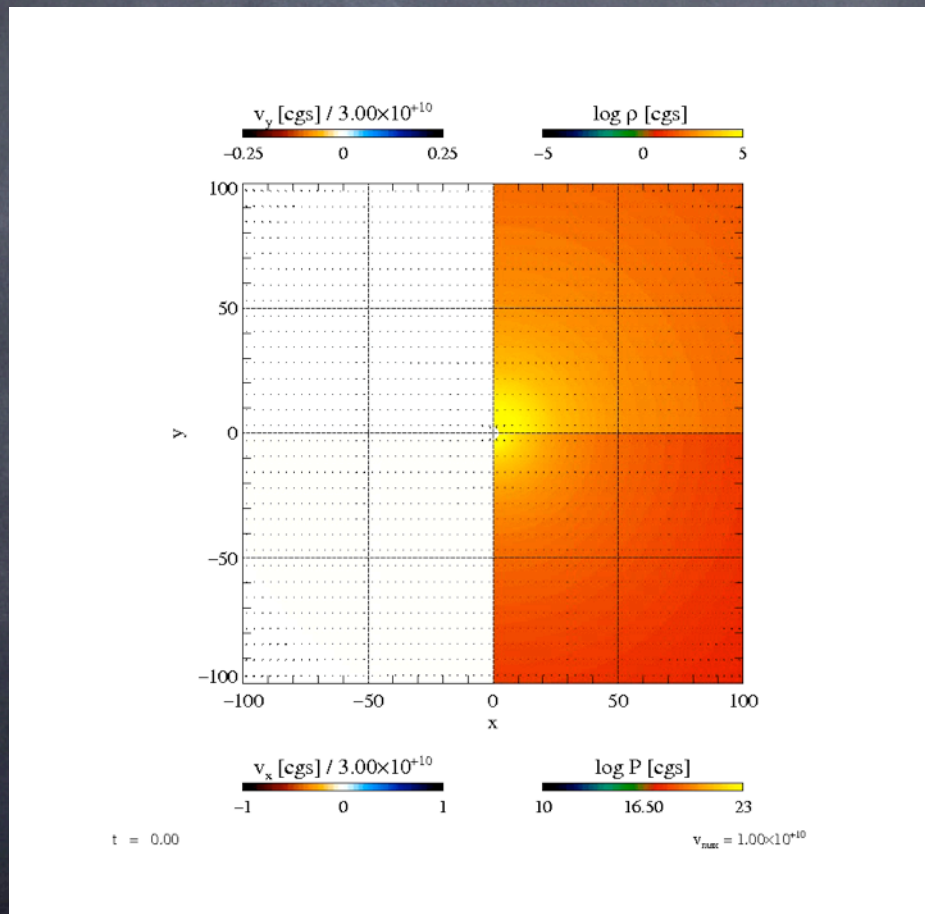
The Jet drills a hole in the star Model



Zhang, Woosley &
MacFadyen 2004

Jet Simulations

(Obergaulinger, Piran + 11)



Opening angle of 15° degrees at 2000 km into a star of 15 solar masses and solar metallicity. Constant energy injection rate, $5 * 10^{50}$ erg /s, through the entire run of the model. Lorentz factor at injection 7

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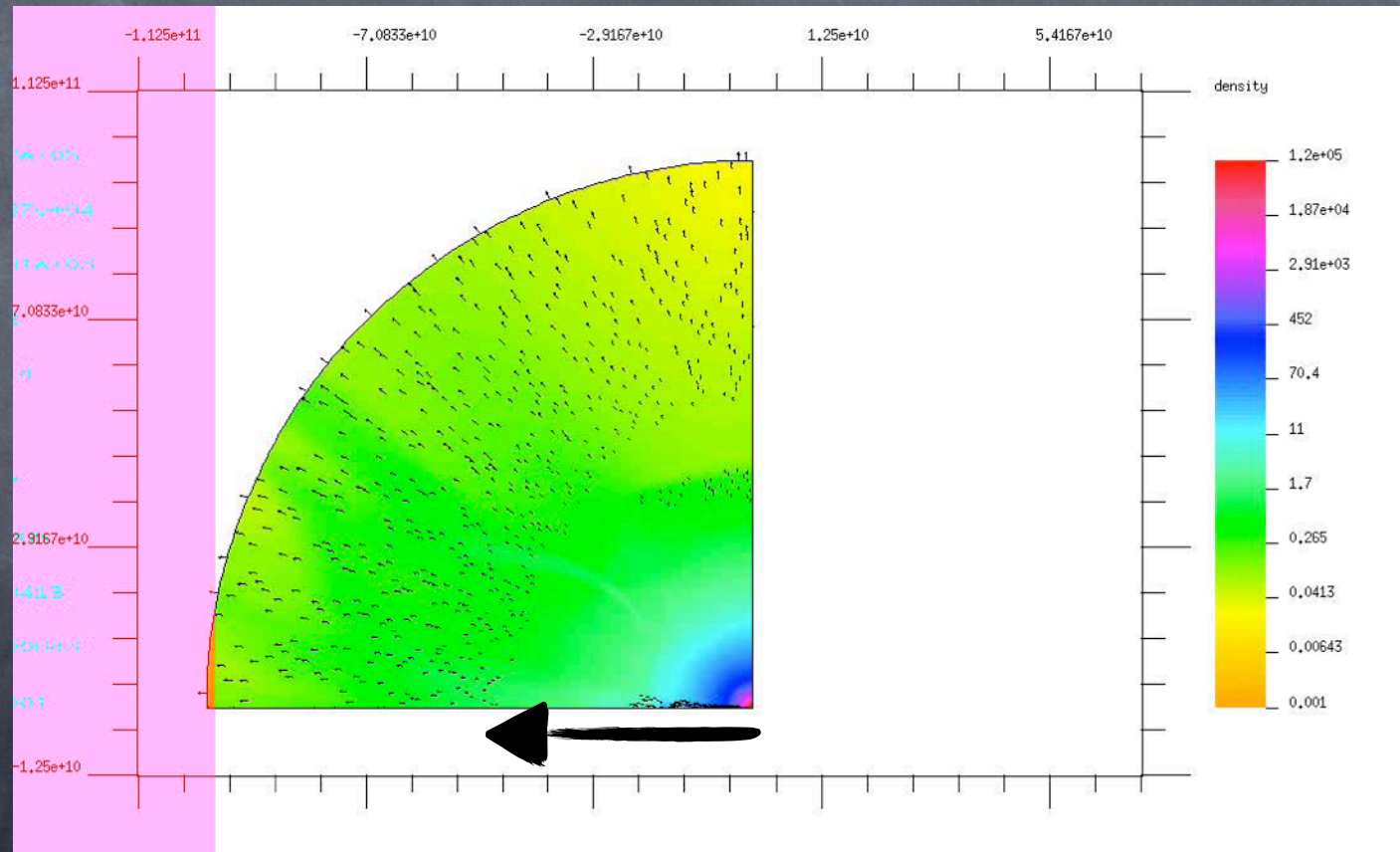
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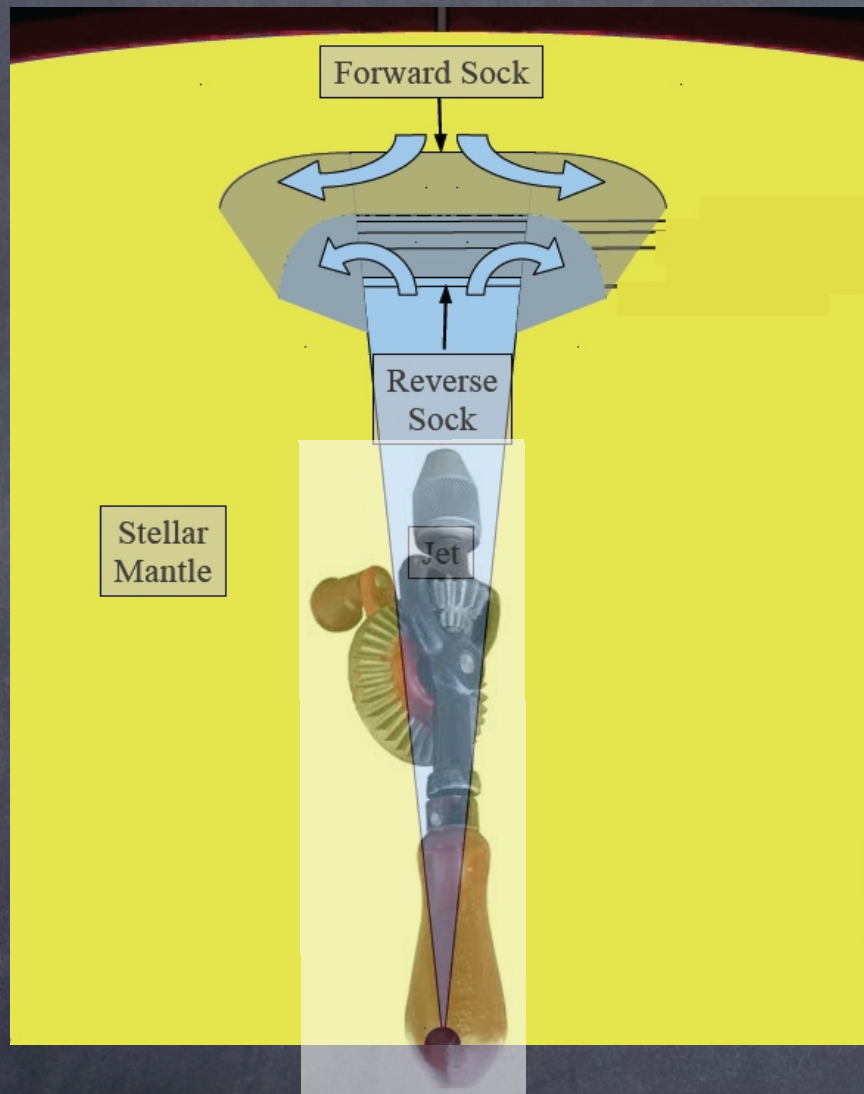
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Another explosion – Disruption of the Stellar envelope by the jet –

Genet, Livne, Obergaulinger & TP 2011

About one
solar mass
is ejected
non
spherically





Bromberg Nakar, TP,
Sari 11 ApJ 2011

- The jet dissipates its energy while propagating.
- The jet is slowed down to about $0.1c$

Jet breakout time

(Bromberg Nakar, TP, Sari 11 ApJ 2011)

$$t_b \simeq 15 \text{ sec} \cdot \left(\frac{L_{iso}}{10^{51} \text{ erg/sec}} \right)^{-1/3} \left(\frac{\theta}{10^\circ} \right)^{2/3} \left(\frac{R_*}{5R_\odot} \right)^{2/3} \left(\frac{M_*}{15M_\odot} \right)^{1/3}$$

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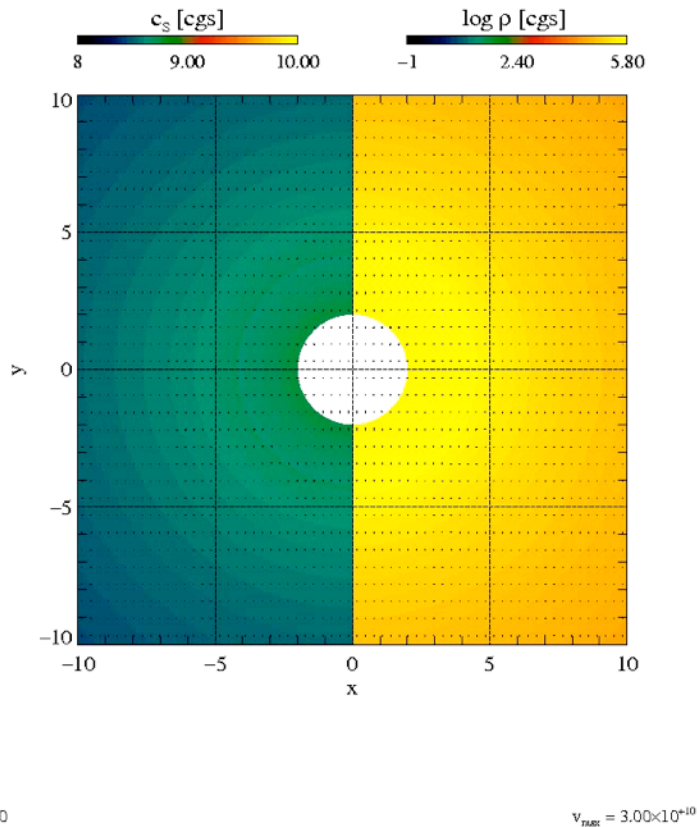
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The engine must be active until
the jet's head breaks out!*

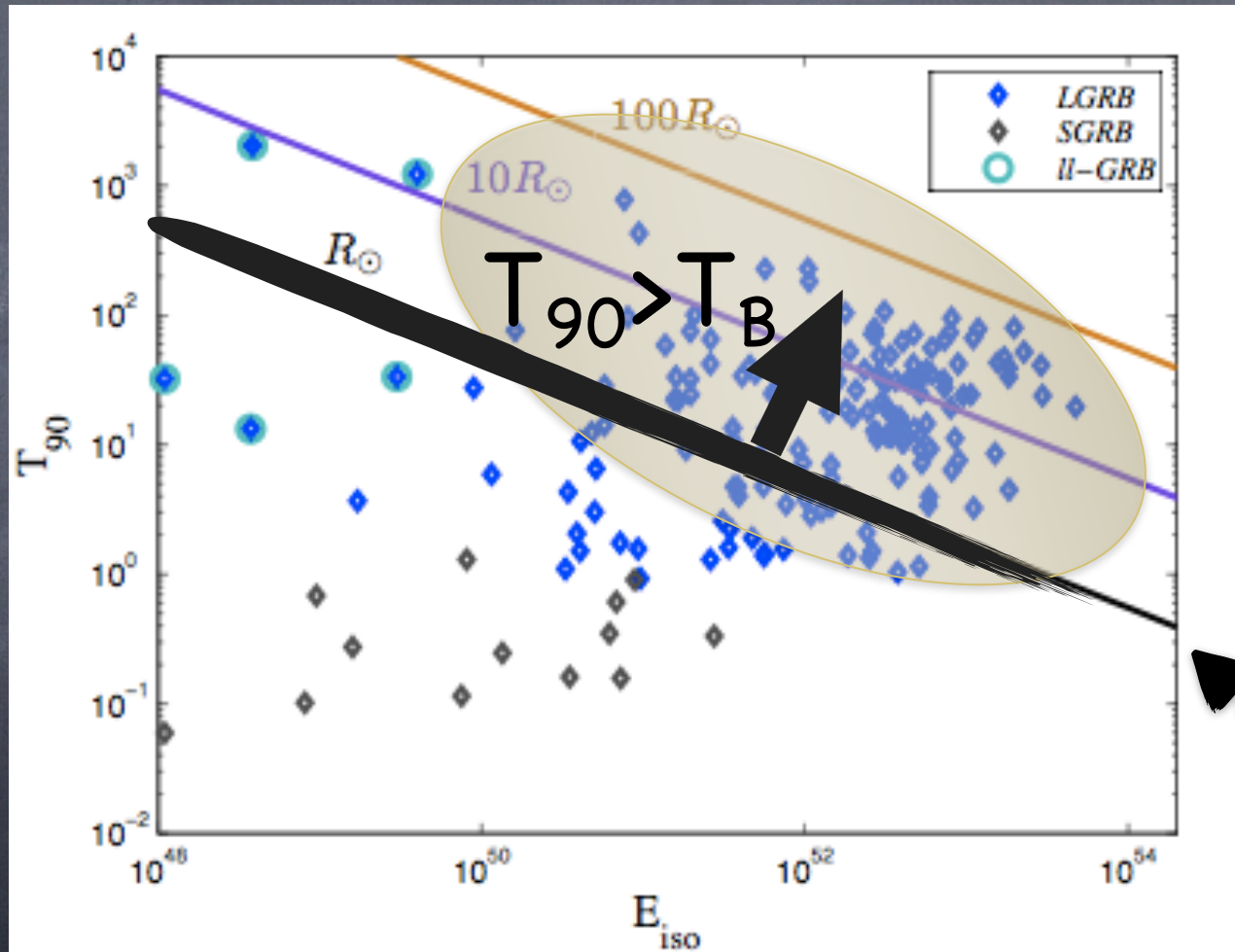
Jet Simulations – A Failed Jet

Jet (Obergaullinger, Piran + 11)

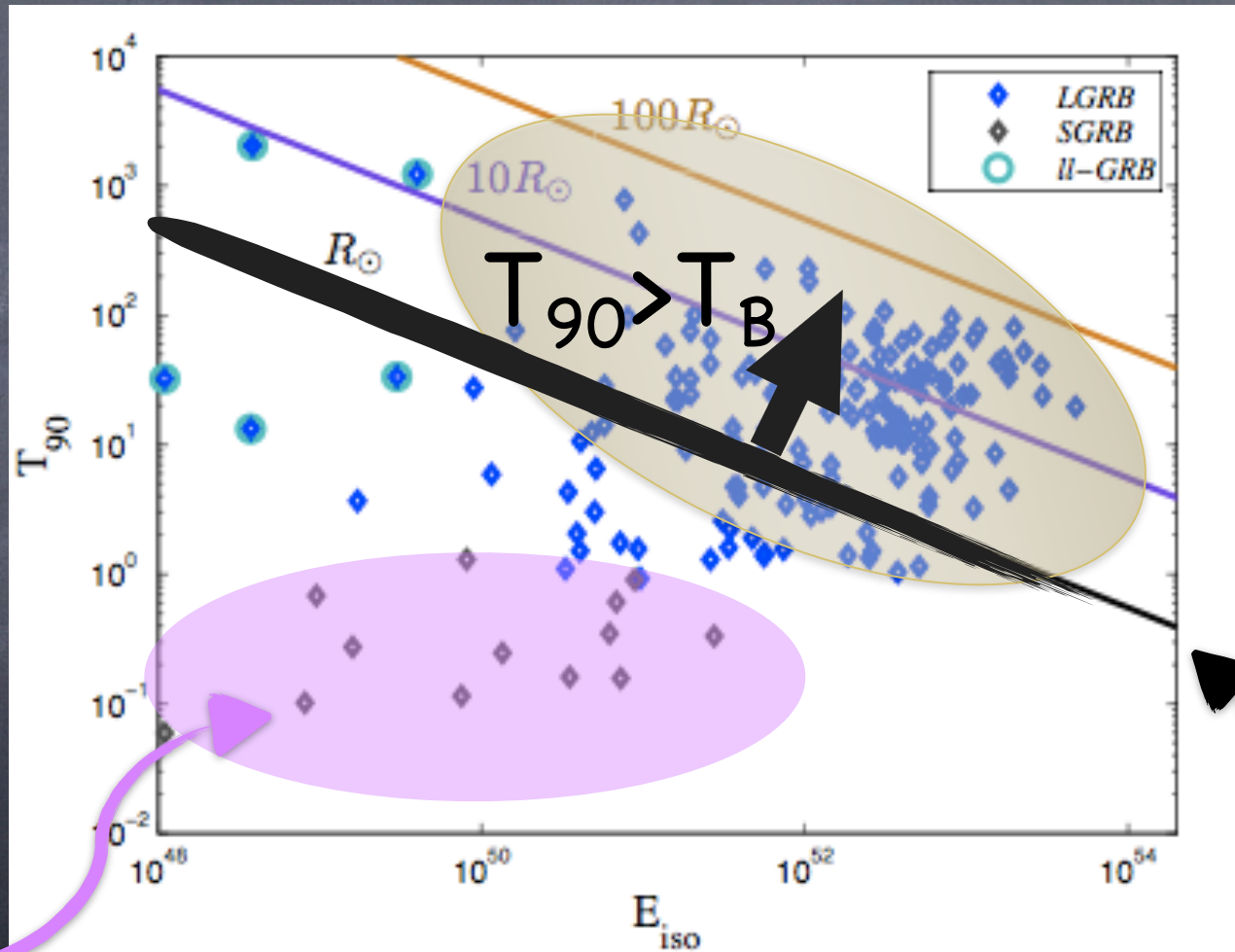


Opening angle of 15° degrees at 2000 km into a star of 15 solar masses and solar metallicity. Constant energy injection rate, 5×10^{50} erg/s, for 2 seconds.

