The particle-without-particle approach to the self-force problem

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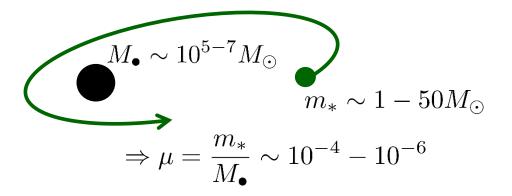


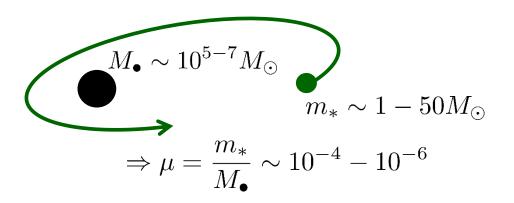


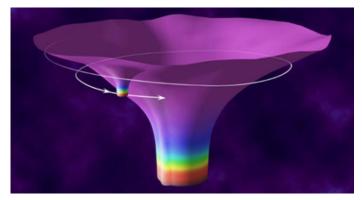






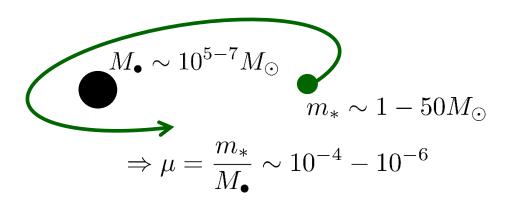


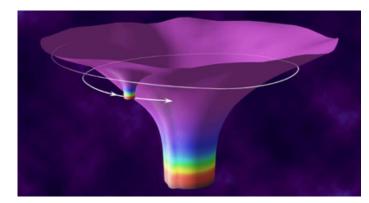




[arXiv:1305.5720]

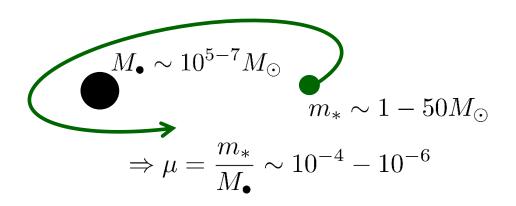
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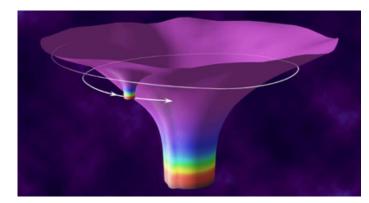




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- The theoretical formalism to compute the self-force has been largely established, e.g. [S. Gralla and R. Wald, CQG 25, 205009 (2008)], but its mathematical implementation is still under development; we use the *Particle-without-Particle (PwP)* method.

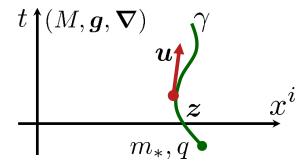




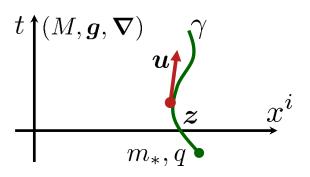
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- A helpful testbed for the gravitational self-force is the *scalar* self-force—we tackle this using the PwP method in the *frequency domain*.

• **Setup**: m_* is a charged scalar particle (with charge q associated to a scalar field Φ) orbiting a non-rotating black hole (of fixed—Schwarzschild—geometry, $(M, \boldsymbol{g}, \boldsymbol{\nabla})$) along a geodesic γ with worldline $\boldsymbol{z}(\tau)$ and 4-velocity $\boldsymbol{u} = \dot{\boldsymbol{z}}$. The EOMs are:



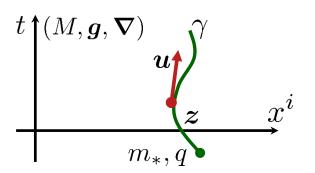
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$$\nabla^2 \Phi = -4\pi q \int_{\gamma} d\tau \, \delta(\boldsymbol{x} - \boldsymbol{z}(\tau)) ,$$
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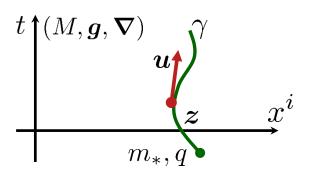