Duration (T_{90}) vs Energy



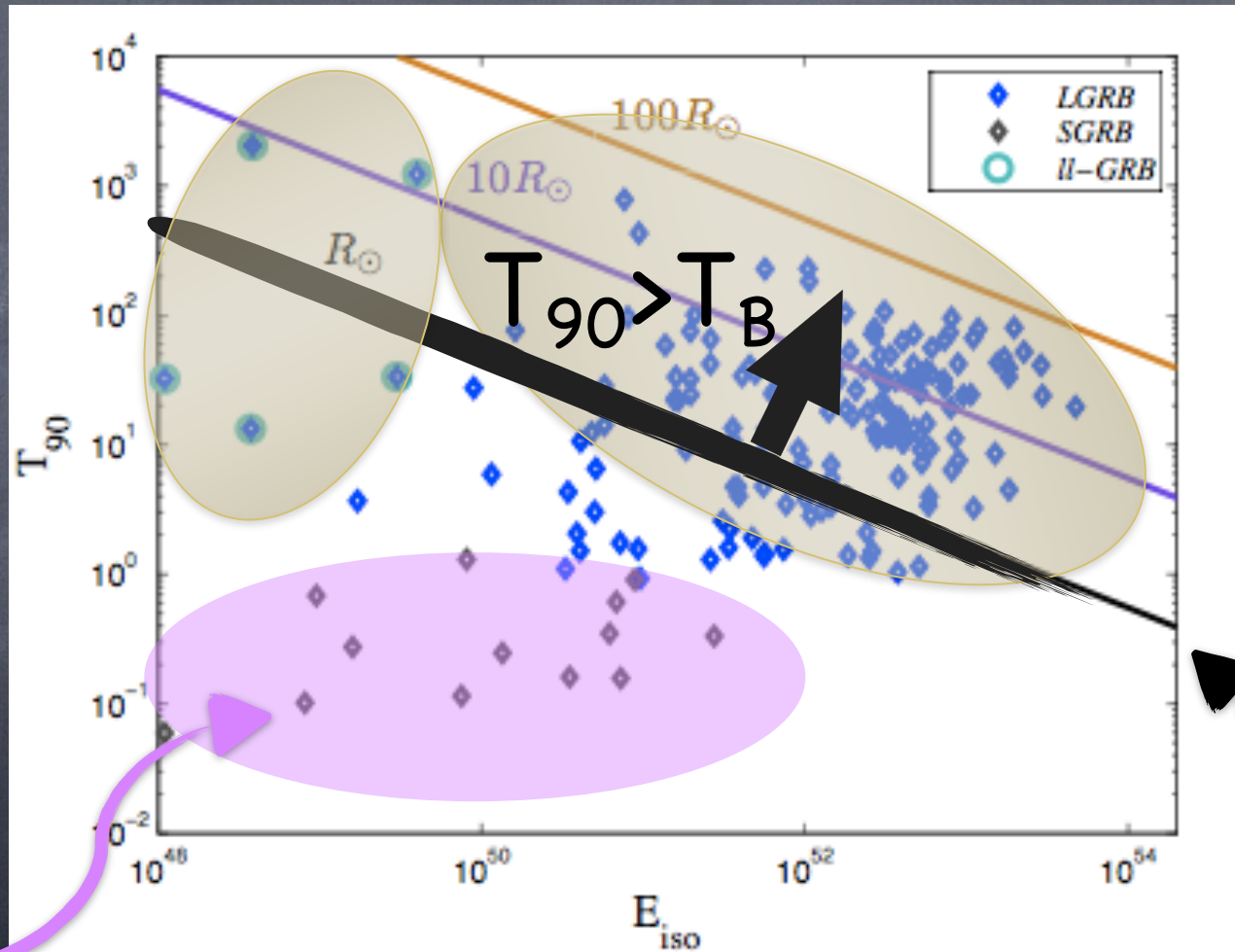
Duration (T_{90}) vs Energy



Short
GRBs

T_B

Duration (T_{90}) vs Energy



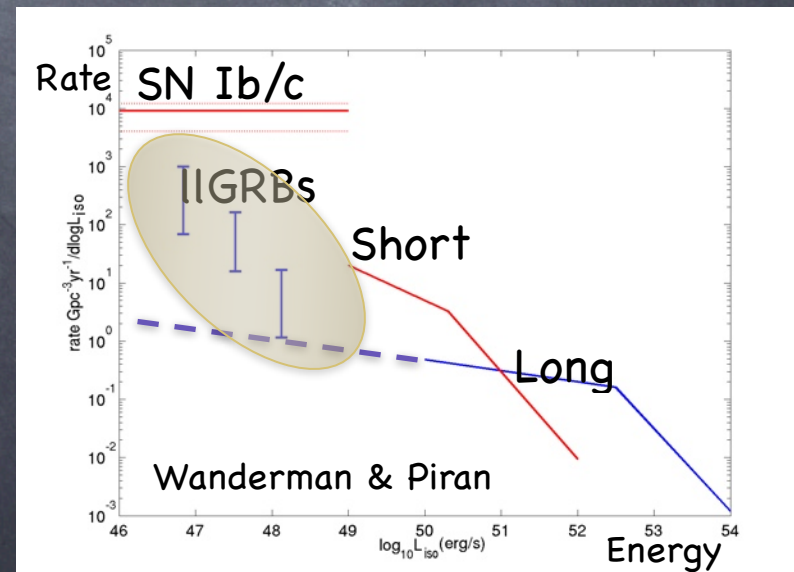
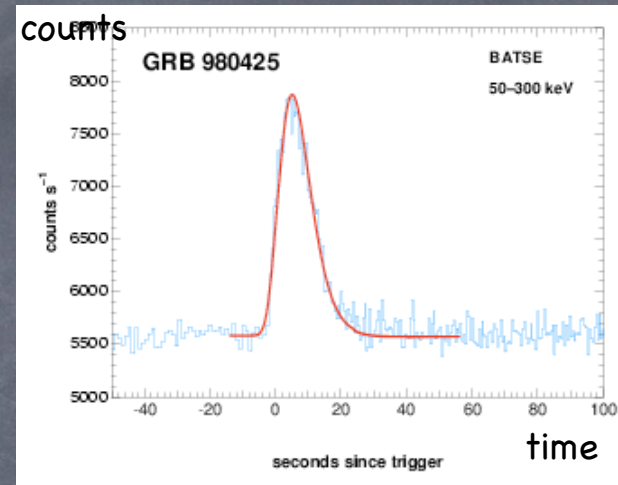
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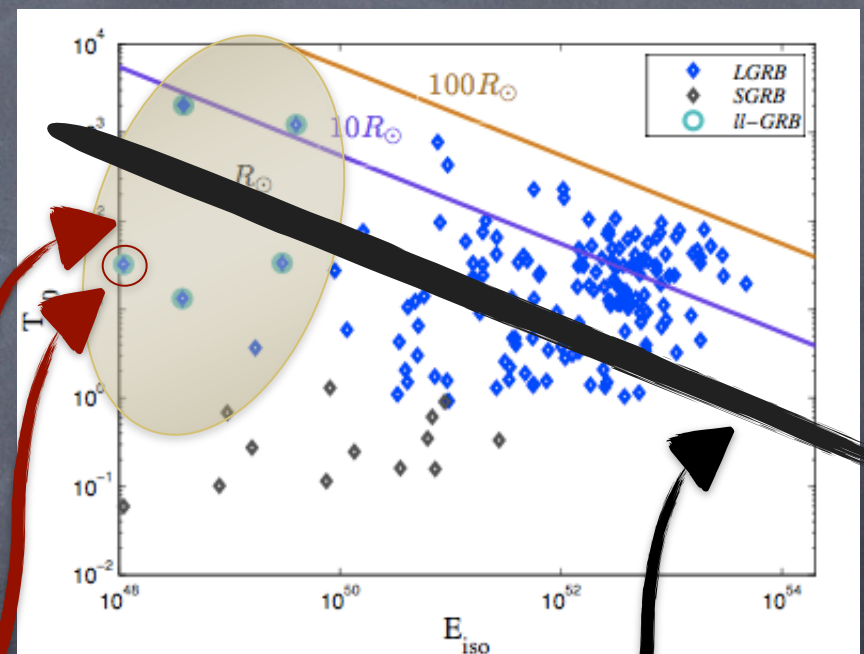
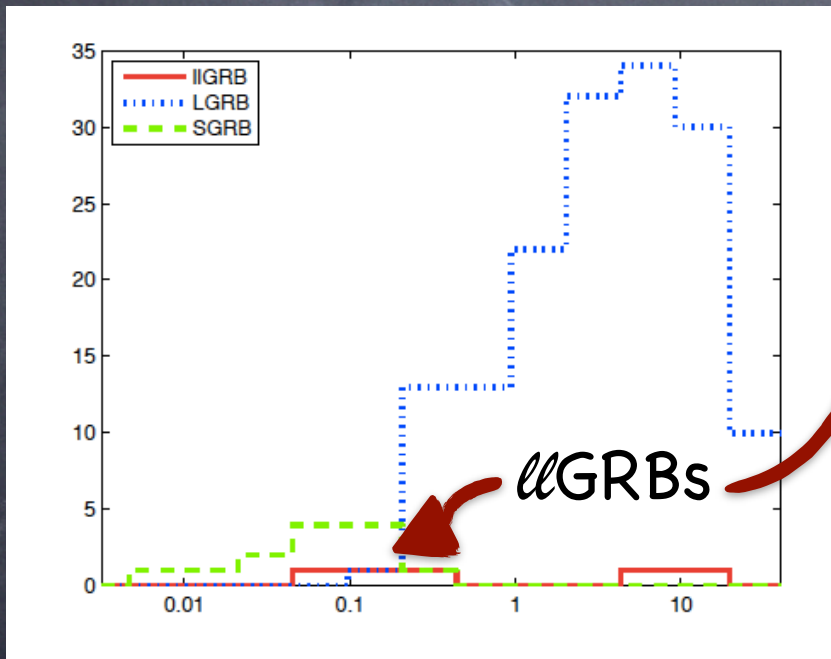
T_B

Low Luminosity GRBs – *ll*GRBs

Bromberg Nakar, TP, 11 ApJL 2011

- Low luminosity GRBs:
 - $E_{\text{iso}} \sim 10^{48} - 10^{49}$ ergs
 - Smooth single peaked light curve.
 - Soft Emission ($E_{\text{peak}} < 150$ keV)
 - Much more numerous than regular long GRBs!
 - *ll*GRBs don't have enough power to penetrate the star

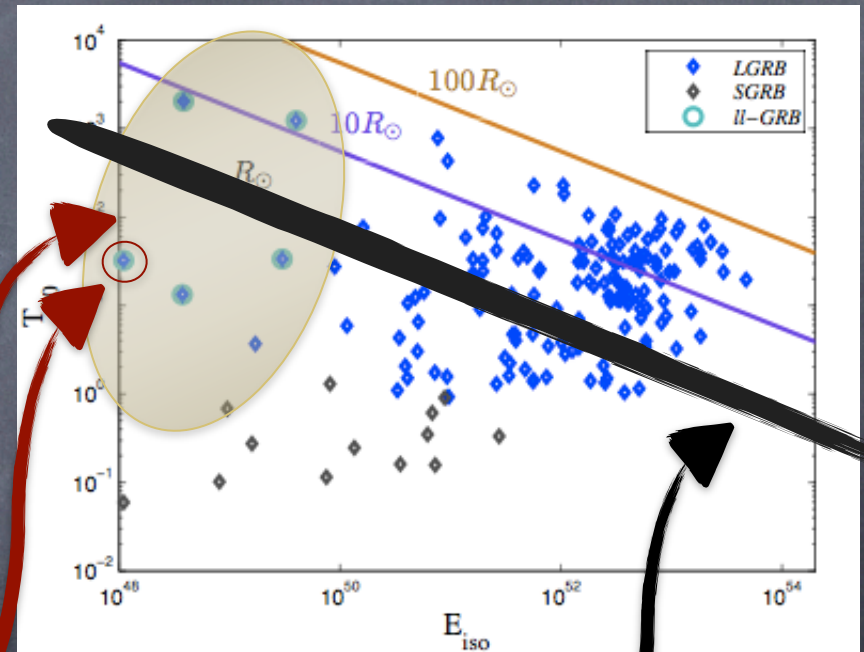
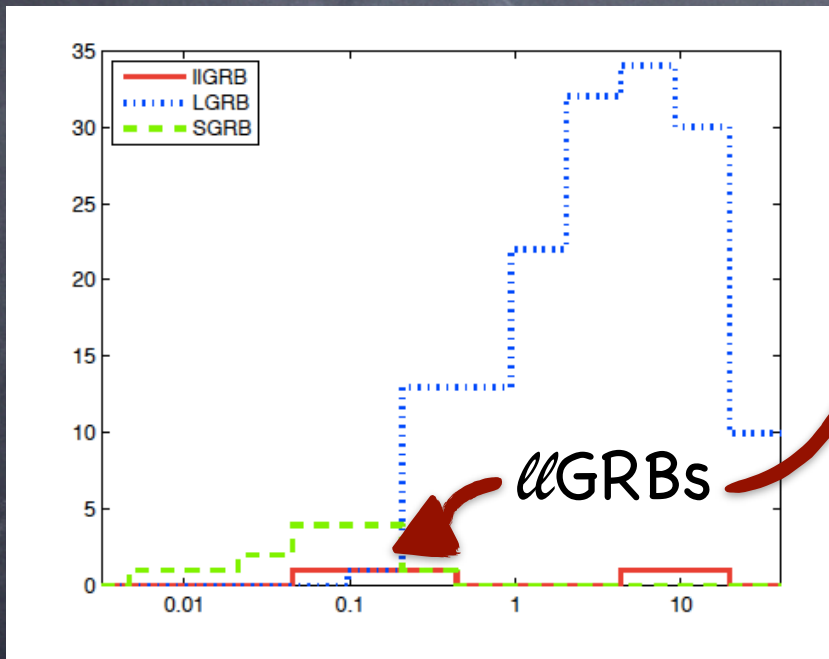




98bw

T_B

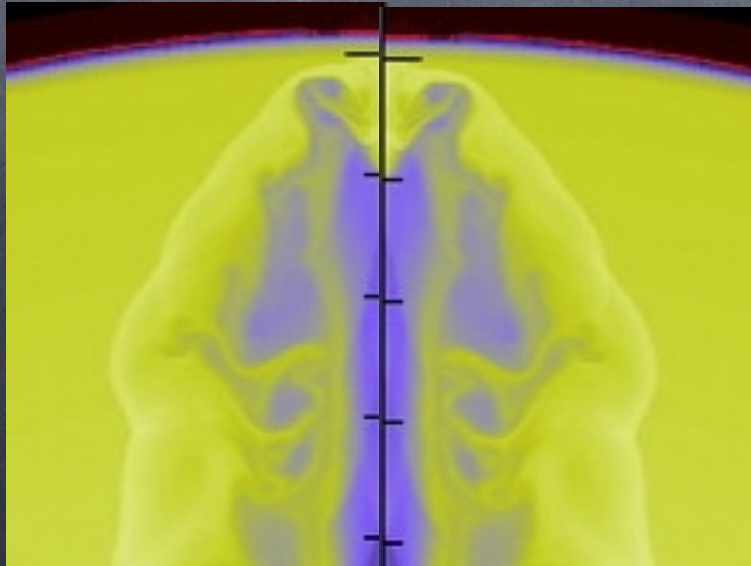
Low luminosity GRBs – ℓ GRBs cannot arise from Collapsars



98bw

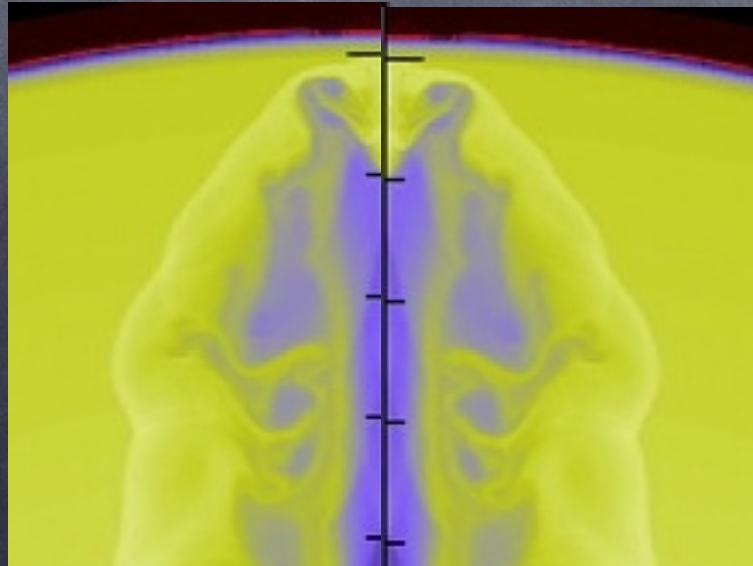
T_B

What makes a ℓ GRBs?



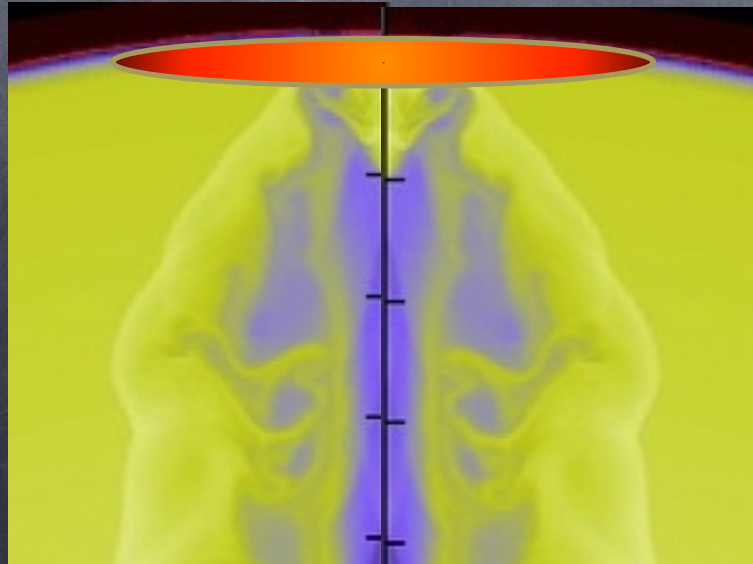
What makes a ℓ GRBs?

- 👁 A weak jet that fails to break out (“a failed GRB”).



What makes a ℓ GRBs?

- A weak jet that fails to break out (“a failed GRB”).
- We observe the shock breakout from the stellar envelope (Colgate, 1967; Katz, Budnik, Waxman, 2010; Nakar & Sari, 2011)



Almost ALL GRBS
accompanied by SNe
are
*ll*GRBs
?

A prediction of the Collapsar model

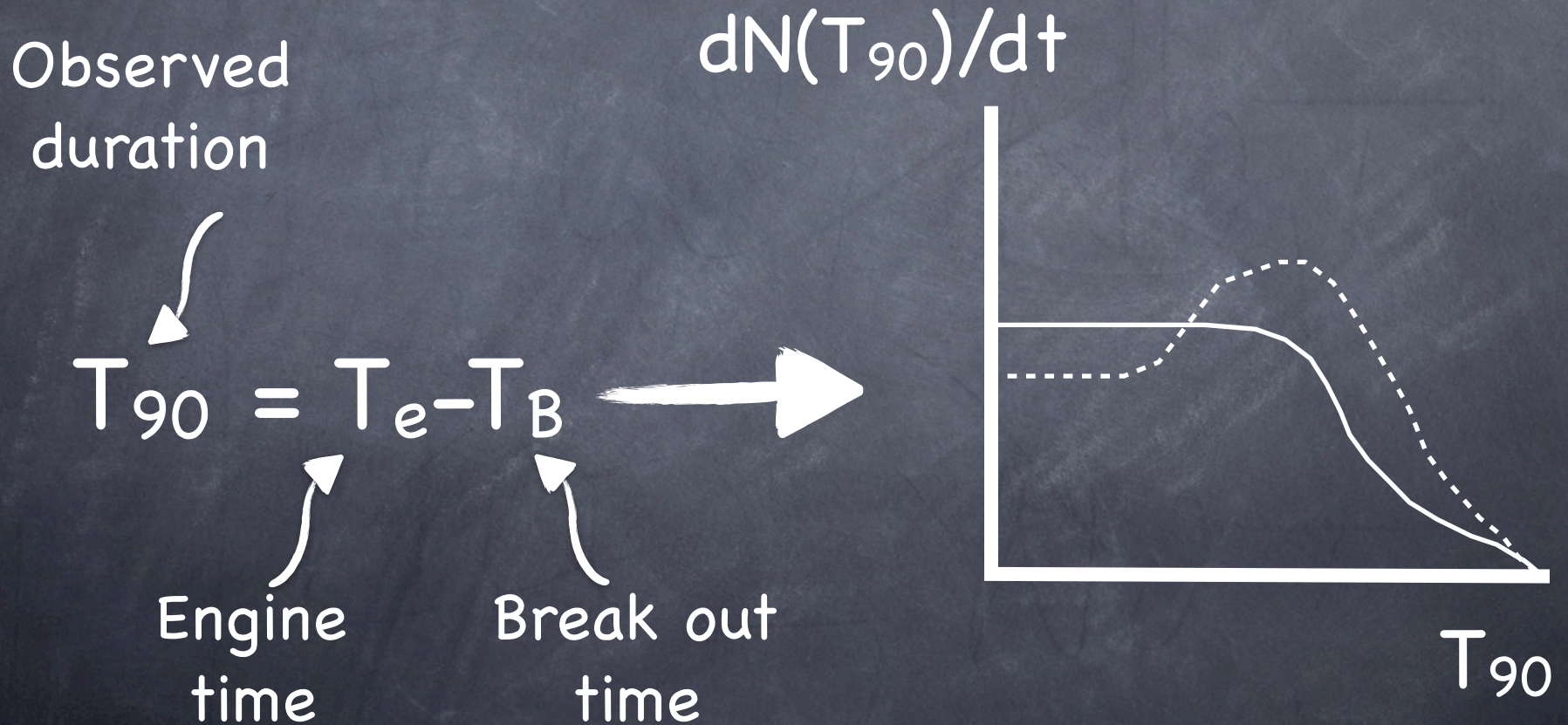
Observed
duration

$$T_{90} = T_e - T_B$$

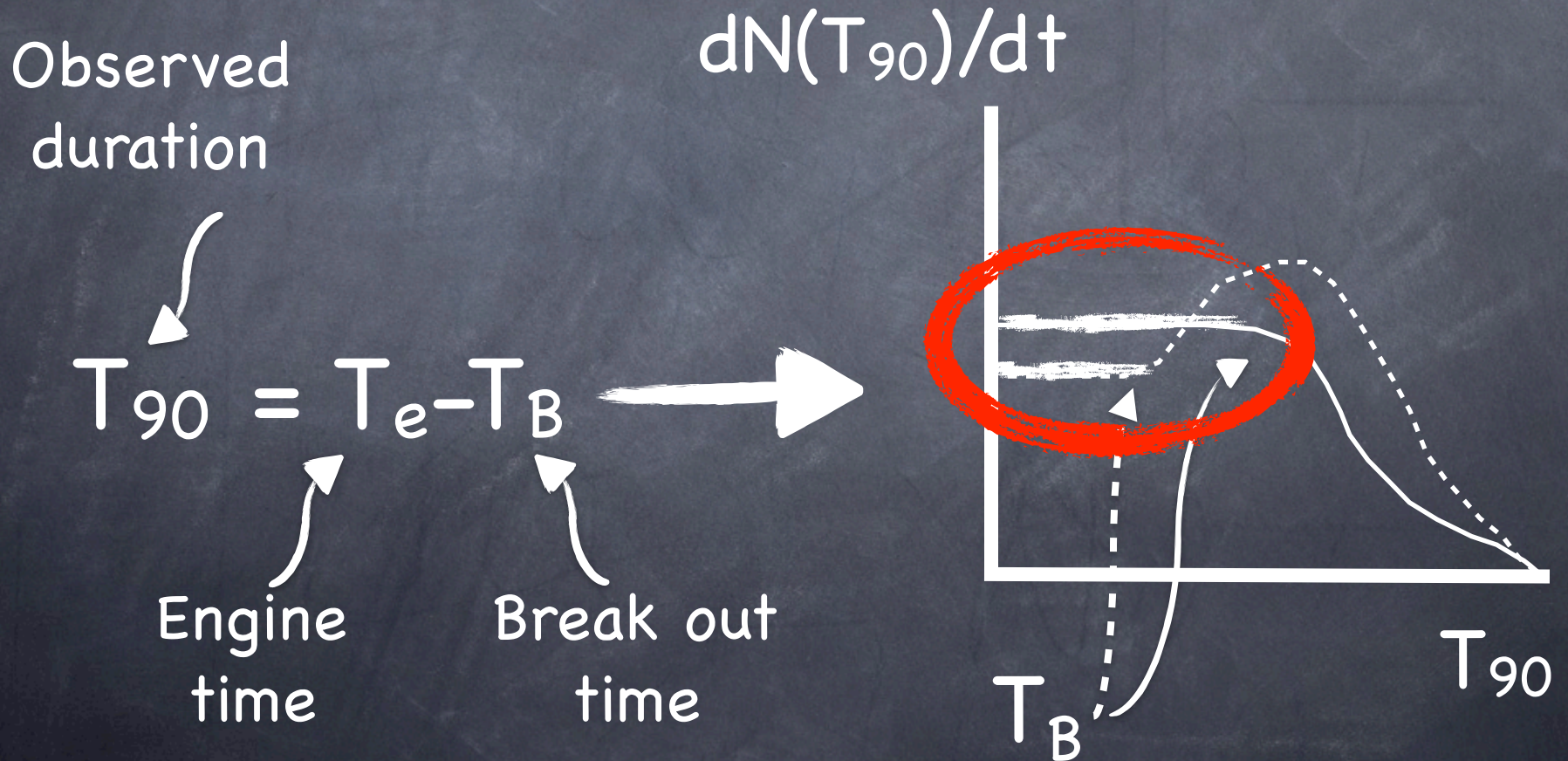
Engine
time

Break out
time

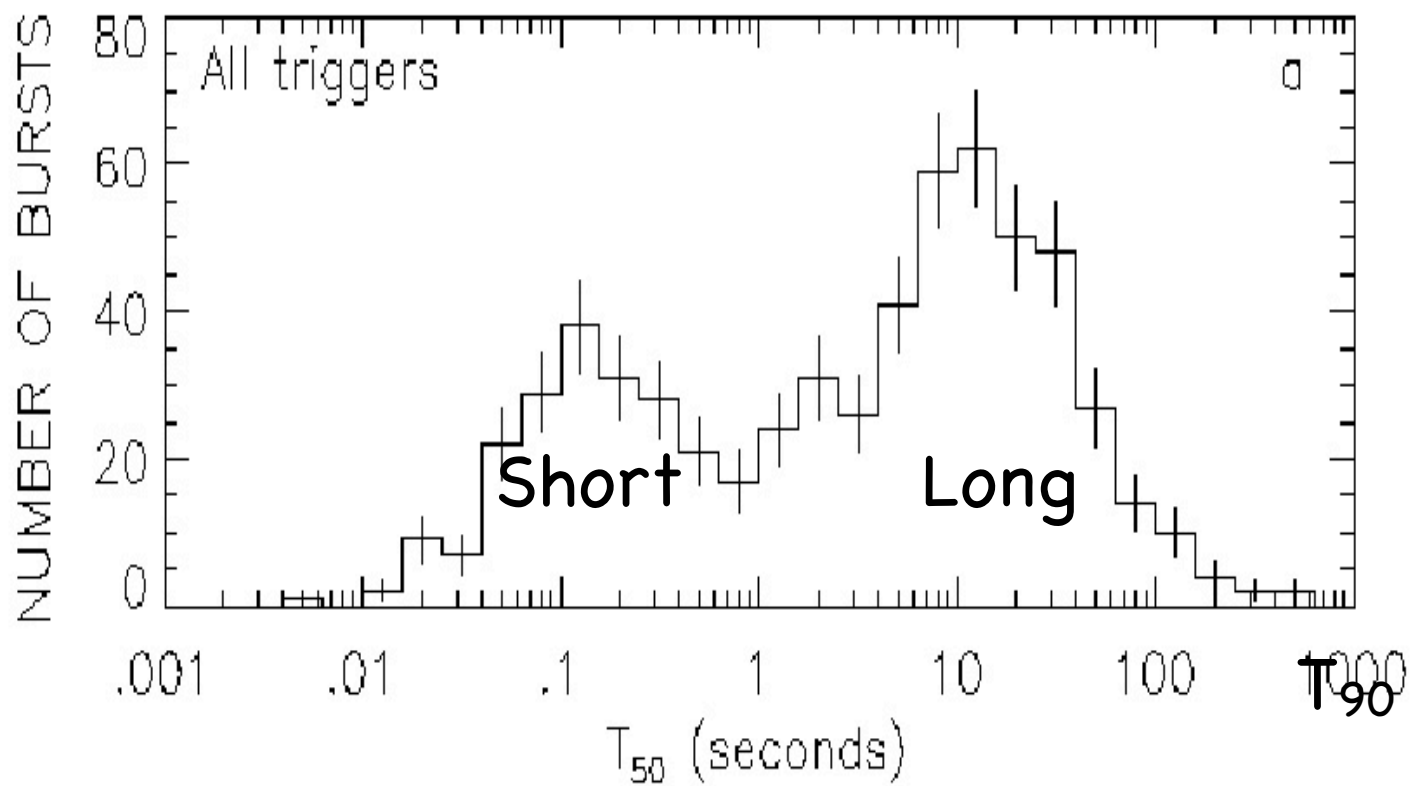
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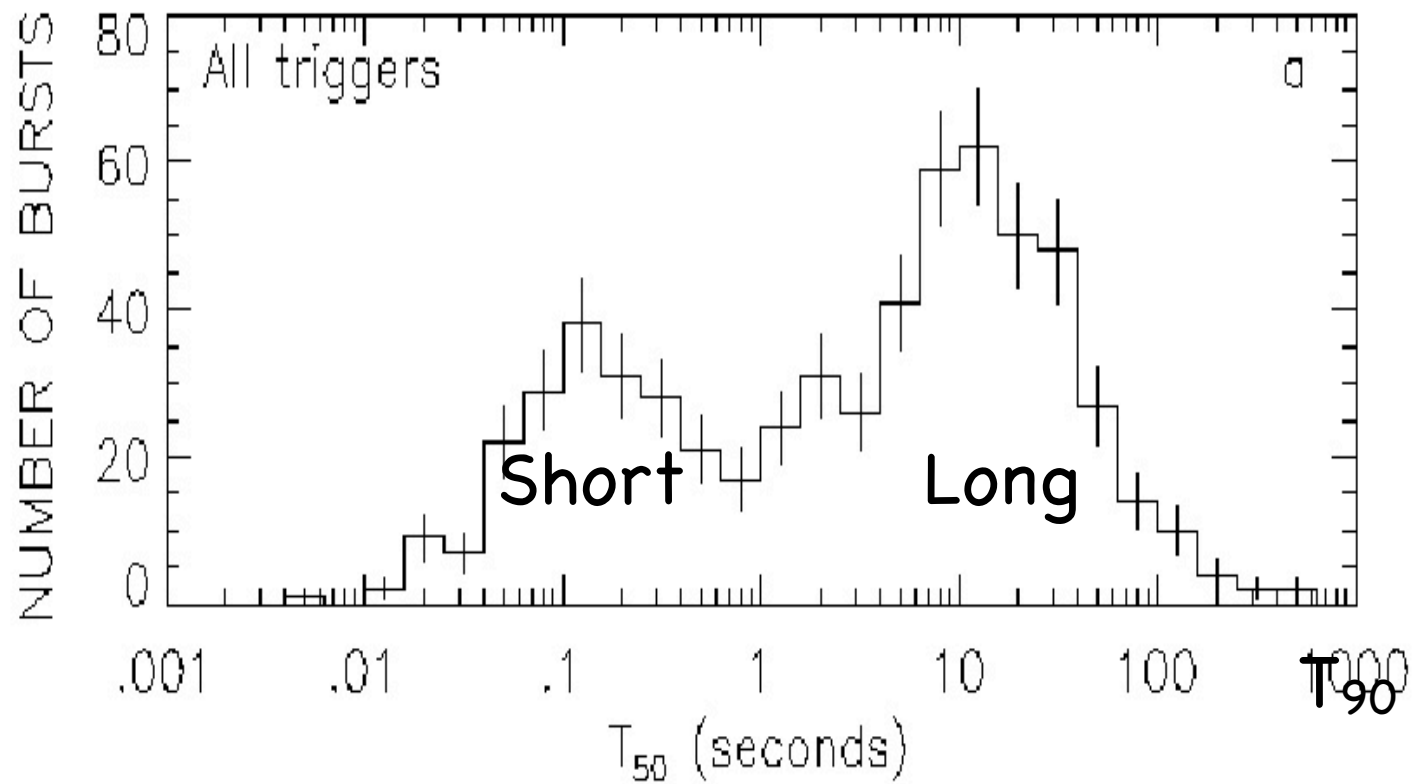


?



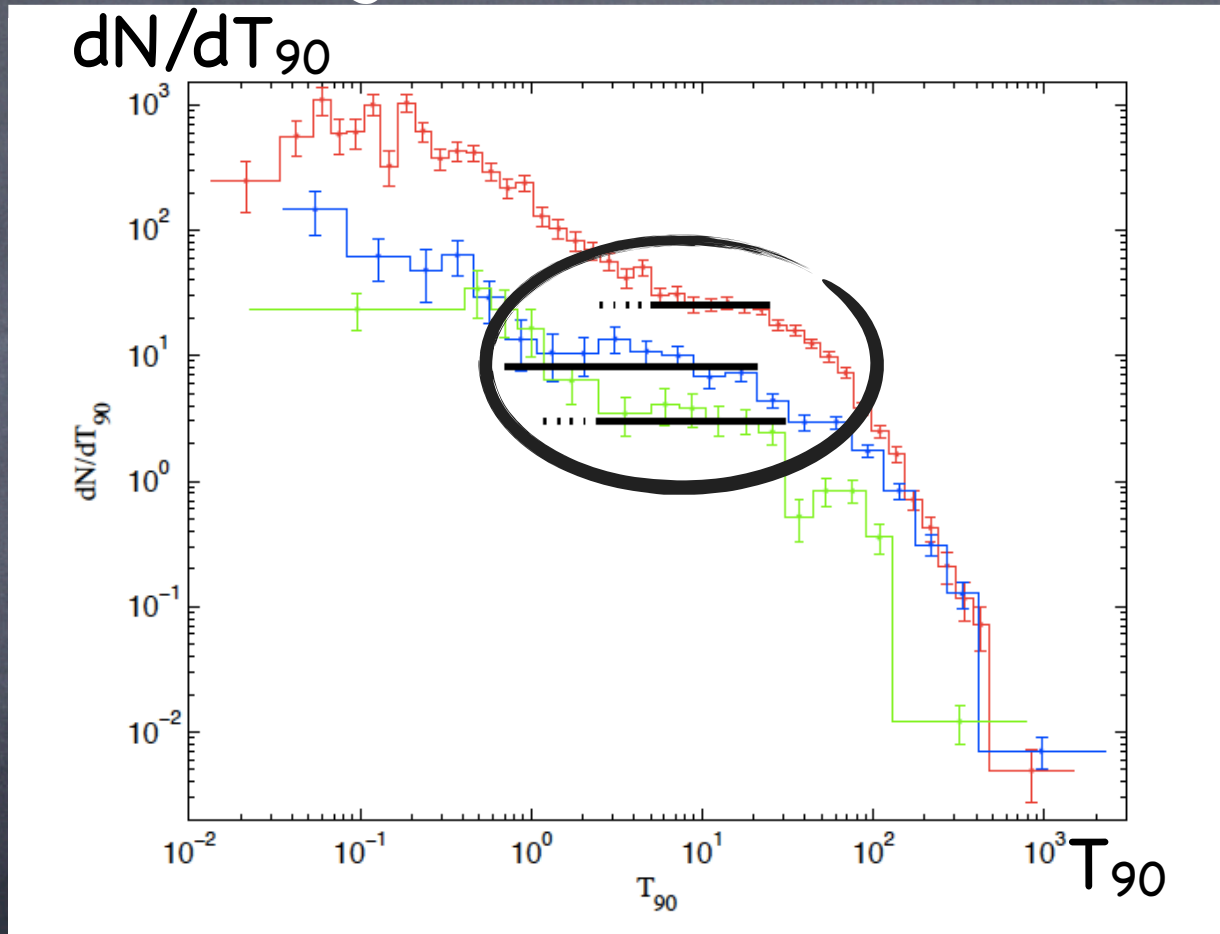
?

$d\log(N)/dT_{90}$



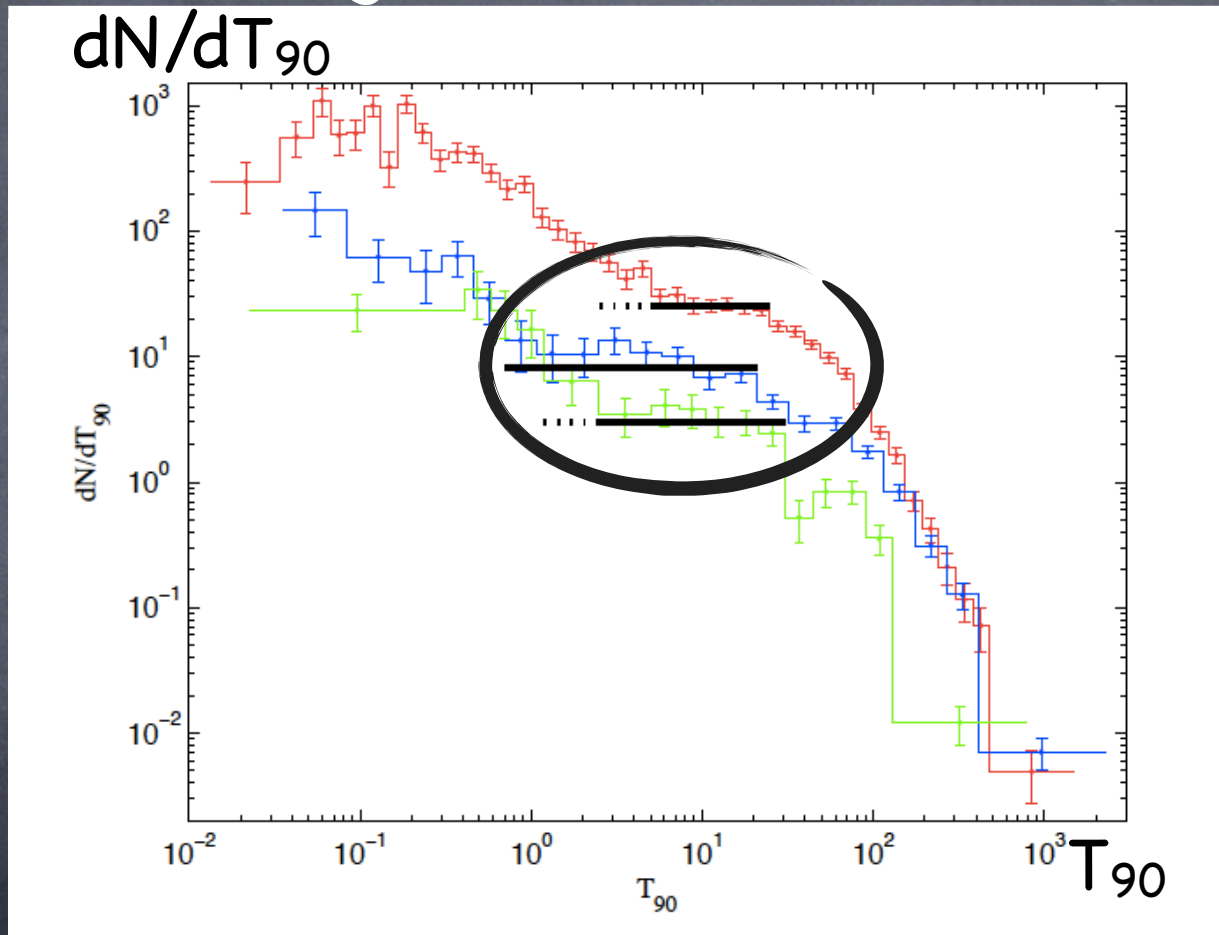
A second look

(Bromberg Nakar, TP & Sari, 2011)