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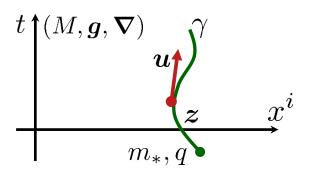
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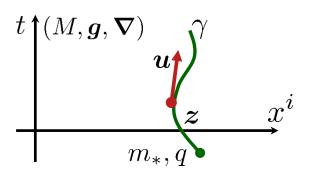
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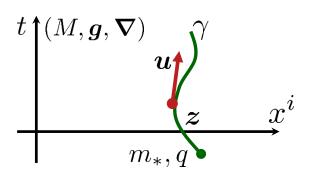
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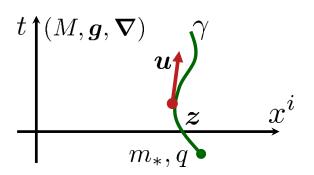
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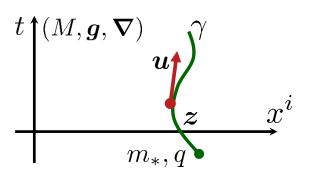
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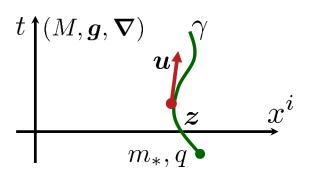
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Regge-Wheeler potential

proportional to q, dependent on \boldsymbol{z}

particle's radial location

redefined ("tortoise") radial coordinate

• Once the field is solved for, its singular part must be subtracted (via "mode-sum regularization" [L. Barack and A. Ori, PRD 61, 061502 (2000)]).

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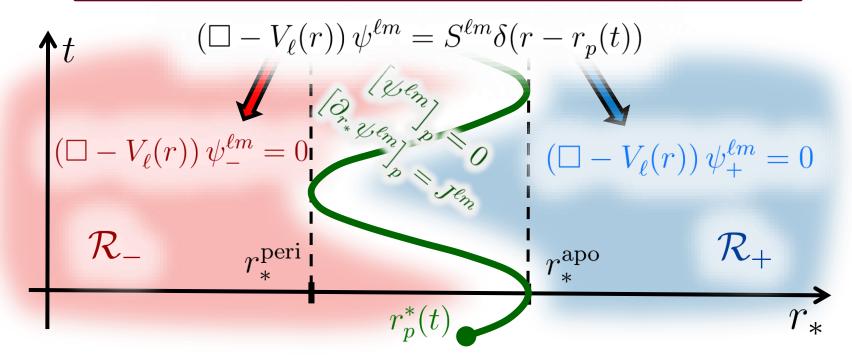
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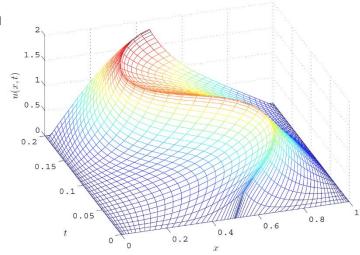
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- **Example**: Heat equation with constant source supported at a sinusoidally moving particle (using a Chebyshev-Lobatto grid in space and a standard finite-difference scheme in time).



• The PwP has already been successfully used to compute the scalar self-force in the *time domain* [P. Cañizares and CFS, PRD 79, 084020 (2009); CQG 28, 134011 (2011)].

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- Move to the *frequency domain*! We have bound orbits, and thus discreet series:

$$\psi_{\pm}^{\ell m}(t,r) = e^{-\mathrm{i}m\omega_{\varphi}t} \sum_{n=-\infty}^{+\infty} e^{-\mathrm{i}n\omega_{r}t} R_{\ell mn}^{\pm}(r).$$

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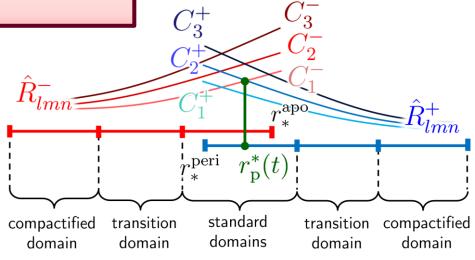
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We use a pseudospectral collocation method to find the homogeneous numerical solutions \hat{R}_{lmn}^- and \hat{R}_{lmn}^+ for arbitrary BCs, then use the jump conditions to get the true solution,

$$R_{lmn} = C_{lmn}^{-} \hat{R}_{lmn}^{-} \Theta_{p}^{-} + C_{lmn}^{+} \hat{R}_{lmn}^{+} \Theta_{p}^{+}$$



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- For future work, another objective is to also extend this method to *rotating black holes (Kerr)*.

Thanks for your attention!