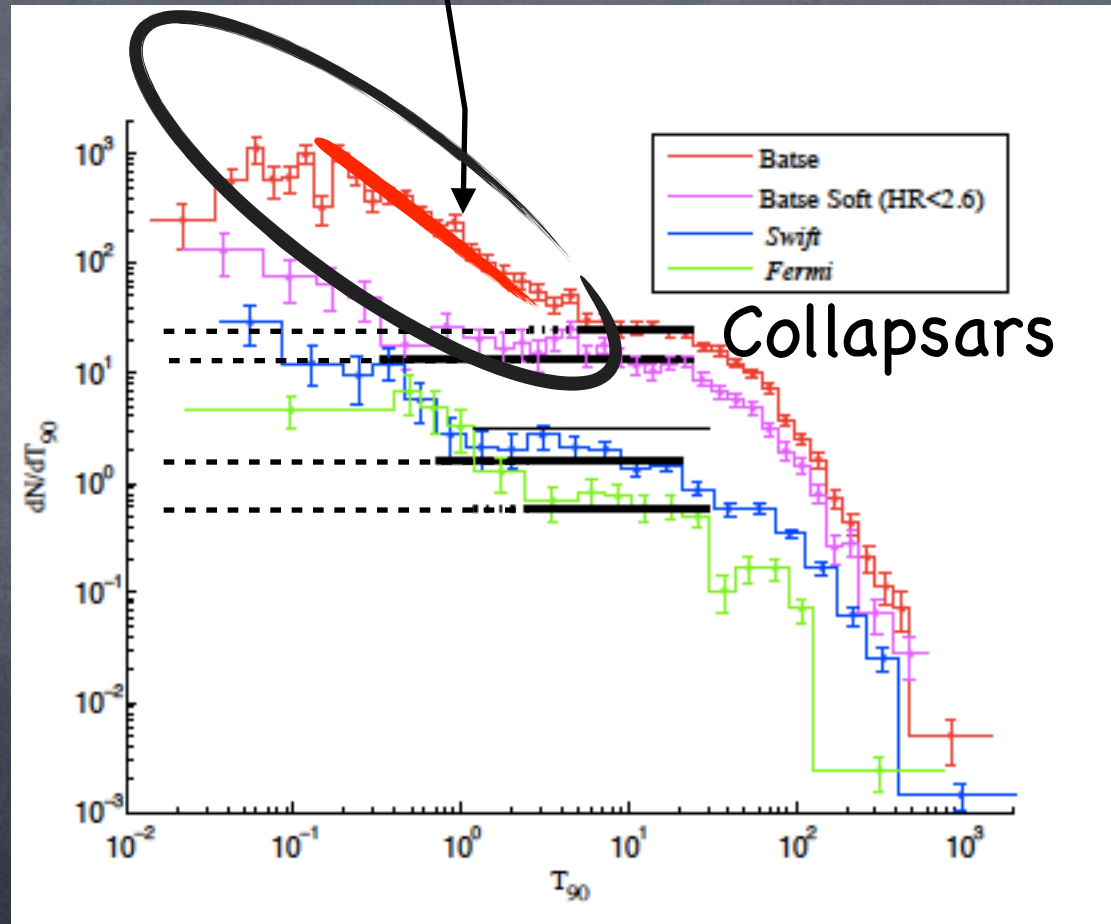
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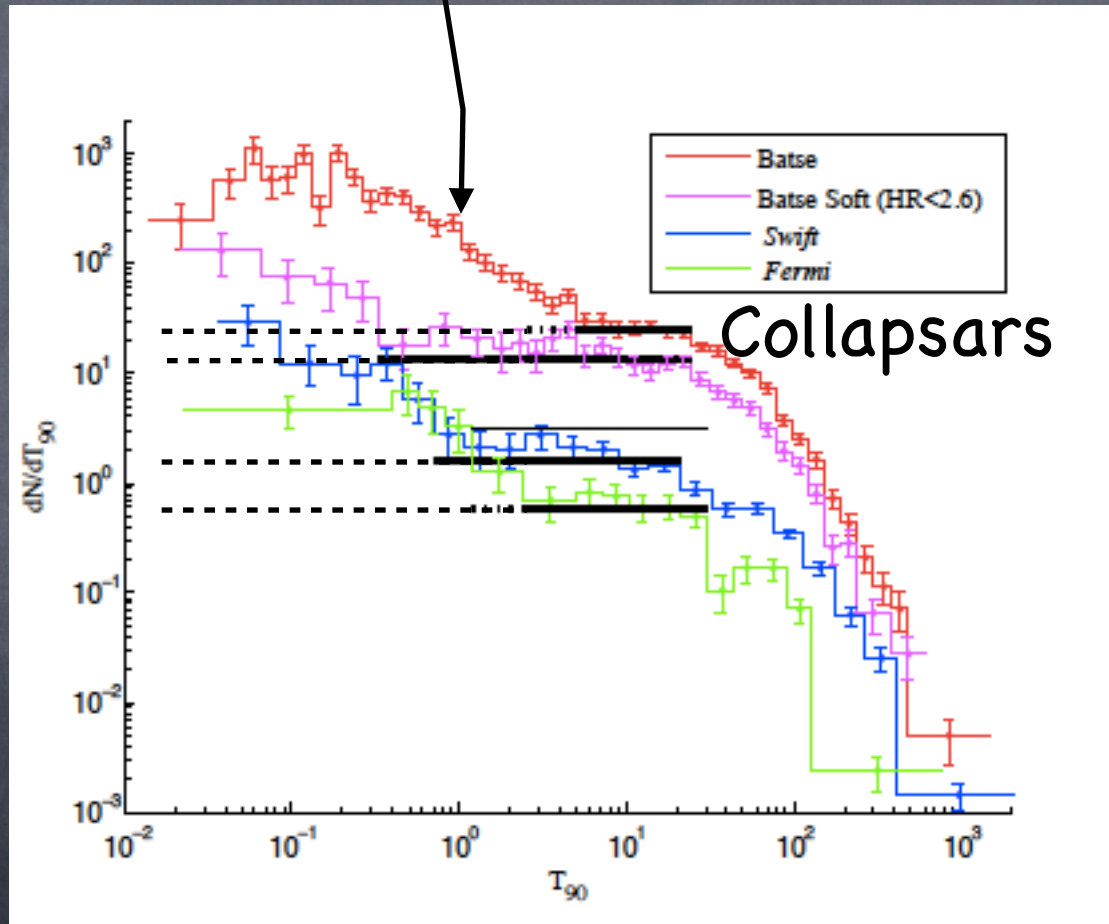


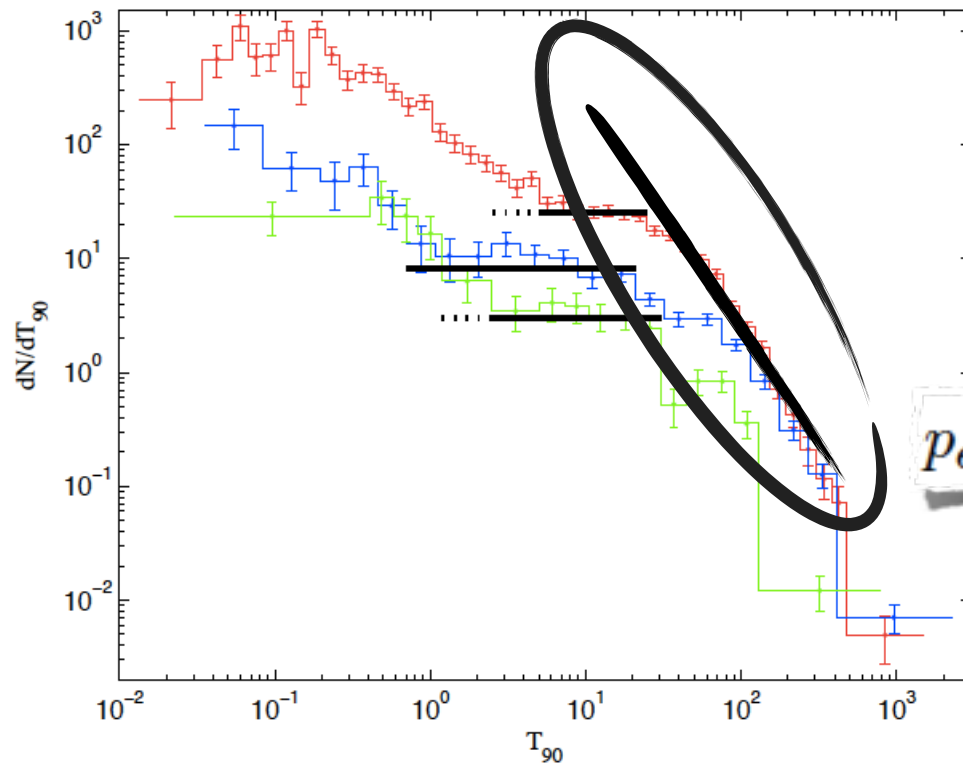
This provides a direct observational proof
of the Collapsar model.

Short (Non-Collapsars) GRBs



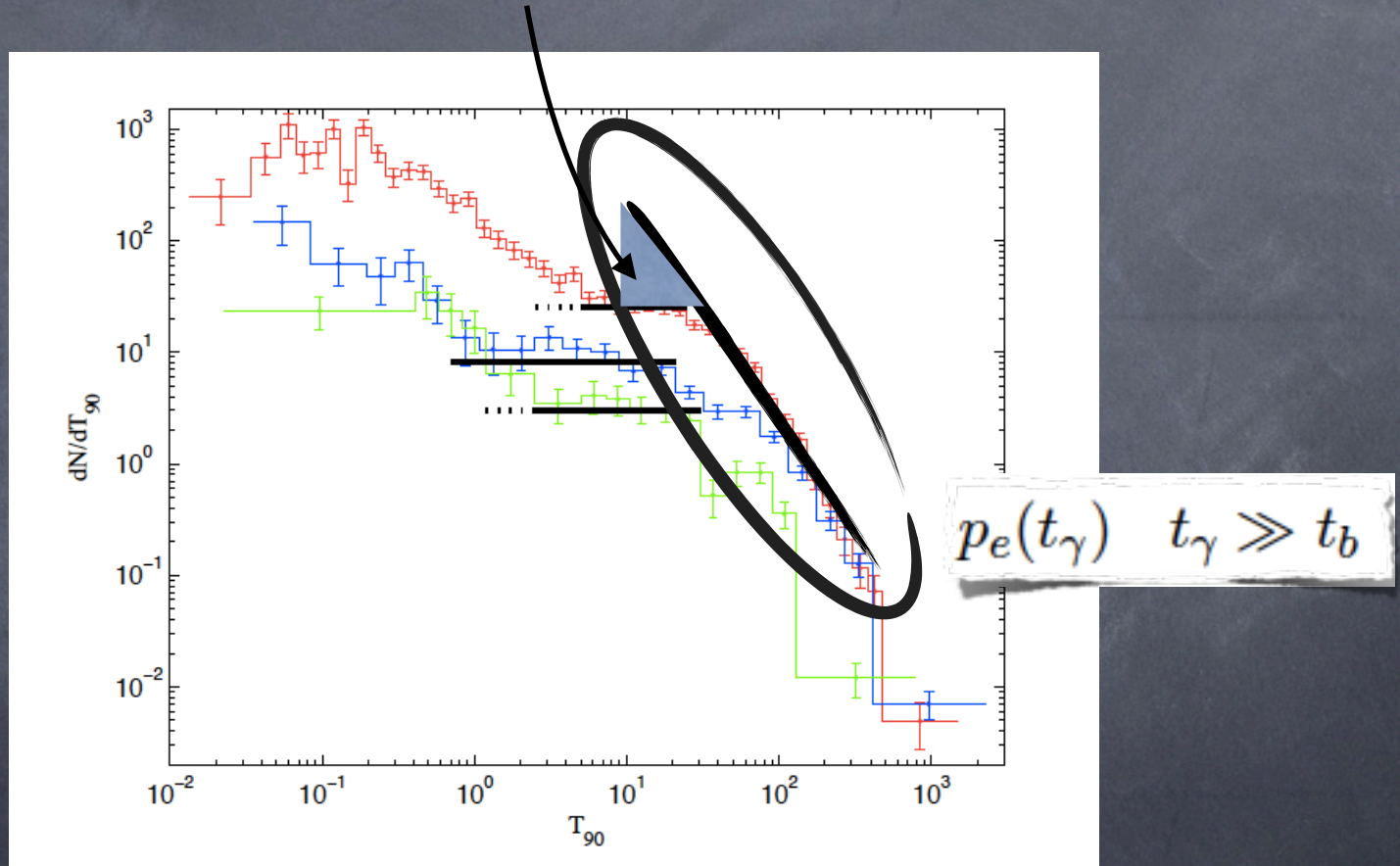
Short (Non-Collapsars) GRBs





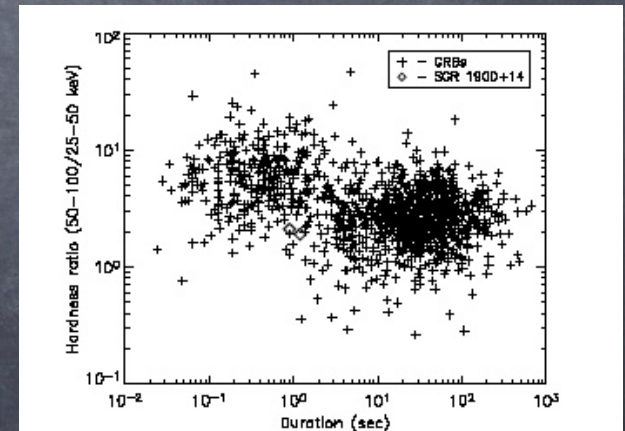
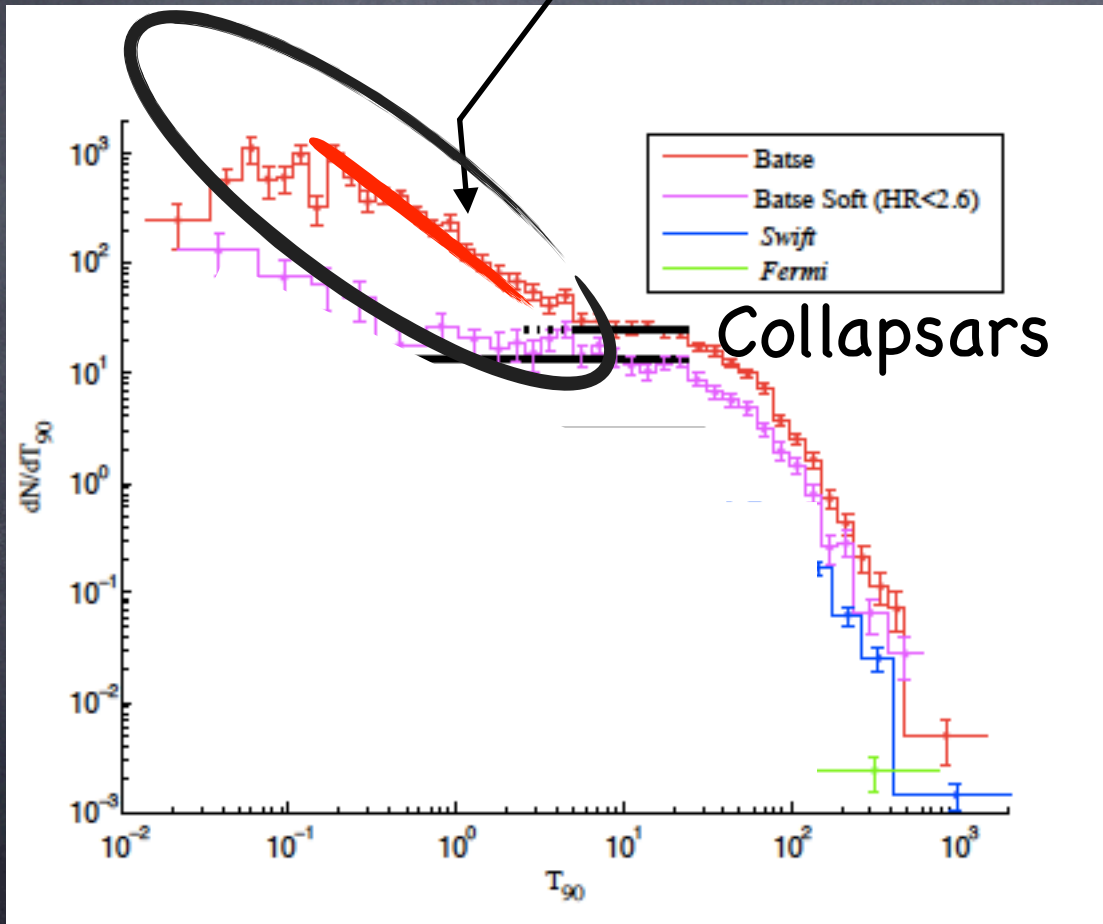
$$p_e(t_\gamma) \quad t_\gamma \gg t_b$$

A large number of “failed or Choked” jets –
a “failed GRB” ✓

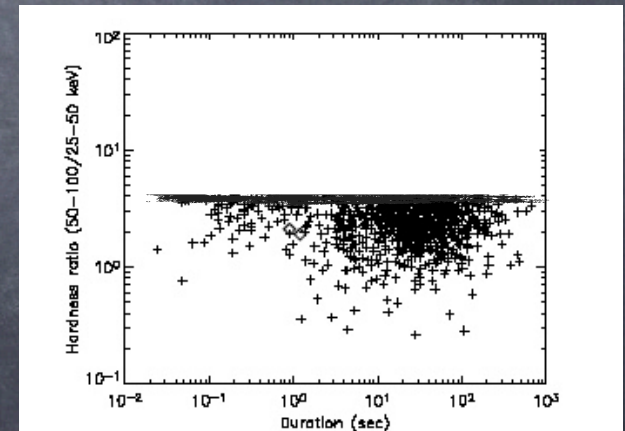
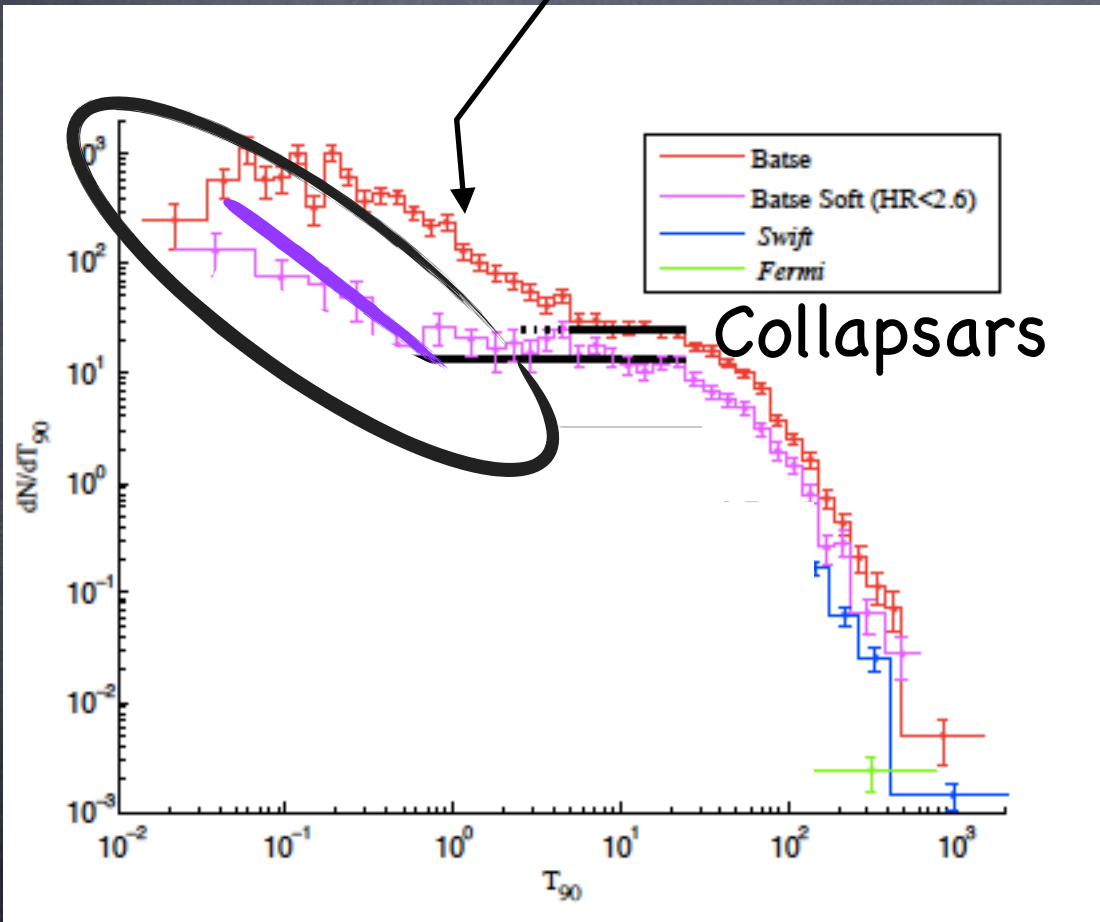


The rate of ℓ GRBs is very large

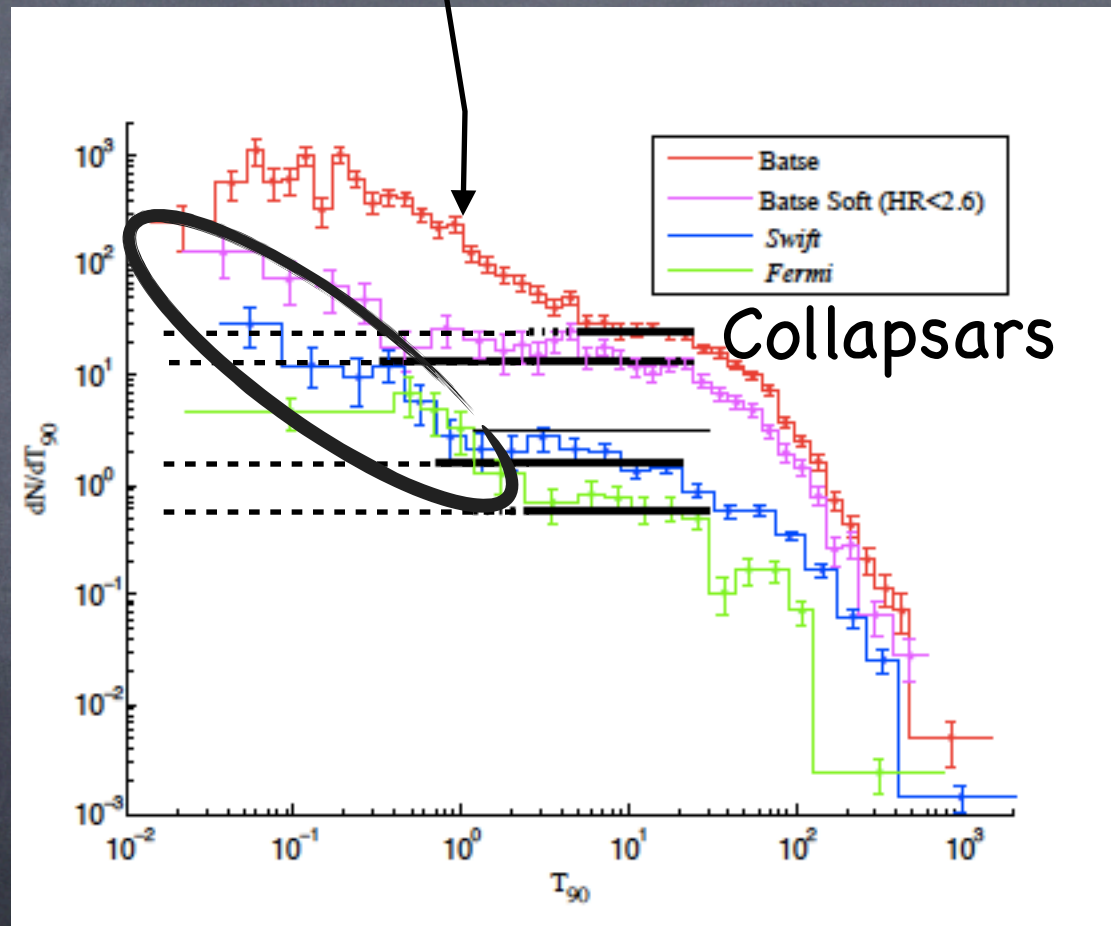
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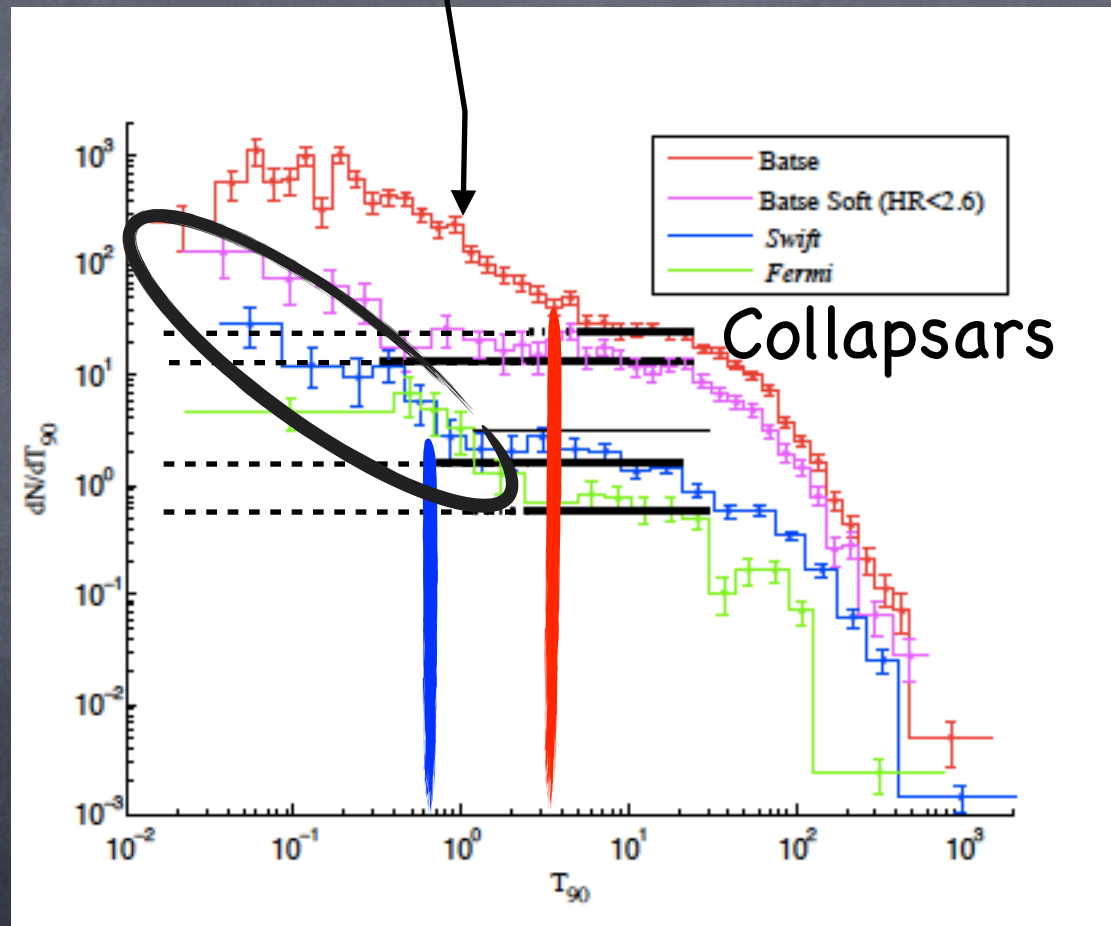
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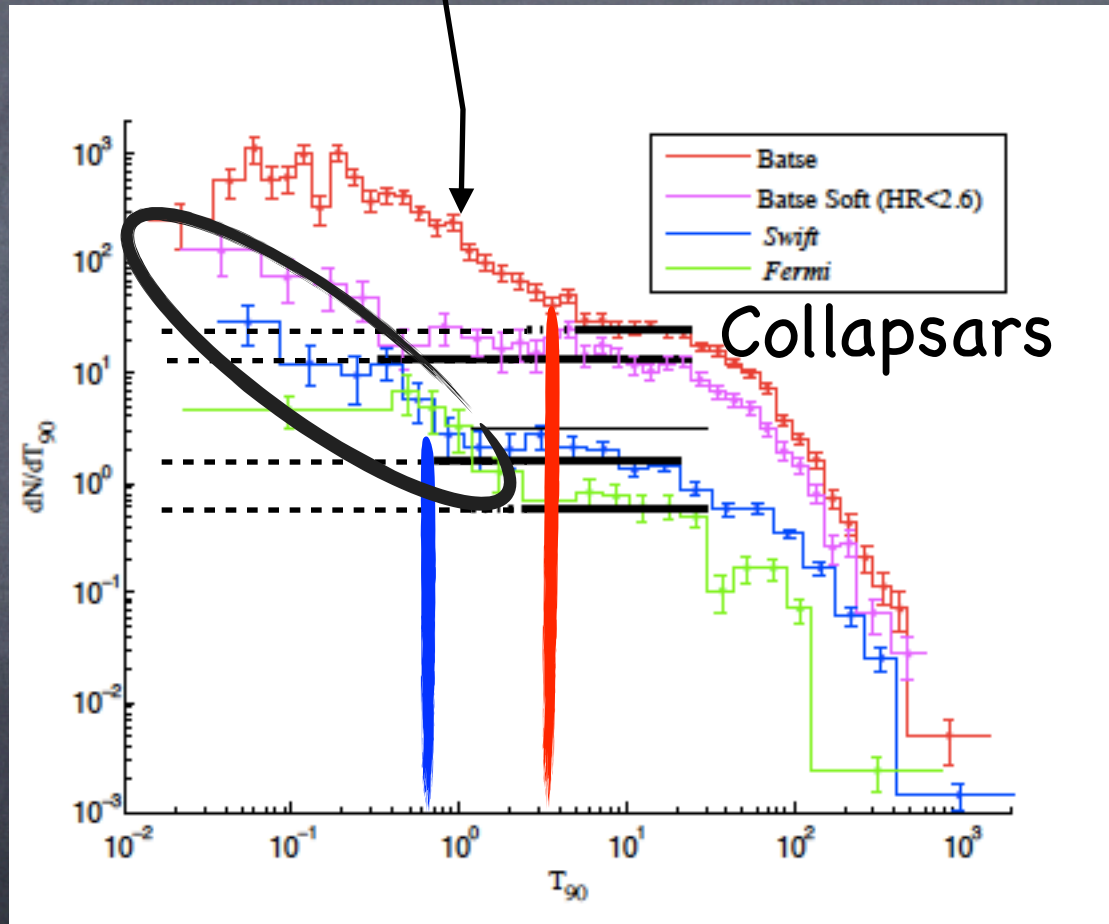
Swift Short (Non-Collapsars) GRBs



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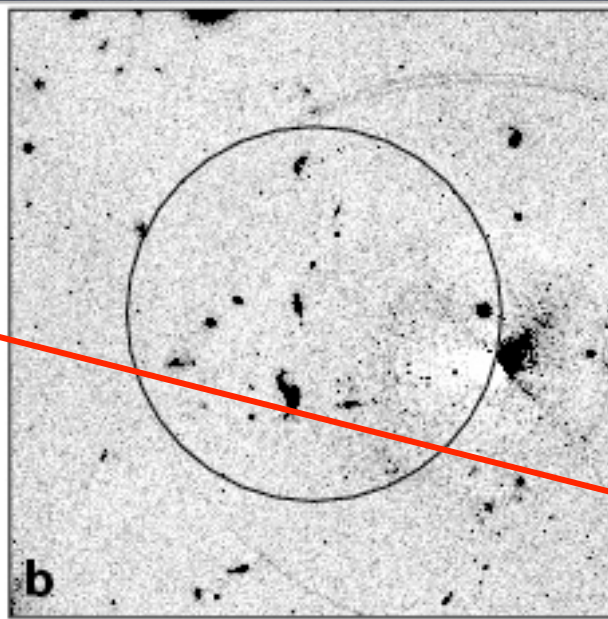
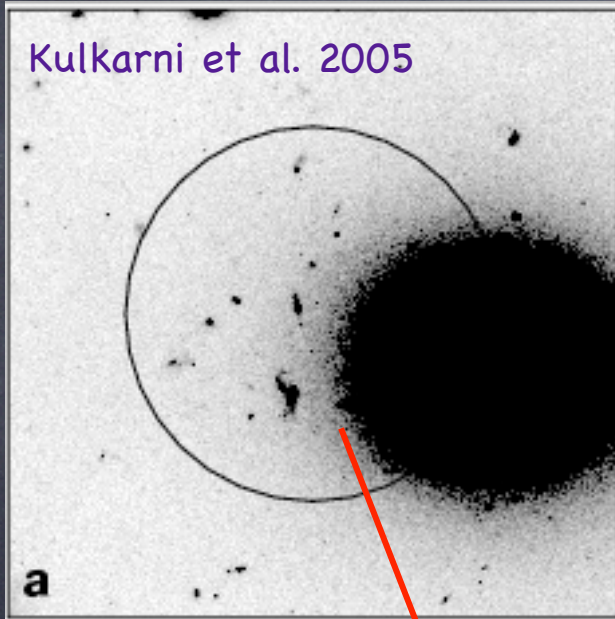
Short Swift GRBs with $T_{90} > 0.7$ sec are not "short"!

Short GRBs – GRB 050509b

Swift/XRT position intersects a bright elliptical at $z = 0.226$

No optical/radio afterglow

Kulkarni et al. 2005

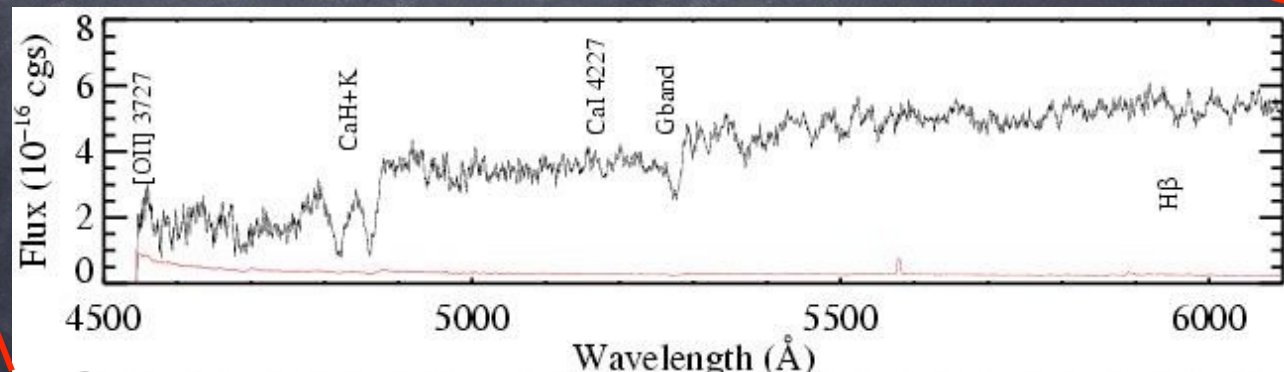


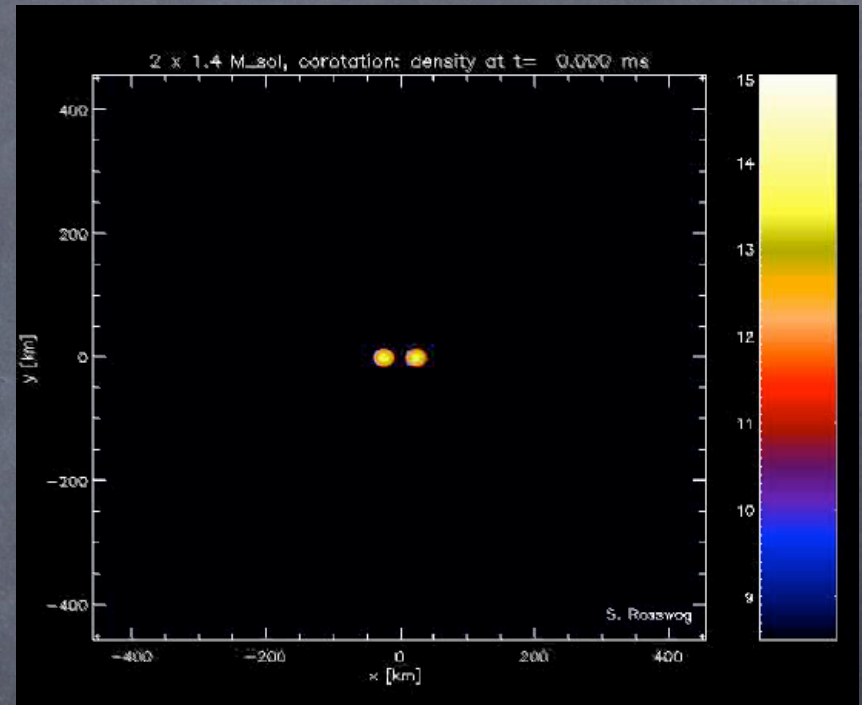
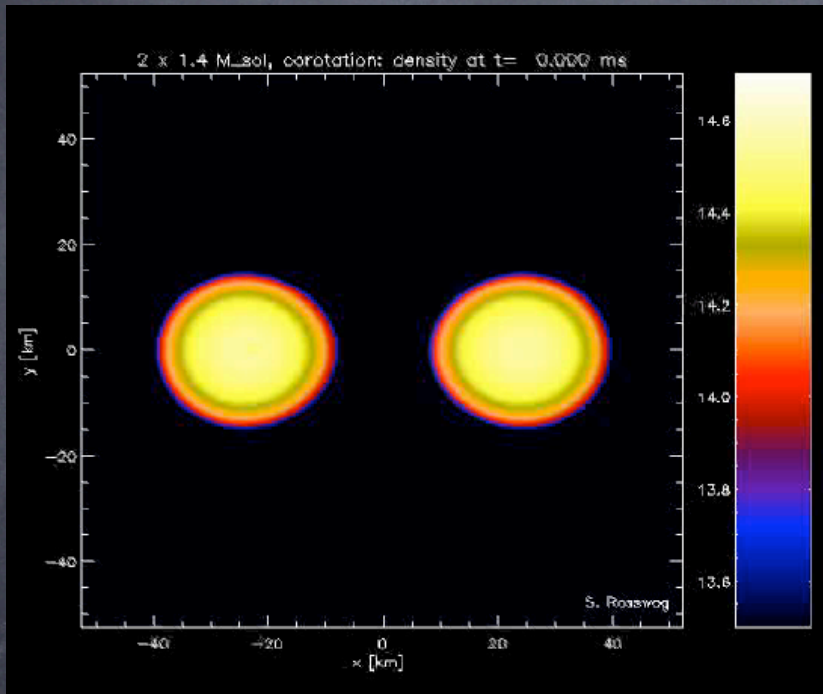
Elliptical host



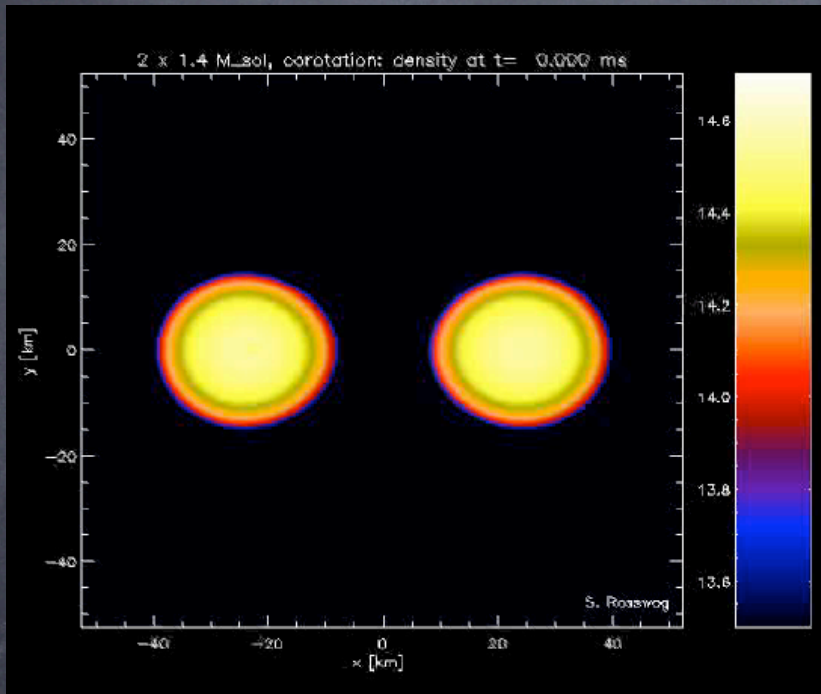
Old stellar
population

Bloom et al. 2005
Castro-Tirado et al.
2005
Gehrels et al. 2005
Hjorth et al. 2005



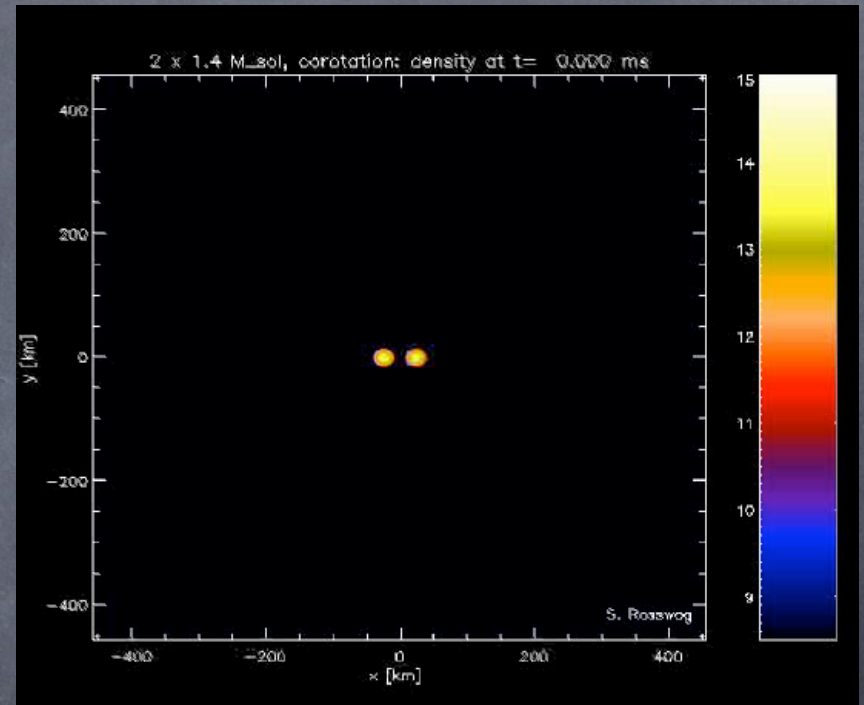


Price & Rosswog

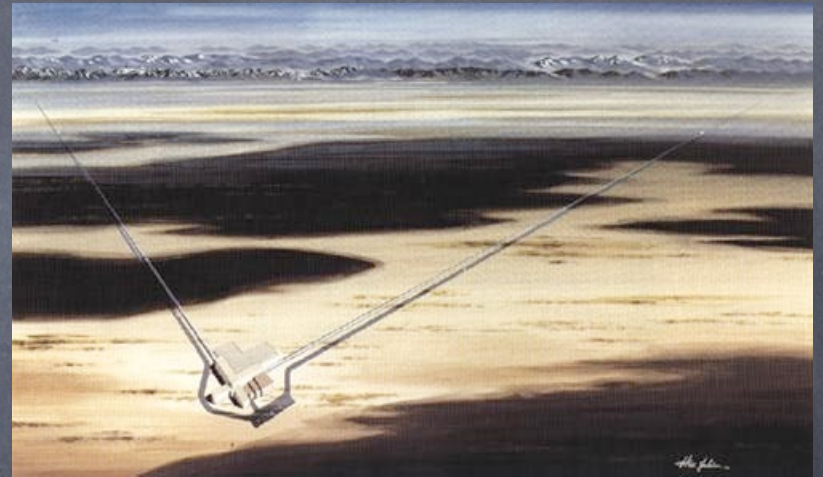
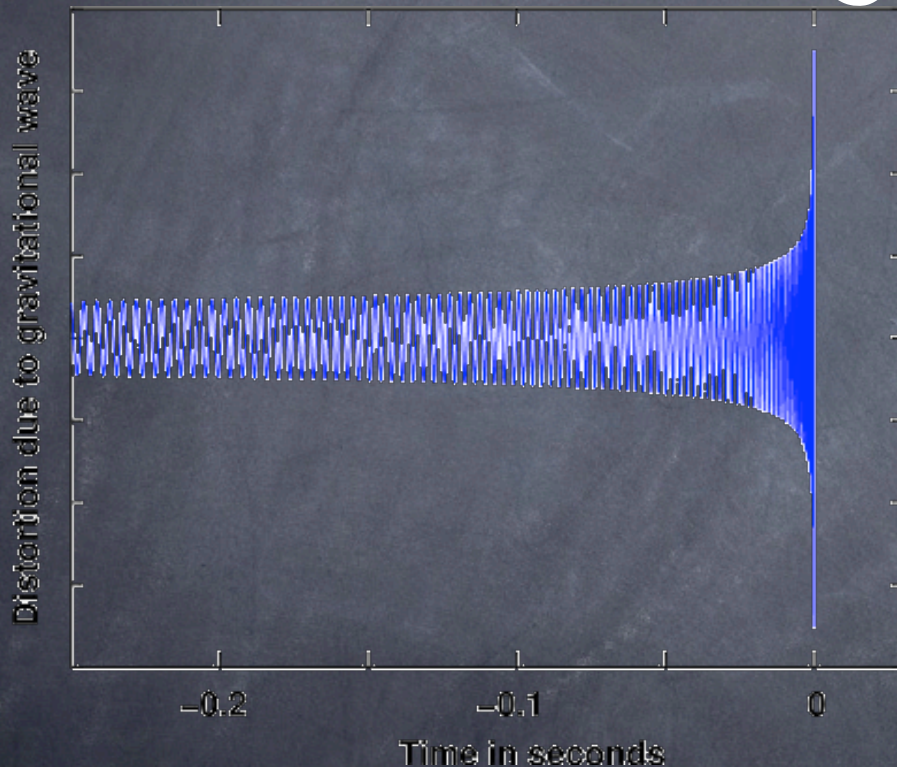


Price & Rosswog

Price & Rosswog



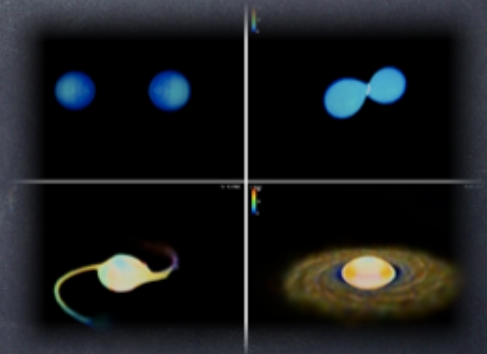
Short GRBs-NS Mergers ?



To be confirmation with detection of
Gravitational radiation

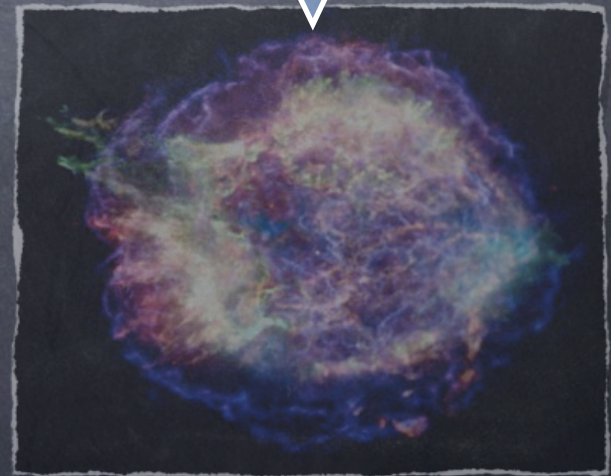
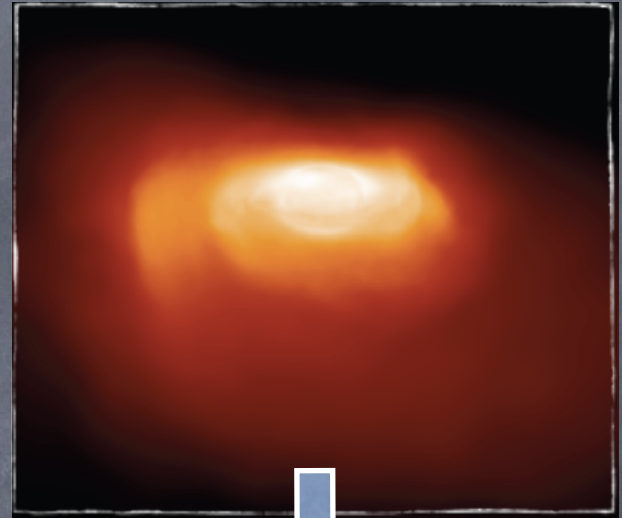
II. Radio Flares from Neutron Star merges – The Electromagnetic signals that follow the Gravitational Waves

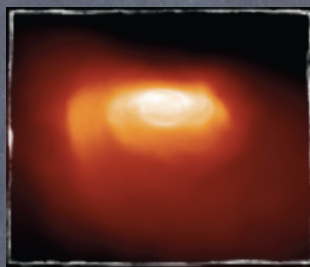
Nakar + TP, Nature 2011, Nakar, TP & Rosswog in
preparation

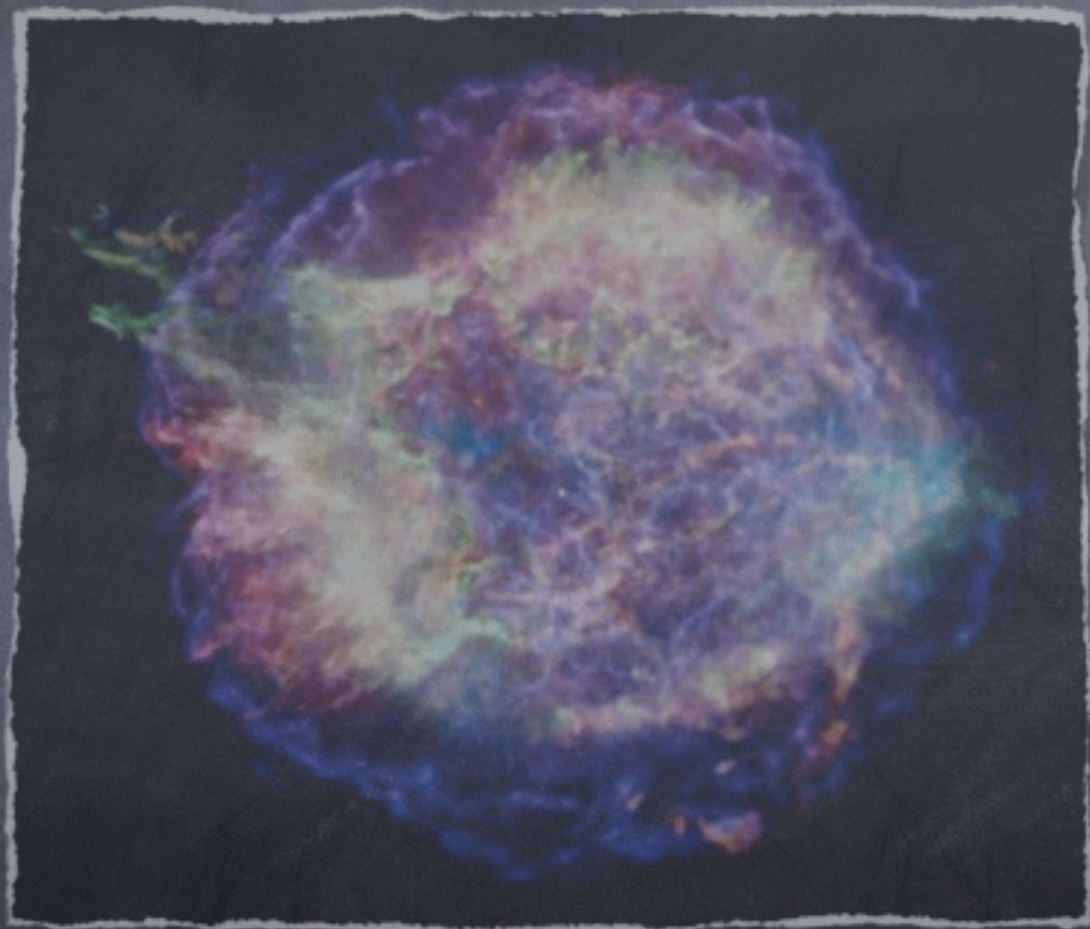


The Model

Numerous numerical simulations show that NS merger eject Sub - or Mildly relativistic outflow with $E \sim 10^{49}$ erg
Lorentz factor $(\Gamma - 1) \approx 1$
Interaction of the outflow with the ISM

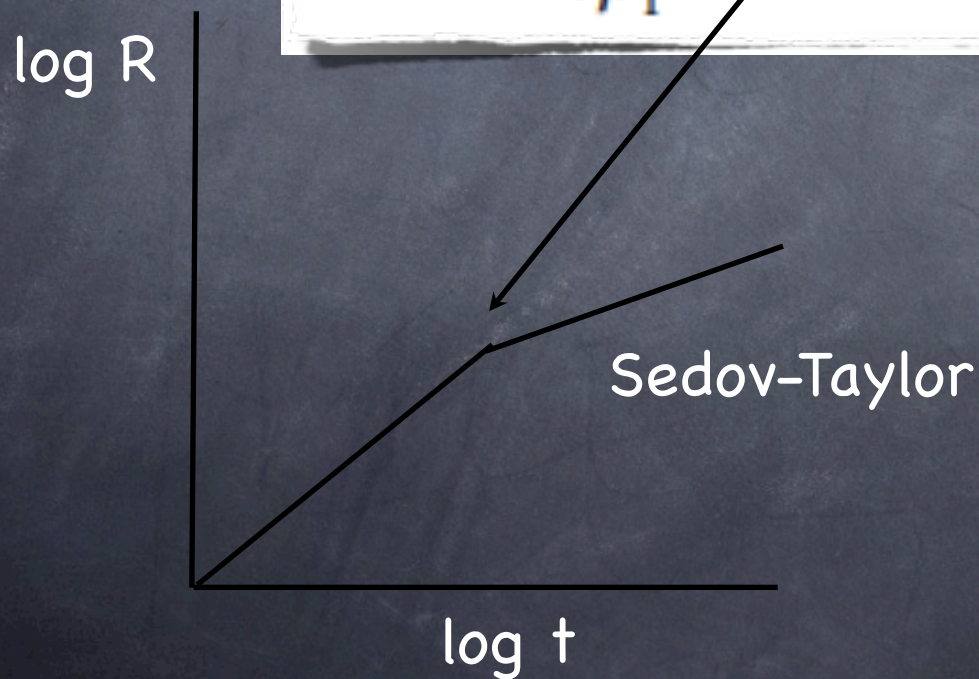






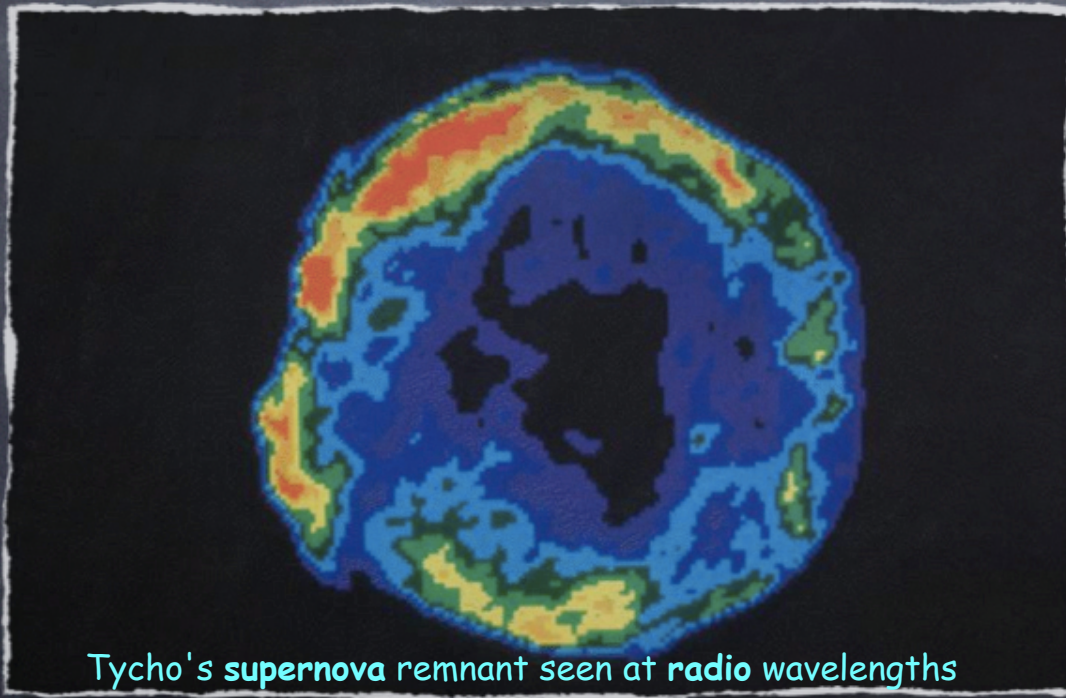
Dynamics

$$t_{\text{dec}} = \frac{R_{\text{dec}}}{c\beta_i} \approx 30 E_{49}^{1/3} n_0^{-1/3} \beta_i^{-5/3} \text{ days}$$



Radio Supernova

e.g. 1998bw (Chevalier 98)



$$e_e = \xi_e e$$

$$e_B = B^2 / 8\pi = \xi_B e$$

$$N(\gamma) \propto \gamma^{-p} \quad \text{for } \gamma > \gamma_m$$

$$p = 2.5 - 3$$

$$\gamma_m = (m_p / m_e) e_e (\Gamma - 1)$$

$$\nu = (3/4\pi) e B \gamma^2$$

$$F_\nu = (\sigma_T c / e) N_e B$$

Frequency and Intensity

(Nakar & TP Nature, 2011)

$$\nu_{m,dec} \equiv \nu_m(t_{dec}) \approx 1 \text{ GHz } n^{1/2} \epsilon_{B,-1}^{1/2} \epsilon_{e,-1}^2 (\Gamma_0 - 1)^{5/2},$$

$$F_{\nu_{obs},peak} [\nu_{obs} > \nu_{m,dec}, \nu_{a,dec}] \approx$$

$$0.3 E_{49} n_0^{\frac{p+1}{4}} \epsilon_{B,-1}^{\frac{p+1}{4}} \epsilon_{e,-1}^{p-1} \beta_i^{\frac{5p-7}{2}} d_{27}^{-2} \left(\frac{\nu_{obs}}{1.4} \right)^{-\frac{p-1}{2}}$$

$$N_{all-sky}(1.4\text{GHz}) \approx 20 E_{49}^{11/6} n^{\frac{9p-1}{24}} \epsilon_{B,-1}^{\frac{3(p+1)}{8}} \epsilon_{e,-1}^{\frac{3(p-1)}{2}} (\Gamma_0 - 1)^{\frac{45p-83}{24}} \mathcal{R}_{300} F_{lim,-1}^{-3/2} .$$

$$N_{all-sky}(1.4\text{GHz}) \approx 20 E_{49}^{11/6} n^{\frac{9p-1}{24}} \epsilon_{B,-1}^{\frac{3(p+1)}{8}} \epsilon_{e,-1}^{\frac{3(p-1)}{2}} (\Gamma_0 - 1)^{\frac{45p-83}{24}} \mathcal{R}_{300} F_{lim,-1}^{-3/2} .$$



$$N_{all-sky}(1.4\text{GHz}) \approx 20 E_{49}^{11/6} n^{\frac{9p-1}{24}} \epsilon_{B,-1}^{\frac{3(p+1)}{8}} \epsilon_{e,-1}^{\frac{3(p-1)}{2}} (\Gamma_0 - 1)^{\frac{45p-83}{24}} \mathcal{R}_{300} F_{lim,-1}^{-3/2} .$$

Detectability


Table 1 | Observing radio flares

Radio facility	Observing frequency (GHz)	Field of view (deg ²)	One-hour r.m.s.* (μJy)	One-hour detection horizon†	
				$\beta_1 \approx 1, E_{49} = 1, n_0 = 1$	$\beta_1 \approx 1, E_{49} = 10, n_0 = 1$
EVLA	1.4	0.25	7	1 Gpc	3.3 Gpc
ASKAP	1.4	30	30	500 Mpc	1.6 Gpc
MeerKAT	1.4	1.5	35	500 Mpc	1.6 Gpc
Apertif	1.4	8	50	400 Mpc	1.25 Gpc
LOFAR	0.15	20	1,000	35 Mpc	90 Mpc

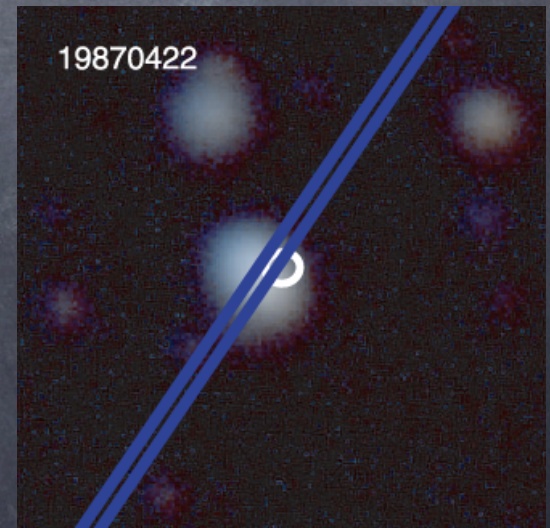
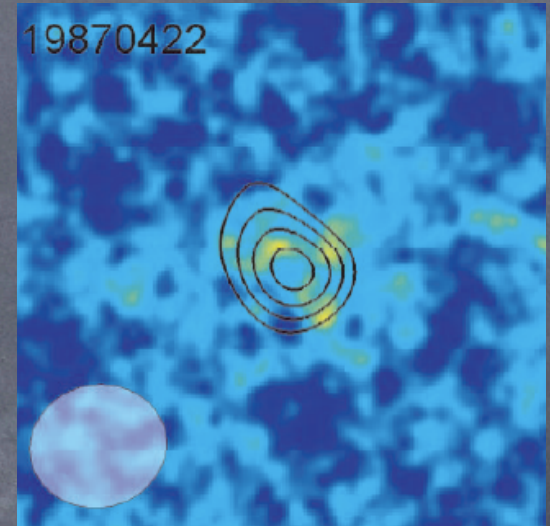
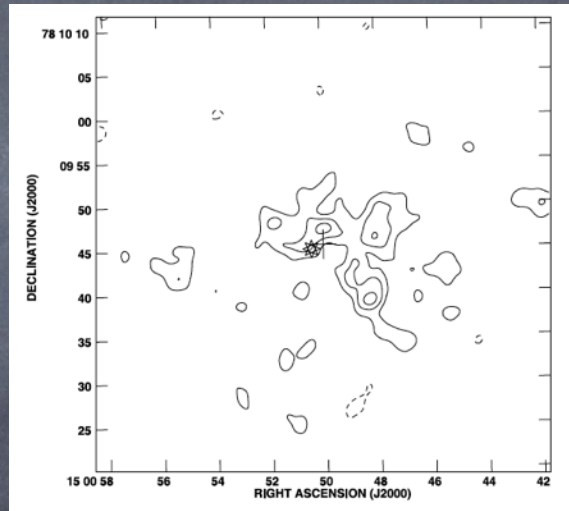
Ten-hour detection horizon	
$\beta_1 = 0.2, E_{49} = 10, n_0 = 1, p = 2.5$	$\beta_1 \approx 1, E_{49} = 1, n_0 = 10^{-3}, p = 2$
370 Mpc	140 Mpc
180 Mpc	70 Mpc
165 Mpc	65 Mpc
140 Mpc	50 Mpc
70 Mpc	20 Mpc

The Bower Transient

19870422



5GHz 0.5mJy
(<0.036 mJy)
 $t_{\text{next}} = 96$ days
1.5" from the
centroid of
MAPS-
P023-0189163
a blue Sc
galaxy at
 $z=0.249$
(1050Mpc) with
current star
formation

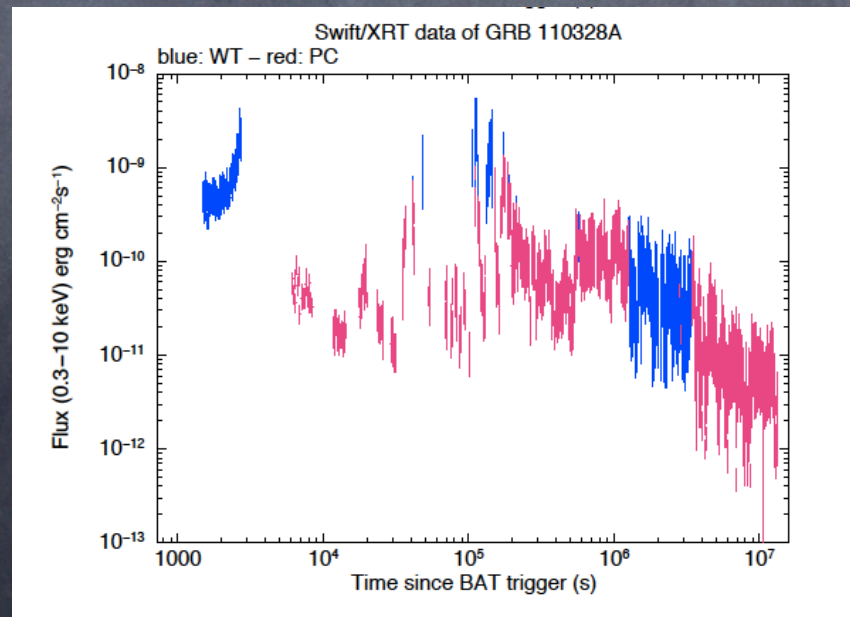




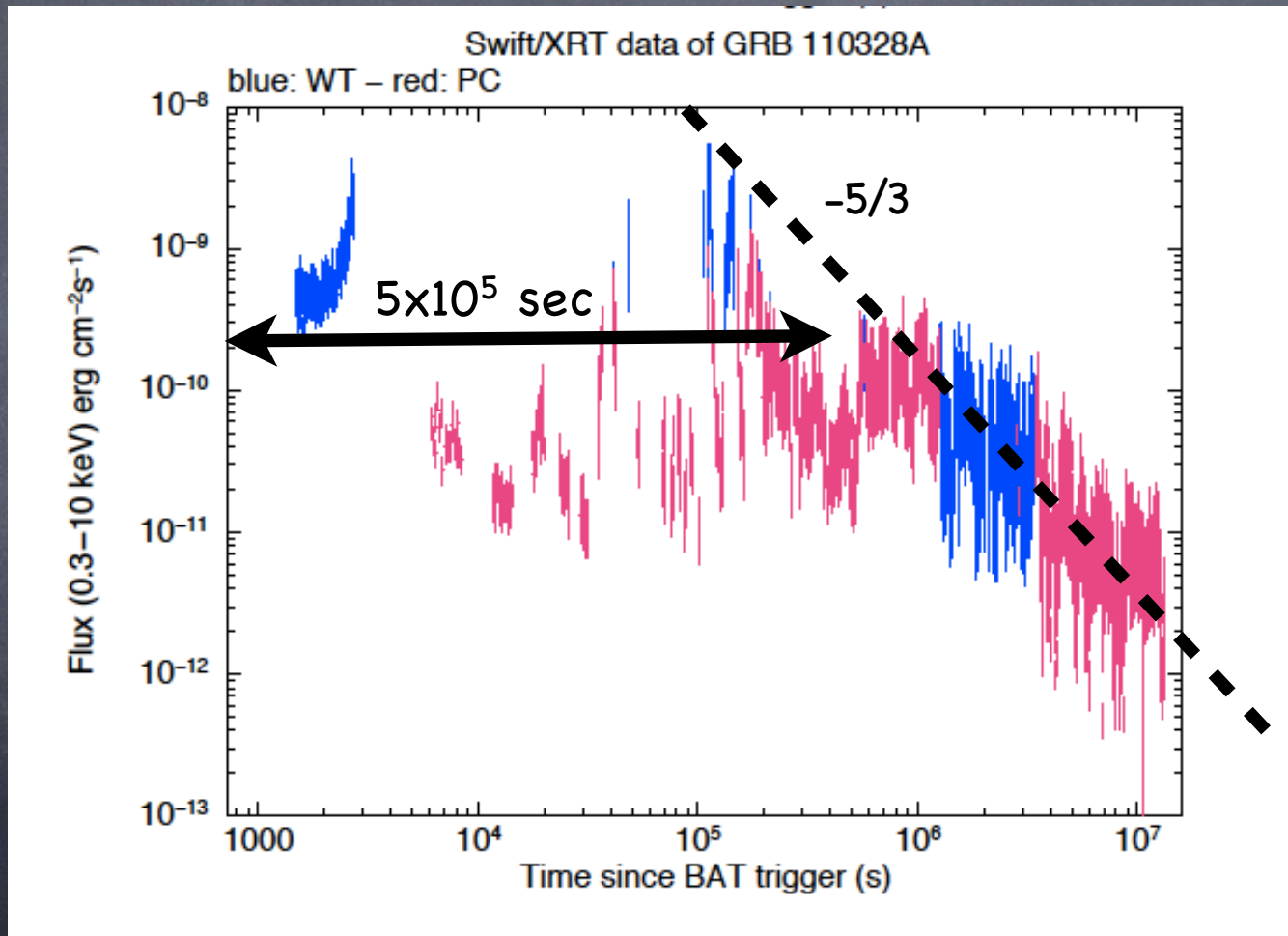
ERROR!

III. Tidal Disruption Events – Swift J1644 (GRB110328) and J2058

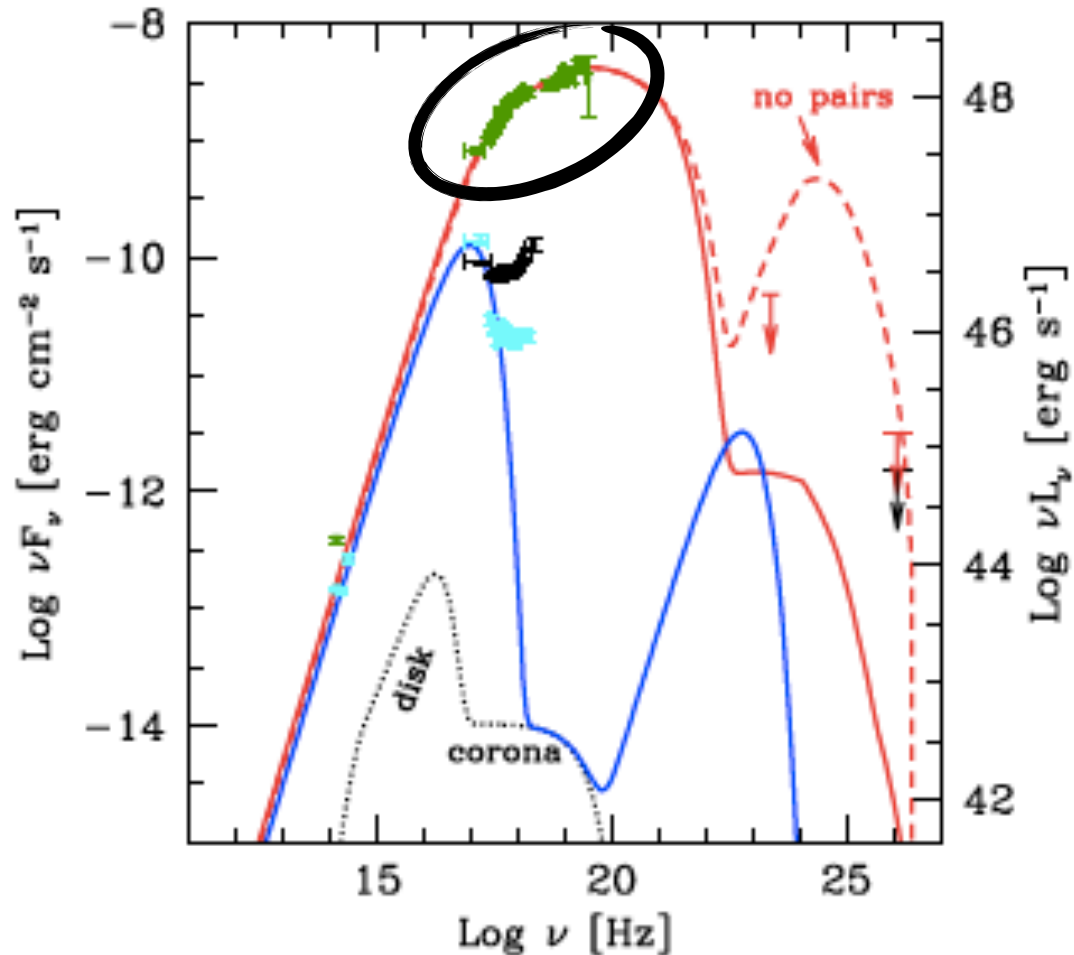
TP + Julian Krolik (ApJ. in press +
arXiv1111.2802)

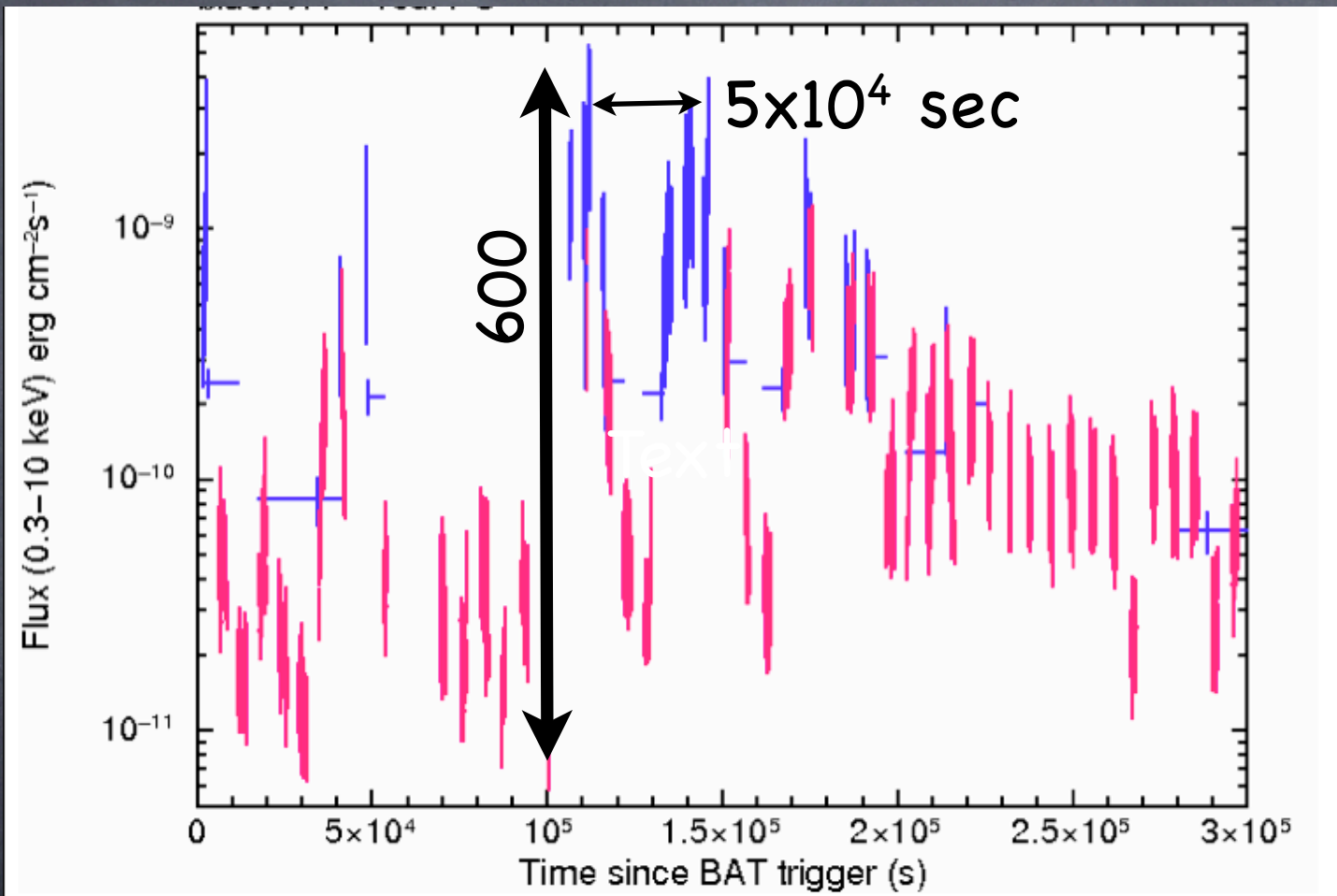


Ligth Curve



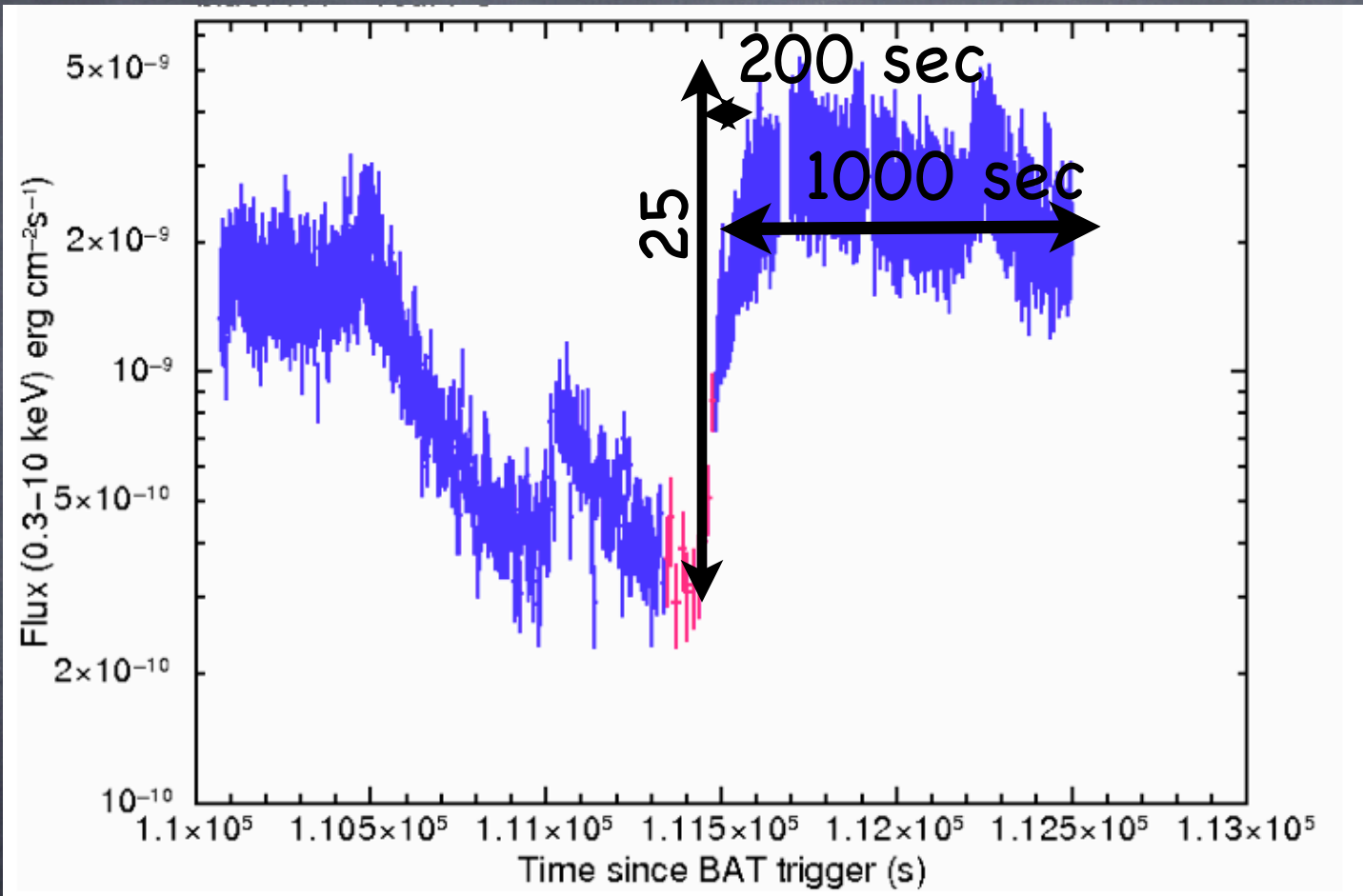
SED (Burrows et al.)





Swift light curve on a linear scale

The Third Flare



Light curve from 1.1×10^5 to 1.13×10^5 sec

Temporal features:

- Strong variability on 100 sec time scale
- Flares last about 1000–2000 sec
- Minima between the flares is a factor of 600 below the maxima
- 3×10^4 sec between flares
- 2×10^5 sec – duration before onset of a gradual decay

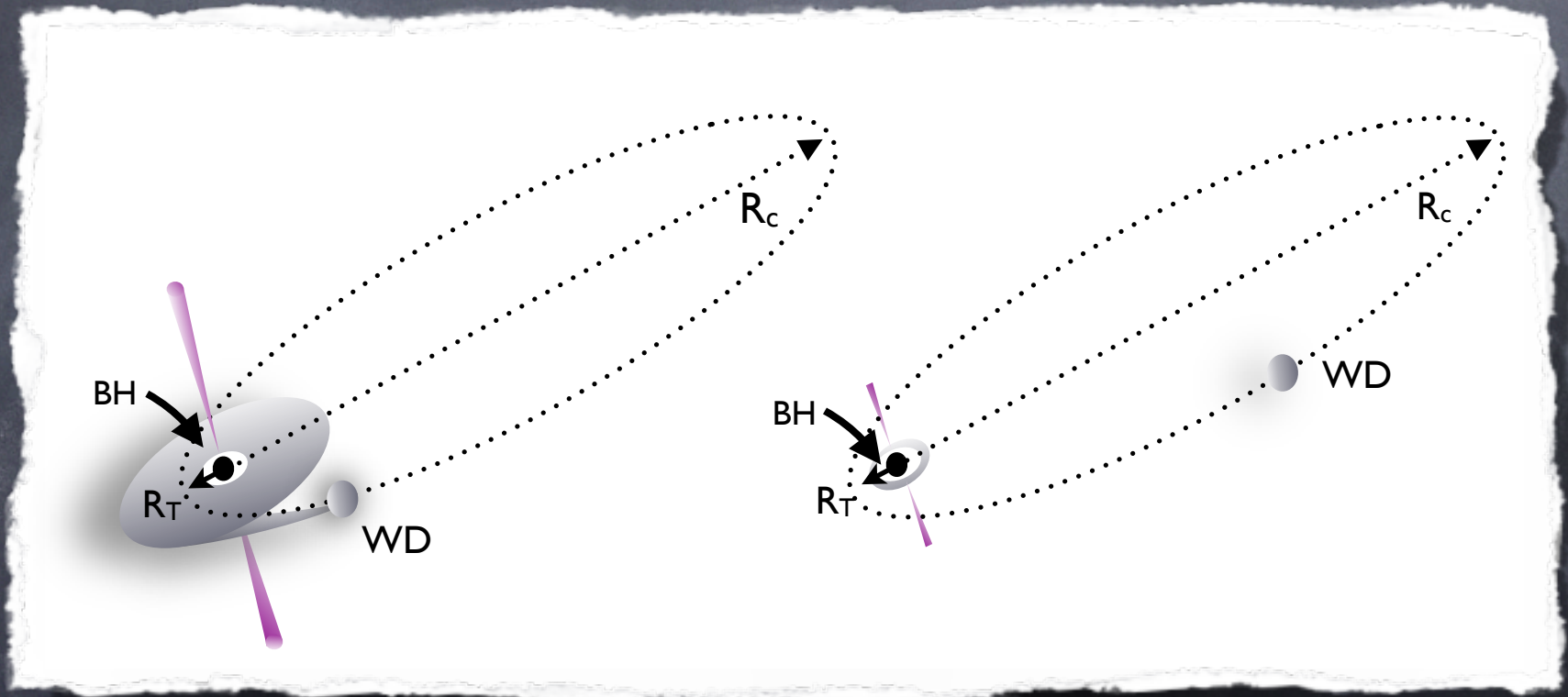
A tidal disruption of a main sequence star

The minimal relevant time scale:

$$P_{\text{orb}}(R_T) \approx 1/\sqrt{(G \rho)} \approx 10^4 \text{ sec}$$

Impossible to get 200 sec variability

A Disruption of a White Dwarf by a $5 \cdot 10^5 M_{\odot}$ black hole



$$P_{\text{orb}}(R_T) \approx 1/\sqrt{(G \rho)} \approx 6 \text{ sec}$$

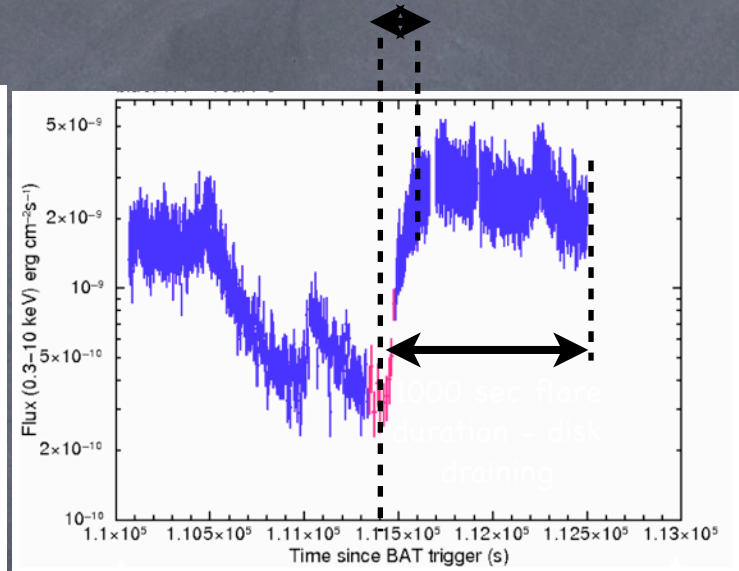
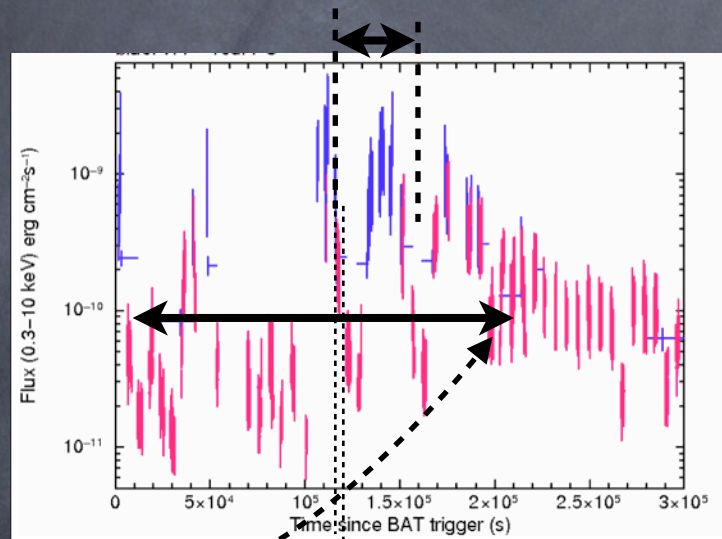
For a White Dwarf

$$P_{\text{orb}}(R_T) \approx 1/\sqrt{(G \rho)} \approx 6 \text{ sec}$$

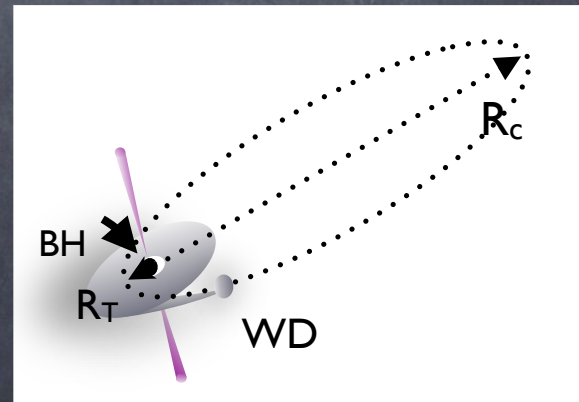
- 100 sec rise time – onset of accretion
- 1000 sec flare duration – the “drainage” time of a small accretion disk forms in a partial disruption event.
- 5×10^4 sec between flares – orbital time
- Precurse three days before the event is the “first” tidal passage

5×10^4 sec-
orbital period of
WD remnant

200 sec rise
time - a few
RT orbits



2×10^5 sec - onset of
 $t^{-5/3}$ decay



Why Jet?

- Blandford Znajek Jet power is determined by Magnetic field (B) on the horizon and the BH's area.
- B is determined by P (Pressure around ISCO)
- P depends on accretion mode – super or sub Eddington which depends on BH mass M_{BH} accretion rate $\dot{M} \propto (M_*/M_{\text{BH}})^{1/2}$
- Thermal (UV) emission also depends on accretion mode.

Sub-Eddington

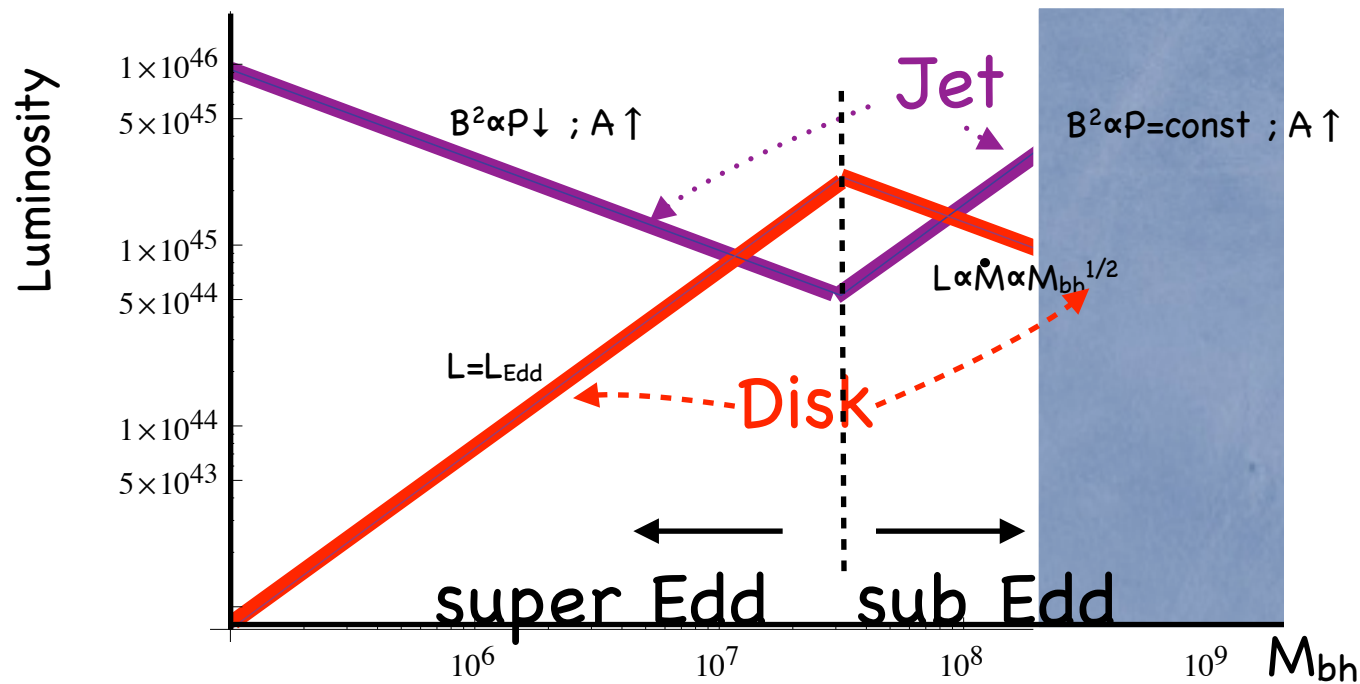
P is independent of \dot{M}

radiation dominate

Super-Eddington

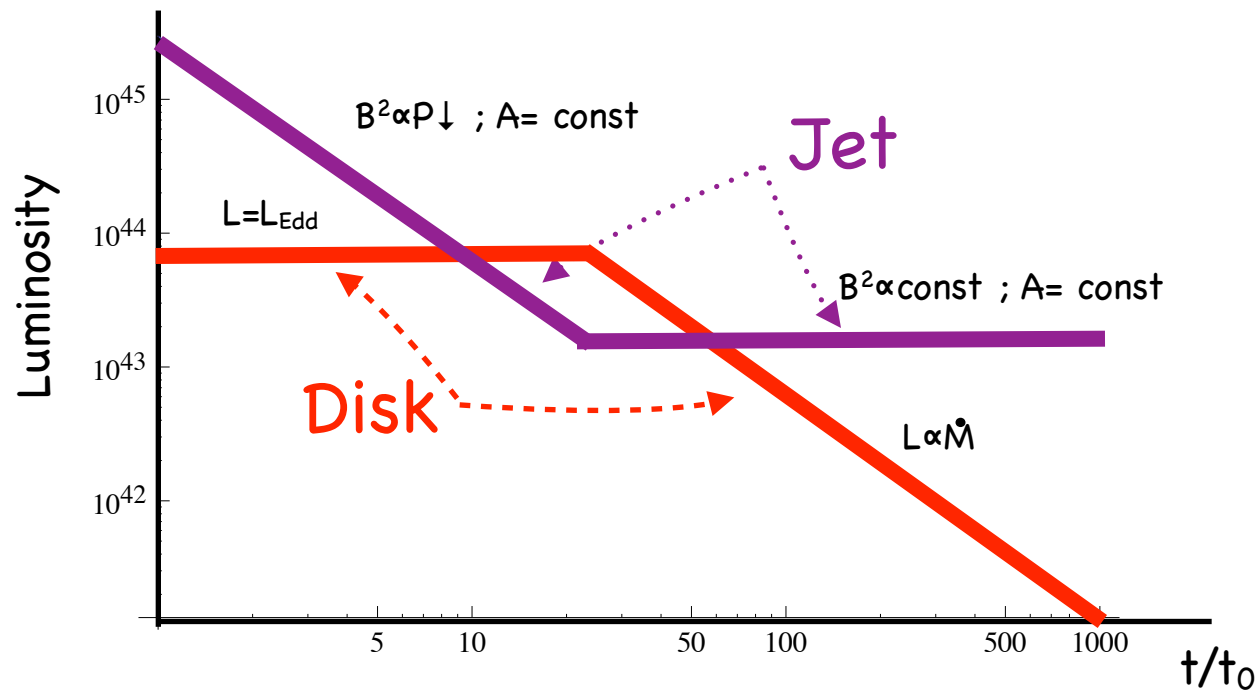


Jet (non thermal) vs Disk (thermal UV) Luminosity

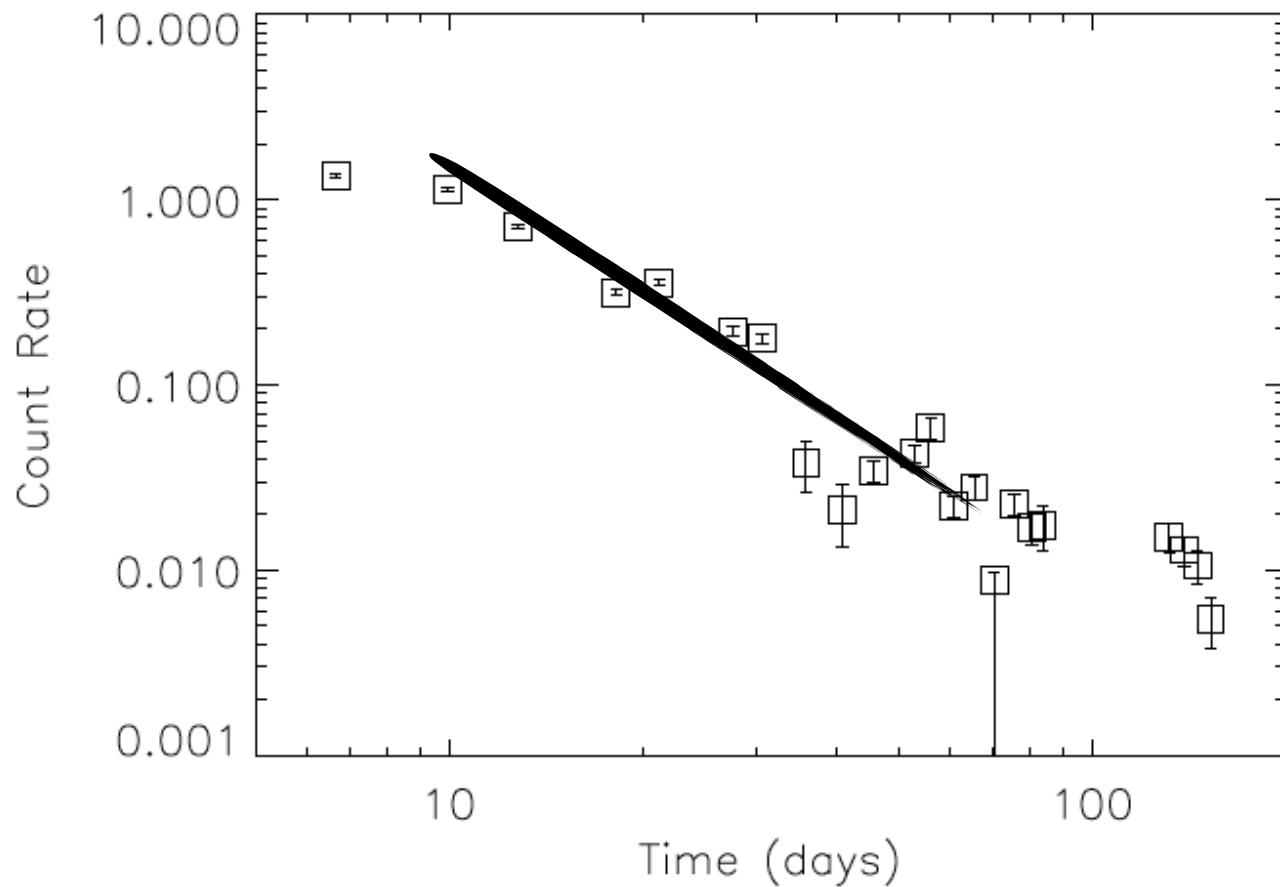


Light curves

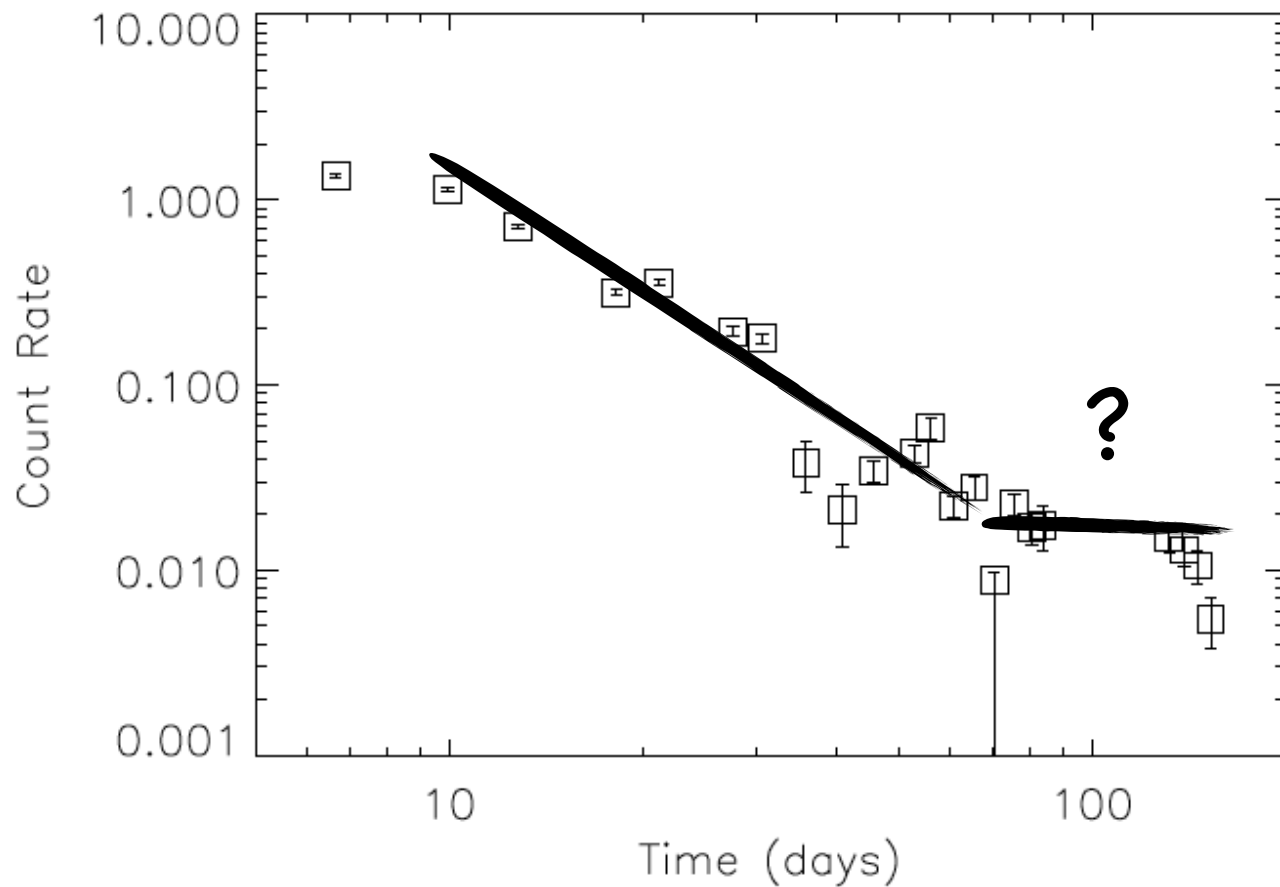
$$t_0 \simeq 1 \times 10^7 (\mathcal{M}_* M_{BH,6})^{1/2} \text{ s},$$



The x-ray light curve of Swift J2508



The x-ray light curve of Swift J2508

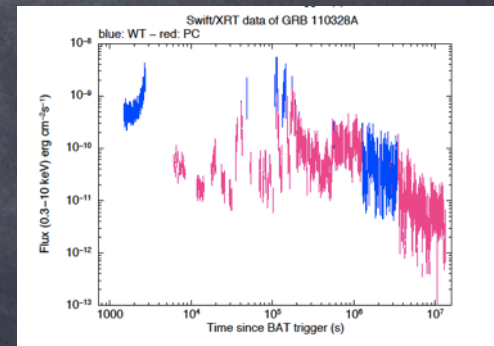
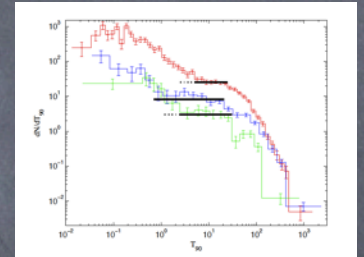


Can we apply a similar
reasoning to Neutrino
Dominated Accretion
Disks in GRBs?

Kawanaka & TP in preparation 2011

Summary

- There is a third population of GRBs – low luminosity GRBs – llGRBs – that arise from a different physical mechanism
- The observed plateau in the duration distribution of LGRBs show that LGRBs arise from Collapsars!
- A large fraction ($\sim 1/3$) of Swift short GRBs are Collapsars (only those with less than 0.7 sec are clearly Non-Collapsars).
- Strong Radio Flares should follow Neutron Star mergers
- Swift J1644 was a disruption of a WD
- Super Edd \rightarrow Jets \rightarrow x-rays in TDEs



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